

# Linking the complementary evaporation relationship with the Budyko framework for ungauged areas in Australia

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**Abstract.** While the calibration-free complementary relationship (CR) has performed excellently in predicting terrestrial evapotranspiration ( $ET_s$ ), how to determine the Priestley-Taylor coefficient ( $\alpha_c$ ) is still questionable. In this work, we evaluated this highly utilizable method, which only requires atmospheric data, with in-situ flux observations and basin-scale water balance estimates ( $ET_{wb}$ ) in Australia, proposing how to constrain it with a traditional Budyko equation for ungauged locations. We found that the CR method with a constant  $\alpha_c$  transferred from fractional wet areas performed poorly in reproducing the mean annual  $ET_{wb}$  in unregulated river basins, and it underperformed sophisticated physical, machine-learning, and land surface models in closing grid-scale water balance. This problem was remedied by linking the CR method with a traditional Budyko equation that allowed an upscaling of the optimal  $\alpha_c$  from gauged basins to ungauged locations. The proposed CR-Budyko framework enabled us to reflect climate conditions in  $\alpha_c$ , leading to more plausible  $ET_s$  estimates in ungauged areas. The spatially varying  $\alpha_c$  conditioned by local climates made the CR method outperformed the three  $ET_s$  models in reproducing the grid-scale  $ET_{wb}$  across the Australian continent. We here argued that the polynomial CR with a constant  $\alpha_c$  could result in biased  $ET_s$ , and it can be constrained by local climate conditions for improvement.

## 1 Introduction

Evapotranspiration ( $ET_s$ ) plays a pivotal role in water and energy exchanges between the land and the atmosphere. On the global scale, more than 60% of terrestrial precipitation (P) returns to the atmosphere through plants' vascular systems and soil pores, while consuming over 70% of surface net radiation (Trenberth et al., 2007; 2009). Since it is tightly coupled with carbon cycles, abnormally low  $ET_s$  would indicate food insecurity and low ecosystem sustainability (Jasechko, 2018; Kyatengerwa et al., 2020; Pareek et al., 2020; Swann et al., 2016). In severe cases,  $ET_s$  limited by deficient soil moisture can lead to extreme heatwaves that further propagate the water deficit in space and time (Miralles et al., 2014; Mueller and Seneviratne, 2012; Schumacher et al., 2022).

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Despite great community efforts for sharing in-situ observations (e.g., Baldocchi, 2020; Novick et al., 2018),  $ET_a$  gauging networks are unevenly established over land surfaces and often subjected to error sources (e.g., unclosed energy balance) and limited data lengths (Ma et al., 2021). Inevitably, modeling approaches are needed to predict  $ET_a$  in ungauged or poorly gauged areas, or to characterize it on a long timescale in a large area. Hence, various approaches have been proposed including physical models (e.g., Martens et al., 2017; Zhang et al., 2016), machine-learning techniques (e.g., Jung et al., 2019; Tramontana et al., 2016), and conceptual land surface schemes (e.g., Guimberteau et al., 2018; Haverd et al., 2018).

Those modeling approaches typically require P data and land surface information (e.g., remote-sensing vegetation indices) to quantify available soil moisture to the vaporization process. However, due in part to uncertainty associated with P data (Sun et al., 2018) and model structures (Samaniego et al., 2017; Zhang et al., 2019), resulting  $ET_a$  estimates have shown substantial disparities. In the comprehensive intercomparison by Pan et al. (2020), for example, the 14 advanced land surface models generated the global mean  $ET_a$  varying widely between  $450 \text{ mm a}^{-1}$  and  $700 \text{ mm a}^{-1}$ . Such a large incongruity in modeled  $ET_a$  was also found by the earlier Global Soil Wetness Project (Schlosser and Gao, 2010), suggesting that an alternative method is necessary to circumvent the uncertainty sources.

A practical method to simulate  $ET_a$  without P data and land-surface schemes is the complementary relationship (CR) of evaporation (Bouchet, 1963). It uses the evident fact that the air over a water-limited surface amplifies its vapor pressure deficit (VPD), while this effect disappears when the same surface is amply wet (Chen and Buchberger, 2018; Ramirez et al., 2005; Zhou et al., 2019). Based on the atmospheric self-adjustment, numerous equations have been formulated to predict  $ET_a$  only using routine meteorological data (e.g., Anayah and Kaluarachchi, 2014; Crago and Crowley, 2005; Crago and Qualls, 2013; Hobbins et al., 2004; Huntington et al., 2011; Kahler and Brutsaert, 2006, among others). In particular, the definitive derivation by Brutsaert (2015) and the following modifications (Crago et al., 2016; Crago and Qualls, 2021; Szilagyi, 2021; Szilagyi et al., 2017) provided strong physical foundations to Bouchet's (1963) early principle. They have excellently predicted  $ET_a$  at various spatial and temporal scales (e.g., Brutsaert et al., 2017, 2020; Crago and Qualls, 2018; Ma et al., 2019, 2021; Ma and Szilagyi, 2019), and allowed users to assess vegetation droughts over national and continental areas (e.g., Kim et al., 2019, 2021; Kyatengerwa et al., 2020).

Nevertheless, the definitive CRs still require at least some  $ET_a$  data to calibrate the parameters that determine the hypothetical wet-surface evaporation ( $ET_w$ ; Qualls and Crago, 2020); thus, they are not fully free of P data or parameterization. For instance, Brutsaert et al. (2020) calibrated the single parameter of Brutsaert's (2015) CR with flux observations and basin-scale P and runoff (Q) data to estimate annual  $ET_a$  across the globe. For evaluating four definitive CRs from Brutsaert's (2015) derivation, Crago et al. (2022) also calibrated their parameters by eddy-covariance flux observations. To date, Szilagyi et al. (2017) has proposed the only CR formulation that purely uses routine meteorological data; however, it depends on a questionable assumption that the parameter for  $ET_w$  is constant over a large continental area, being counterfactual to experimental studies on the Priestley and Taylor (1972) coefficient (e.g., Assouline et al., 2016; Baldocchi et al., 2016; Parlange and Katul, 1992; Wang et al., 2014). Given the complex space-time links between climate, soil, and vegetation (Hagedorn et

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al., 2019; Mekonnen et al., 2019; Rodriguez-Iturbe, 2000), the aerodynamic component of  $ET_w$  is unlikely represented by a fixed fraction of the net radiation.

Owing to the data required for parameter calibration, the state-of-the-art CR formulations might not be applicable in ungauged locations. In part, this problem can be mended by an additional constraint for determining the essential parameters, and the traditional Budyko framework can come into play. A Budyko function (e.g., Fu, 1981; Yang et al., 2008) explains the mean ratio of  $ET_a$  to P (i.e., surface water balance) simply by climatological aridity and a few implicit parameters, simultaneously closing the surface energy budget (Mianabadi et al., 2020). Although Bouchet's principle has often been linked with the water balance describe by Budyko functions (e.g., Carmona et al., 2016; Chen and Buchberger, 2018; Lhomme and Moussa, 2016; Zhang and Burtsaert, 2021), this theoretical link has been ignored when predicting  $ET_a$  by the definitive CRs. Kim and Chun (2021) explicitly showed that the atmospheric self-adjustment is tightly coupled with the climatological aridity within a Budyko function. This implicates that the optimal parameter for a definitive CR should vary with climates rather than staying constant.

In this work, we showed that a Budyko equation could become an important physical constraint when predicting  $ET_a$  by a definitive CR over a continental area. Here, a practical approach was proposed to determine the parameters reasonably in ungauged locations via a case study for the Australian continent, where the performance of the CR method remained unknown in many parts. Based on the analytical relationship between the CR and the Budyko framework, we showed why the parameter of the CR is not independent of local climate conditions, and addressed how to reflect spatially varying climates in its essential parameter.

## 2 Methodology and data

### 2.1 The polynomial CR by Szilagyi et al. (2017)

For the case study, we employed the calibration-free CR formulated by Szilagyi et al. (2017). It describes the atmospheric self-adjustment to surface moisture conditions using three evaporation rates, namely,  $ET_a$ ,  $ET_w$ , and the potential evaporation ( $ET_p$ ).  $ET_a$  is the actual moisture flux from a land surface to the atmosphere, and  $ET_w$  is the hypothetical  $ET_a$  rate that should occur with ample water availability.  $ET_p$  is the atmospheric capacity to receive water vapor that responds actively to soil moisture conditions. By defining the two dimensionless variables,  $x \equiv ET_p/ET_p$  and  $y \equiv ET_a/ET_p$ , Szilagyi et al. (2017) derived a polynomial function from four definitive boundary conditions.

Under ample water conditions,  $ET_p$  does not deviate from  $ET_w$  and  $ET_a$  (i.e.,  $ET_p = ET_w = ET_a$ ); hence, the corresponding zero-order boundary condition is (i)  $y = 1$  for  $x = 1$ . In contrast,  $ET_a$  must be nil over a desiccated surface (i.e.,  $y = 0$ ), and by energy balance, the surface net radiation should be fully transformed to the sensible heat flux. Then, the atmospheric VPD would be amplified at the maximum level with the same net radiation and wind speed. Defining the maximum  $ET_p$  rate as  $E_{pmax}$ , another zero-order boundary condition is given as (ii)  $y = 0$  for  $x = x_{min} \equiv ET_w/E_{pmax}$ . When  $x = 1$  (i.e., ample water), changes in  $ET_a$  would be controlled by changes in  $ET_w$ , yielding a first-order boundary condition as: (iii)

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dy/dx = 1 for x = 1. Over a desiccated surface,  $ET_a$  stays at zero even when  $ET_w$  or  $ET_p$  changes; thus, another first-order boundary condition becomes (iv) dy/dx = 0 for x = 0. The simplest polynomial equation satisfying the four boundary conditions is:

$$y = 2X^2 - X^3, \quad (1a)$$

where, X rescales the variable x into [0, 1] as:

$$X = \frac{x-x_{\min}}{1-x_{\min}} = \frac{E_{p\max}-ET_p-ET_w}{E_{p\max}-ET_w-ET_p} \quad (1b)$$

Eq. (1) allows users to estimate  $ET_a$  with no land-surface information, because  $ET_p$ ,  $ET_w$ , and  $E_{p\max}$  are all obtainable from a set of net radiation, air temperature, dew-point temperature, and wind speed data.  $ET_p$  and  $E_{p\max}$  can be estimated by the Penman (1948) equation:

$$ET_p = \frac{\Delta(T_a) R_n}{\Delta(T_a) + \gamma \lambda_v} + \frac{\gamma}{\Delta(T_a) + \gamma} f_u VPD, \quad (2)$$

$$E_{p\max} = \frac{\Delta(T_{dry}) R_n}{\Delta(T_{dry}) + \gamma \lambda_v} + \frac{\gamma}{\Delta(T_{dry}) + \gamma} f_u e_s(T_{dry}), \quad (3)$$

where,  $\Delta(\cdot)$  is the slope of the saturation vapor pressure curve (kPa °C<sup>-1</sup>),  $T_a$  is the mean air temperature (°C),  $\gamma$  is the psychrometric constant (kPa °C<sup>-1</sup>),  $R_n$  is the surface net radiation less the soil heat flux (MJ m<sup>-2</sup> d<sup>-1</sup>),  $\lambda_v$  is the latent heat of vaporization (MJ kg<sup>-1</sup>),  $f_u = 2.6 (1 + 0.54 u_2)$  is the Rome wind function (mm d<sup>-1</sup> kPa<sup>-1</sup>), where  $u_2$  is the 2-m wind speed (m s<sup>-1</sup>), and VPD is calculated by  $e_s(T_a)$  minus  $e_s(T_{dew})$ , where  $e_s(\cdot)$  is the saturation vapor pressure (kPa) and  $T_{dew}$  is the dew point temperature (°C).

$T_{dry}$  in Eq. (3) is the air temperature (°C) at which the lower atmosphere is devoid of humidity presumably by the adiabatic drying process:

$$T_{dry} = T_{wb} + \frac{e_s(T_{wb})}{\gamma} = T_a + \frac{e_s(T_{dew})}{\gamma}, \quad (4)$$

where,  $T_{wb}$  is the wet-bulb temperature (°C) at which the saturation vapor pressure curve intersects with the adiabatic wetting line. Thus, it is obtained by:

$$\gamma \frac{T_{wb} - T_{avg}}{e_s(T_{wb}) - e_a} = -1. \quad (5)$$

To estimate  $ET_w$  in Eq. (1b), the Priestly-Taylor (1972) equation has been a typical choice (e.g., Brutsaert, 2015; Crago et al., 2016; Han and Tian, 2018; Szilagyi et al., 2017):

$$ET_w = \alpha_e \frac{\Delta(T_w) R_n}{\Delta(T_w) + \gamma \lambda_v}, \quad (6)$$

where,  $\alpha_e$  is the Priestley-Taylor coefficient ranging usually within [1.10, 1.32] (Szilagyi et al., 2017), and  $T_w$  is the wet-environment air temperature (°C).  $T_w$  can be approximated with the wet-surface temperature ( $T_{ws}$ ), because the vertical air temperature gradient is negligible under a wet environment. Given its independence on areal extent (Szilagyi and Schepers, 2014),  $T_{ws}$  can be approximated by the implicit Bowen ratio ( $\beta$ ) of a small wet patch:

$$\beta = \frac{R_n - ET_p}{ET_p} \approx \gamma \frac{T_{ws} - T_a}{e_s(T_{ws}) - e_s(T_{dew})}. \quad (7)$$

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Eq. (7) assumes that the available radiation for the wet patch is close to that of the drying surface (Szilagyi et al., 2017).  $T_{ws}$  might be higher than  $T_a$  when the air is close to saturation. In such a case,  $T_{ws}$  should be capped by  $T_a$  when calculating  $ET_w$ .

The single parameter of the polynomial CR, i.e.,  $\alpha_e$ , is analytically obtainable by inserting the Priestley-Taylor equation into the Bowen ratio of a wet environment (Szilagyi et al., 2017):

$$\alpha_e = \frac{[\Delta(T_a) + \gamma](e_s(T_{ws}) - e_s(T_{dew}))}{\Delta(T_a)[(e_s(T_{ws}) - e_s(T_{dew})) + \gamma(T_{ws} - T_a)]} \quad (8)$$

where,  $\alpha_e$  must fall within the theoretical limit of  $[1, 1 + \gamma/\Delta(T_a)]$  (Priestley and Taylor, 1972).

## 2.2 The analytical relationship between the polynomial CR and a Budvko function

Since Eq. (8) is applicable only in a wet environment, Szilagyi et al. (2017) identified wet locations in a continental area based on the fact that the air close to saturation should have high relative humidity (RH) with  $T_{ws} > T_a$ . Thus, they calculated  $\alpha_e$  values at locations with  $RH > 90\%$  and  $T_{ws} > T_a + 2$  °C, and the average value was used to predict  $ET_a$  for a continental area. However, the spatially constant  $\alpha_e$  is unlikely suitable in such a large area under diverse climates, because the equilibrium between the atmosphere and the underlying surface is intertwined with the partitioning of P to  $ET_a$  and Q over the surface.

Kim and Chun (2021) analytically related Eq. (1) with the traditional Turc-Mezentsev equation, and found that the self-adjustment of  $ET_a$  (i.e.,  $x$ ) is tightly linked with climatological aridity and land properties. For the independence between P and 'the possible maximum  $ET_a$ ' of the Budyko framework, Kim and Chun (2021) reformulated the traditional equation with  $\Phi_0 \equiv ET_w/P$  instead of the commonly used aridity index ( $\Phi \equiv ET_p/P$ ) as:

$$\frac{ET_a}{P} = \frac{ET_w}{P} \left[ \frac{1}{1 + \left(\frac{ET_w}{P}\right)^n} \right]^{\frac{1}{n}} = \frac{xET_p}{P} \left[ \frac{1}{1 + \left(\frac{xET_p}{P}\right)^n} \right]^{\frac{1}{n}} \quad (9)$$

where, the parameter  $n$  implicitly represents the factors affecting the P partitioning other than the climatic drivers. By dividing Eq. (9) by  $\Phi$ , it is found that the Budyko equation (9) is intertwined with the Eq. (1a):

$$y = \frac{ET_a}{ET_p} = 2X^2 - X^3 = \left[ \frac{x^n}{1 + x^n \Phi^n} \right]^{\frac{1}{n}} \quad (10)$$

Eq. (10) implies that the self-adjustment of  $ET_a$  (i.e.,  $x$ ) is tightly related with the climatic condition (i.e.,  $\Phi$ ), and the land property (i.e.,  $n$ ).

While the  $x$  and  $n$  can be achievable from a set of  $ET_a$ ,  $ET_p$ ,  $E_{pmax}$ , and P values by inverting Eq. (10), such an approach is not applicable in locations with no  $ET_a$  data. To quantify  $x$  values only using  $ET_p$ ,  $E_{pmax}$ , and P, Kim and Chun (2021) developed a regression equation between  $x$  and  $\Phi$ ,  $x_{min}$ , and  $n$  values from the 513 gauged river basins over the world. We used the same regression-based regionalization. Considering  $x_{min} = xET_p/E_{pmax}$ , the non-linear Eq. (10) can be approximated by a multiple regression as:

$$x = b_0 + b_1 \ln(\Phi) + b_2 \ln(ET_p/E_{pmax}) + b_3 \ln(n) \quad (11)$$

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where,  $x_e$  is the approximate ratio of  $ET_w$  to  $ET_p$ , and  $b_0$ ,  $b_1$ , and  $b_2$  are the intercept and the regression coefficients, respectively.

Since the implicit parameter  $n$  is unavailable in ungauged locations, Eq. (11) needs to be further simplified by neglecting the last term:

$$x_e \approx c_0 + c_1 \ln(\Phi) + c_2 \ln(ET_p/E_{pmax}), \quad (12)$$

where,  $c_0$ ,  $c_1$ , and  $c_2$  are the intercept and the coefficients of the approximated regression.

If  $x_e$  is known by the regression Eq. (12), the parameter  $\alpha_e$  can be estimated using the Priestley-Taylor equation as:

$$\alpha_e = x_e \frac{ET_p}{ET_{eq}} \quad (13a)$$

$$ET_{eq} = \frac{\Delta_e(T_w) - E_n}{\Delta_e(T_w) + \gamma \lambda_v} \quad (13b)$$

where,  $\alpha_e$  is the Priestley-Taylor coefficient that approximately satisfies the CR and the Budyko equations together, and  $ET_{eq}$  is the equilibrium evapotranspiration (mm d<sup>-1</sup>) at which VPD is nil under a wet environment. It should be noted that  $P$ ,  $ET_p$ ,  $E_{pmax}$ , and  $ET_{eq}$  within Eqs. (9)-(13) must be on a timescale where the Turc-Mezentsev equation is valid (typically longer than a year), and  $\alpha_e$  is still bounded by [1, 1+ $\gamma/\Delta(T_s)$ ].

### 2.3 Atmospheric forcing, eddy-covariance, and runoff data

We examined the CR-Budyko combined framework in the Australian continent lying within [10°-45°S, 113°-155°E]. The required atmospheric forcing data ( $R_n$ ,  $T_{ds}$ ,  $T_{dew}$ , and  $u_2$ ) were collected from the advanced ERA5-Land reanalysis archive (Muñoz-Sabater et al., 2021) of the European Centre for Medium-Range Weather Forecasts (<https://cds.climate.copernicus.eu>; last access on Dec-10/2021). The monthly averages of surface latent and sensible heat fluxes, 2-m air temperature, 2-m dew-point temperature, and 10-m U and V wind speed components at 0.1°×0.1° were downloaded for 1981-2020.  $R_n$  was calculated by summing the two heat fluxes, and the 10-m wind speed components were converted to  $u_2$  using the logarithmic wind profile (Allen et al., 1998).

We also collected the Australian edition of the Catchment Attributes and Meteorology for Large sample Studies (CAMELS; Fowler et al., 2021) series of datasets (available at <https://doi.org/10.1594/PANGAEA.921850>; last access on Sep-27/2021). The CAMELS datasets comprise daily time series of 19 hydrometeorological variables at 222 unregulated river basins in Australia up to 2014, and we selected the 71 basins larger than 500 km<sup>2</sup> to contain at least five CR  $ET_a$  estimates within the boundaries. The water-balance  $ET_a$  ( $ET_{wb}$ ) (i.e.,  $ET_{wb} \approx \Sigma P - \Sigma Q$ ) of each basin was calculated for the two periods of 1981-1997 and 1998-2014. The mean annual  $ET_{wb}$  for the former period was used for the regressions with Eqs. (11) and (12), and the predicted  $ET_a$  was evaluated against the latter.

As a point-scale evaluation dataset, the annual flux observations were taken from the 15 eddy-covariance stations (Table 1) of the FLUXNET2015 archive (<https://fluxnet.org/>; last access on Jul-1/2021). We chose the flux towers with 2 or more annual means, and adopted the energy-balance-corrected latent heat flux observations with the quality measures

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'LE\_F\_MDS\_QC', higher than 0.70. Given the fine resolution of the ERA5-Land forcing data, we believed that the ET<sub>a</sub> estimates by CR could be directly compared with the point-scale observations.

In addition, as a grid-scale evaluation reference, the SILO P data at 0.01°×0.01° were collected from the Queensland government (<https://www.longpaddock.qld.gov.au/silo/gridded-data>; last access on Jun-01/2021) together with the Global RUNoff (GRUN) ENSEMBLE (Ghiggi et al., 2021) (<https://doi.org/10.6084/m9.figshare.12794075>; last access on Oct-1/2021). The global Q data were produced at 0.5°×0.5° using a machine-learning algorithm trained by in-situ streamflow observations, and potential errors were reduced by simulations with 21 sets of atmospheric forcing (Ghiggi et al., 2021). The SILO P was used to calculate  $\Phi = P/ET_p$  at each grid of the forcing data. After bilinearly unifying the resolutions of SILO P and GRUN Q data, we also calculated the mean annual ET<sub>wb</sub> for 1998-2014 at 0.5°×0.5° over the entire Australian continent.

Against the grid-scale ET<sub>wb</sub> estimates, performance of the polynomial CR was also compared with three ET<sub>a</sub> products from a physical, a machine-learning, and a land-surface model. The physical model was the Global Land Evaporation Amsterdam Model (GLEAM) v3.2 (Martens et al., 2017; <https://www.gleam.eu>; last access on Jun-03/2020) based on the Priestley-Taylor equation constrained by microwave-derived soil moisture, surface temperature, and vegetation optical depth. The machine-learning ET<sub>a</sub> product was the FluxCom (<http://www.fluxcom.org/>; last access Mar-18/2019) that upscaled in-situ observations at 224 eddy-covariance towers using 11 algorithms (Jung et al., 2019). We used the version forced by the CRUNCEPv8 that has the longest data length from 1950 to 2016. The land-surface-model product was the ERA5-Land monthly ET<sub>a</sub> (<https://cds.climate.copernicus.eu>; last access on Jul-7/2021) simulated by the advanced Hydrology Tiled ECMWF Scheme for Surface Exchanges over Land scheme (Balsamo et al., 2015). All the modeled ET<sub>a</sub> datasets were bilinearly gridded to 0.5°×0.5° for 1998-2014 to be compared with the grid-scale ET<sub>wb</sub> data.

### 3 Results

#### 3.1 Performance of the calibration-free CR in Australia

Figure 1a depicts the spatial distribution of the inverted aridity index ( $\Phi^{-1} = P/ET_p$ ) that can traditionally categorize climate conditions. The mean ratios between SILO P and ET<sub>p</sub> for 1998-2014 indicated that 83% of the Australian land surfaces were under arid ( $\Phi^{-1} < 0.2$ ) and semi-arid climates ( $0.2 < \Phi^{-1} < 0.5$ ). Semi-humid ( $0.5 < \Phi^{-1} < 0.65$ ) and humid climates ( $\Phi^{-1} > 0.65$ ) were only found in the northern and southeastern coastal areas and the southwestern edge where major cities and agricultural lands have developed. Despite the high aridity, hyper-arid climates ( $\Phi^{-1} < 0.05$ ) were not found in Australia.

We first examined the calibration-free approach by Szilagyi et al. (2017) that only uses the meteorological forcing inputs. The blue-colored points in Figure 1a are the locations with RH > 90% and  $T_{ws} > T_a + 2^\circ\text{C}$ , at which the  $\alpha_c$  values from Eq. (8) were within  $1.15 \pm 0.047$  (mean  $\pm$  standard deviation). Though the two conditions were met in some mountainous areas in the southeastern part, we excluded them because unexpectedly high  $\alpha_c$  values were obtained. The mean  $\alpha_c = 1.15$  fell within

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the theoretical limits, and was equal to the value used in the prior studies in China (Ma et al., 2019) and the conterminous U.S. (Ma and Szilagyi, 2019).

Using the CR with  $\alpha_c = 1.15$ , we predicted  $ET_a$  over the entire Australian continent (Figure 1b). The distribution of the resulting mean  $ET_a$  for 1998-2014 was coherent with that of  $\Phi^{-1}$ . The mean CR  $ET_a$  ranged in  $262 \pm 85.3$  mm  $a^{-1}$  and  $547 \pm 173$  mm  $a^{-1}$  under arid and semi-arid climates, respectively. On the other hand, CR  $ET_a$  in semi-humid and humid locations were much higher in  $886 \pm 187$  mm  $a^{-1}$  and  $1,010 \pm 213$  mm  $a^{-1}$ , respectively. The calibration-free CR predicted the continental mean  $ET_a$  as high as  $489$  mm  $a^{-1}$  for 1981-2012, and it was about 11.3% higher than the estimate for the same period ( $439$  mm  $a^{-1}$ ) by Zhang et al. (2016). The mean fraction of  $ET_a$  to P for 1998-2014 (97%) was larger than the typical  $ET_a$  value in Australia (~90%; Glenn et al., 2011), implicating that the constant  $\alpha_c = 1.15$  seemed to make the CR overrate  $ET_a$ .

The overestimation of the calibration-free CR was confirmed by the flux observations and the basin-scale  $ET_{wb}$  (Figure 2). The percent bias (p-bias) of CR  $ET_a$  to the point-scale annual  $ET_a$  was +10.4%, while it became more than doubled when compared to the basin-scale  $ET_{wb}$ . Though the Pearson correlation coefficients (Pearson r) were significantly high between the CR  $ET_a$  and the two evaluation references, the low Nash-Sutcliffe efficiency (NSE) to  $ET_{wb}$  implicates that the CR method could perform poorly in wet river basins. The regression slopes in Figure 2 also indicate that the calibration-free CR tends to increasingly overestimate as climate becomes wetter. The root mean square error (RMSE) of CR  $ET_a$  to  $ET_{wb}$  was higher than to the point observations. Although it appeared to perform acceptably at the 15 flux towers, the CR method produced considerable biases in the 71 CAMELS basins. The performance measures were not as excellent as the same CR method had shown in the U.S. (Ma et al., 2021; Ma and Szilagyi, 2019; Kim et al., 2019) and in China (Ma et al., 2019).

One may argue that the mean  $\alpha_c$  derived from fractional wet areas is unlikely representative of the large Australian continent, and this might introduce the biases to CR  $ET_a$  estimates. Hence, we re-simulated CR  $ET_a$  with Ma et al.'s (2021) estimate ( $\alpha_c = 1.10$ ) from a global-scale analysis. Figure 3a shows that the predicted  $ET_a$  became nearly unbiased at the 15 flux tower locations, and seemingly suggests that the decreased  $\alpha_c$  could become a solution to improving the CR method. Nevertheless, the fixed  $\alpha_c$  still made the CR overestimate  $ET_a$  in the CAMELS basins under (semi-)humid climates, albeit slightly ameliorated (Figure 3b).

### 3.2 The empirical relationship between $x$ to climate conditions

Figures 2 and 3 imply that the calibration-free CR with a fixed  $\alpha_c$  was unlikely good at closing local water balance particularly in (semi-)humid river basins. To resolve this problem with the CR-Budyko framework, first we estimated the climatological  $x$  and the parameter  $n$  of the CAMELS basins using Eq. (10) with the mean annual  $ET_{wb}$ , P,  $ET_p$ , and  $E_{pmax}$  for 1981-1997. Figure 4a-c illustrates the scatter plots between the resultant  $x$  and corresponding  $\Phi$ ,  $ET_p/E_{pmax}$ , and  $n$  values. The Pearson r between the  $x$  and the other three variables was -0.88, -0.59, and 0.44, respectively (significant at 1% level), suggesting that the self-adjustment of  $ET_p$  is not only correlated with climate conditions, but with land surface properties at least in part. By regressing between the  $x$  values and the log-transformed  $\Phi$ ,  $ET_p/E_{pmax}$  and  $n$ , we obtained an empirical relationship that enables to spatially predict the mean annual ratio of  $ET_w$  to  $ET_p$  as:

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$$x = 0.949 - 0.204 \ln(\Phi) + 0.231 \ln(ET_p/E_{pmax}) + 0.0712 \ln(n). \quad (14)$$

The regression coefficients were all significant at 1% level, and the coefficient of determination ( $R^2$ ) was 0.98. The regression equation was further approximated by discarding  $n$  from the explanatory variables:

$$x = 1.023 - 0.220 \ln(\Phi) + 0.210 \ln(ET_p/E_{pmax}). \quad (15)$$

The  $R^2$  value of Eq. (15) declined to 0.93. We found that the simple regression between  $x$  and  $\Phi$  further reduced  $R^2$  to 0.90. While the heterogeneous land properties exert non-negligible influences, the regression analyses imply that the climatic condition dominantly explains the spatial variation of the atmospheric self-adjustment.

Eq. (15) performed excellently in reproducing the  $x$  values from CR with  $\Phi$  and  $ET_p/E_{pmax}$  (Figure 4d). The NSE, RMSE, Pearson  $r$ , and  $p$ -bias between the predicted  $x$  and the  $x$  from CR were 0.93, 0.03, 0.96, and 0.0%, respectively.

### 3.3 Evaluation of the CR and the advanced models against the grid $ET_{wb}$

By multiplying  $x$  to the mean annual ratio between  $ET_p$  and  $ET_{eq}$ , we determined  $\alpha_e$  across the Australian land surfaces. The resulting  $\alpha_e$  values ranged within  $1.13 \pm 0.114$ , and the median value was almost equal to Ma et al.'s (2021) global estimate (1.10). They were relatively high in the northwestern and the northern part, while being below the mean in the southern and the eastern parts (Figure 5a). On 19% of the surfaces,  $\alpha_e$  values were unity, and thus they might become below the theoretical limit unless bounded.

We again generated CR  $ET_a$  using the spatially varying  $\alpha_e$  values (Figure 5b). The mean CR  $ET_a$  for 1998-2014 ranged in  $249 \pm 78.8 \text{ mm a}^{-1}$  and  $530 \pm 172.0 \text{ mm a}^{-1}$  under arid and semi-arid climates, while it decreased to  $305.2 \pm 209 \text{ mm a}^{-1}$  and  $932 \pm 239 \text{ mm a}^{-1}$  in semi-humid and humid regions, respectively. The flux observations were still acceptably regenerated with the less biases than in the case of  $\alpha_e = 1.15$  (Figure 6a). The  $\alpha_e$  based on the Budyko framework significantly reduced the biases introduced by the constant  $\alpha_e$  in (semi-)humid basins. Albeit some biases remained, the water-balance  $ET_{wb}$  for 1998-2014 in the CAMELS basins were better reproduced by using the spatially varying  $\alpha_e$  (Figure 6b).

To confirm the improved performance of the combined CR-Budyko method across Australia, we resampled the new CR  $ET_a$  estimates to  $0.5^\circ \times 0.5^\circ$  and compared them with the grid  $ET_{wb}$  data. The  $ET_a$  products by GLEAM, FluxCom, and ERA5-Land were evaluated with the grid evaluation reference. As shown, the CR method with a constant  $\alpha_e = 1.15$  overrated the mean annual  $ET_a$  along the eastern and the northern coastlines (Figure 7b), underperforming the physical, the machine-learning, and the land surface models (Figure 8a). Although the smaller  $\alpha_e = 1.10$  made the CR method perform better, its predictability was still poorer than the three advanced models, and the residual variation was as large as in the case of  $\alpha_e = 1.15$  (Figure 8b).

In constant, when employing the  $\alpha_e$  conditioned by local climate conditions, the same CR formulation could alleviate the overestimation along the coastlines (Figure 7c). The Budyko-function-based  $\alpha_e$  led the CR  $ET_a$  estimates to neatly agree with the grid  $ET_{wb}$ , and the residual variance was much smaller than in the case of  $\alpha_e = 1.10$  (Figure 8c). The CR method with  $\alpha_e$  clearly outperformed the three advanced models in reproducing the grid  $ET_{wb}$  estimates (Figure 8d-f). Although the

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referenced, grid  $ET_{wb}$  has some error sources associated with upscaling of P and  $Q_e$ , our comparative evaluation suggests that conditioning  $\alpha_e$  with local climate conditions could substantially reduce the uncertainty of CR  $ET_a$  estimates in ungauged areas.

#### 4 Discussion

##### 4.1 Constraining the CR with the Budyko framework for ungauged areas

The CR explains the dynamic equilibrium between the atmospheric  $ET_a$  and the underlying moisture conditions, while the Budyko framework describes the steady-state water balance with climatic controls (i.e., P and  $ET_p$ ). The analytical link between the CR and the Budyko equations, hence, implies that the atmospheric self-adjustment needs to be conditioned by the long-term climate conditions. Constraining the Turc-Mezentsev equation by the polynomial CR, Kim and Chun (2021) found that Q changes would be more sensitive to climatic changes than when they were not linked. In the opposite direction, the CR can be constrained by the Budyko equation to determine its essential parameter

In Crago and Qualls (2018), the optimal  $\alpha_e$  for the linear CR of Crago et al. (2016) varied largely between 1.00 and 1.43. This point-scale experiment has already suggested that a constant  $\alpha_e$  is unlikely suitable for definitive CRs to predict  $ET_a$  in Australia. The ratio between the aerodynamic and the radiation components of  $ET_w$  is evidently affected by the heat entrainment from the top of the boundary layer (Baldocchi et al., 2016), the dissimilarity between heat and water vapor sources (Assouline et al., 2016), the large-scale synoptic changes (Guo et al., 2015), and the horizontal advection of dry air mass (Jury and Tanner, 1975). More recently, Han et al. (2021) proved the non-linear dependence of  $ET_w$  on  $ET_{eq}$ , and Yang and Roderick (2019) showed  $\alpha_e$  changing with  $R_a$  over ocean surfaces. Hence, the constant  $\alpha_e$  assumption underpinning the calibration-free CR is counterintuitive to the theoretical and empirical evidence. Although Ma et al. (2021) found some global applicability of the calibration-free CR, its performance remains unknown in most of the Australian surfaces and in many ungauged basins over the world.

Since  $ET_a$  plays a pivotal role in the terrestrial water and energy balances, the partitioning of  $R_a$  into the latent and the sensible heat fluxes cannot be independent of the partitioning of P into  $ET_a$  and Q. On a mean annual scale, P and  $ET_p$  are the major determinants of the P partitioning, and thus the parameter  $\alpha_a$  might not be independent of P. Given the large variability of P, assuming a fixed  $\alpha_e$  across a continental area may introduce considerable biases to CR  $ET_a$  estimates. Thus, discarding available P data may not be a good choice when predicting  $ET_a$  by the CR method in ungauged areas. It is noteworthy that  $\Phi$  dominantly explained the spatial variation of the mean annual x of the 71 CAMELS basins, and the  $\alpha_e$  values conditioned by local climates were of a large spatial variation. This suggests that the CR with a constant  $\alpha_e$  may produce unreliable  $ET_a$  estimates in ungauged locations.

Nonetheless, the low performance with a constant  $\alpha_e$  does not indicate that the CR method underperforms the sophisticated  $ET_a$  models. The simple polynomial CR seemed to outperform the advanced the advanced physical, machine-learning, and land surface models, when its parameter was conditioned by local climates. The proposed CR-Budyko framework

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enabled to regionalize the optimal  $\alpha_c$  for the CR method from gauged basins to ungauged locations in an empirical manner. It should be highlighted that the CR with spatially varying  $\alpha_c$  produced the much smaller residual variance than the three advanced models.

#### 4.2 Remaining issues and caveats

In seven Australian eddy-covariance flux towers, Crago et al. (2022) found that the optimal  $\alpha_c$  for the polynomial CR was 1.35 for predicting daily  $ET_a$  in the dimensionless form (i.e.,  $y = ET_p/ET_p$ ). However, it was increased to 1.42, 1.45, 1.47, and 1.50 to simulate the dimensional latent heat fluxes at daily, weekly, monthly, and annual timescales, respectively. This implies that the timescale would largely affect the optimal  $\alpha_c$  for the definitive CRs. Though the stationary Budyko equation can become a constraint at a mean-annual scale, how to capture the scale-dependence of  $\alpha_c$  is a remaining question.

Further questions also can arise as to how to quantify  $ET_p$  and  $E_{p,max}$ . For example, the  $\alpha_c$  values from  $ET_p$  with the Rome wind function rely upon an unrealistic assumption that the aerodynamic resistance on a vegetated surface is equivalent to that of an open-water surface. It is still unknown if this assumption is practically valid, because the Penman equation with the Rome wind function may result in unrealistically high  $ET_p$  even on a large wet area (McMahon et al., 2013). Given the importance of the aerodynamic resistance in modulating surface temperature (Chen et al., 2020), ignoring its variability may become a significant error source for the CR method at both annual and sub-annual timescales.

In addition, there are some caveats in our case study. We employed the meteorological data different from those used in Ma et al. (2021). The ERA5-Land dataset is a downscaled version of the ERA5 data (Hersbach et al., 2020) by which Ma et al. (2021) predicted  $ET_a$  globally. Ma et al. (2021) incorporated remotely sensed albedo and emissivity together with a correction factor when calculating  $R_n$ , whereas we used the sum of the ERA5-Land latent and sensible heat fluxes. Those input differences may lead to differences in CR  $ET_a$  estimates.

The gridded GRUN  $Q_{too}$  has some uncertainty sources, though it is the ensemble of many runoff simulations from 21 different atmospheric forcing inputs. In the machine-learning process by Ghiggi et al. (2021), some  $Q$  observations affected by human activities (e.g., dam regulation and return flows from groundwater abstraction) might not be excluded, potentially disrupting the empirical relationship between atmospheric forcing and natural flows. In addition, the uncertainty of SILO  $P$  might be non-negligible in areas with limited weather stations and in mountainous areas (Fu et al., 2022). Though we reduced the potential errors in the gridded  $P$  and  $Q$  datasets by temporal averaging, the grid-scale  $ET_{wb}$  estimates should be treated as plausible values rather than exact observations.

#### 5 Conclusions

Via a case study in Australia, we showed that the polynomial CR by Szilagyi et al. (2017) is unlikely to perform well when local climate conditions are neglected. The assumption of a constant Priestley-Taylor coefficient cannot reflect the long-term water balance; thereby, produced biased CR  $ET_a$  estimates. We resolved this problem by conditioning the CR with the

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In seven Australian eddy-covariance flux towers, Crago et al. (2022) found that the optimal  $\alpha_c$  for the CR of Szilagyi et al. (2017) was 1.35 when predicting daily  $ET_a$  in the dimensionless form (i.e.,  $y = ET_p/ET_p$ ). However, it should be increased to 1.42, 1.45, 1.47, and 1.50 to simulate dimensional latent heat fluxes at daily, weekly, monthly, and annual timescales, respectively. In Crago and Qualls (2018), the optimal  $\alpha_c$  for the kindred linear CR of Crago et al. (2022) was 1.35 [190]

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traditional Budyko equation, and it allowed a reasonable determination of the essential parameter in ungauged locations. The following conclusions are worth emphasizing:

- (1) The constant Priestley-Taylor coefficient transferred from fractional wet locations could make the CR method perform poorly in closing local water balance. The too simple approach could make the CR method underperform the widely used physical, machine-learning, and land surface models.
- (2) The Budyko framework could become an additional constraint to determine the degree of  $ET_p$  adjustment at the mean annual scale. It allows upscaling of the Priestley-Taylor coefficients from gauged to ungauged locations.
- (3) The Priestley-Taylor coefficients conditioned by local climates made the CR method better close the basin-scale water balance. The varying Priestley-Taylor coefficients seemed to make the CR method outperform the advanced  $ET_a$  models.

#### Author contributions

DK, MC, and JAC organized this study together. DK built the research framework, simulated  $ET_a$  with the CR method, and drafted the manuscript. JAC processed the modeled  $ET_a$  datasets and reviewed the draft, and MC actively participated in discussing the results.

#### Competing interests

The authors declare no competing interests.

#### Code availability

The Python scripts that implement the CR method are available upon request from the leading author (daeha.kim@jbnu.ac.kr).

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**Table 1. List of the chosen FLUXNET2015 sites**

Site ID	Lon. (°E)	Lat. (°S)	Data period	Site ID	Lon. (°E)	Lat. (°S)	Data period
AU-ASM	133.25	22.28	2010-2014	<del>AU-Rig</del>	<del>145.58</del>	<del>36.65</del>	<del>2011-2014</del>
AU-Cpr	140.59	34.00	2010-2014	<del>AU-Stp</del>	<del>133.35</del>	<del>17.15</del>	<del>2008-2014</del>
AU-DaP	131.32	14.06	2007-2013	<del>AU-TTE</del>	<del>133.64</del>	<del>22.29</del>	<del>2012-2014</del>
AU-DaS	131.39	14.16	2008-2014	<del>AU-Tum</del>	<del>148.15</del>	<del>35.66</del>	<del>2001-2014</del>
AU-Dry	132.37	15.26	2008-2014	<del>AU-Wac</del>	<del>145.19</del>	<del>37.43</del>	<del>2005-2008</del>
AU-Emr	148.47	23.86	2011-2013	<del>AU-Whr</del>	<del>145.03</del>	<del>36.67</del>	<del>2011-2014</del>
<del>AU-Gin</del>	<del>115.71</del>	<del>31.38</del>	<del>2011-2014</del>	<del>AU-Wom</del>	<del>144.09</del>	<del>37.42</del>	<del>2010-2014</del>
<del>AU-How</del>	<del>131.15</del>	<del>12.49</del>	<del>2001-2014</del>				

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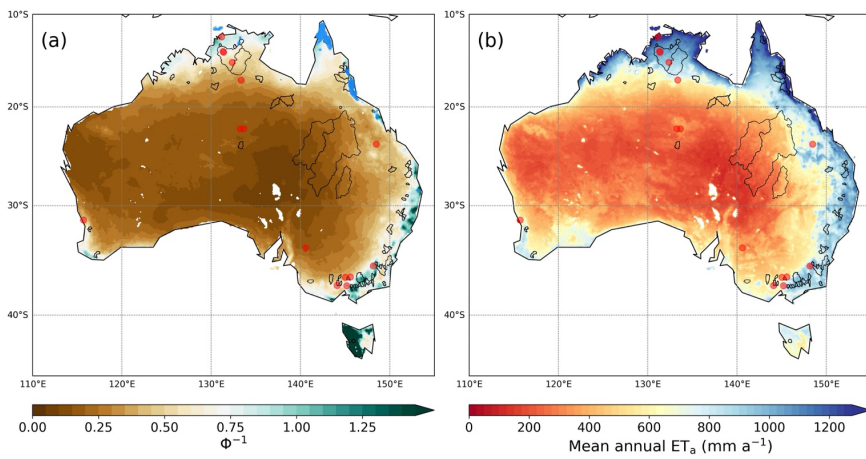
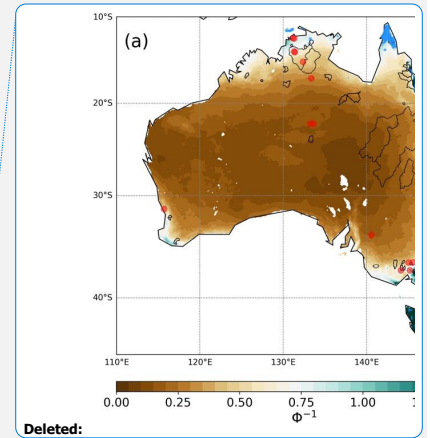


Figure 1: Spatial distributions of (a) the reciprocals of aridity index and (b) the mean annual  $ET_a$  for 1998-2014 predicted by the CR with  $\alpha_c = 1.15$ . The red circles and the gray polygons are the locations of 15 flux towers and the boundaries of 71 CAMELS basins. The blue-colored points in (a) indicate the wet cells with  $RH > 90\%$  and  $T_{ws} > T_a + 2^\circ C$ . CR  $ET_a$  was calculated at the grid cells where the land fraction was larger than 50%.



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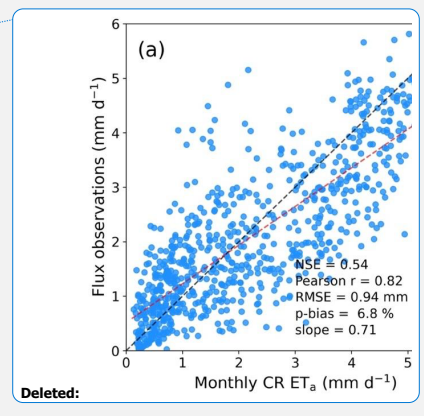
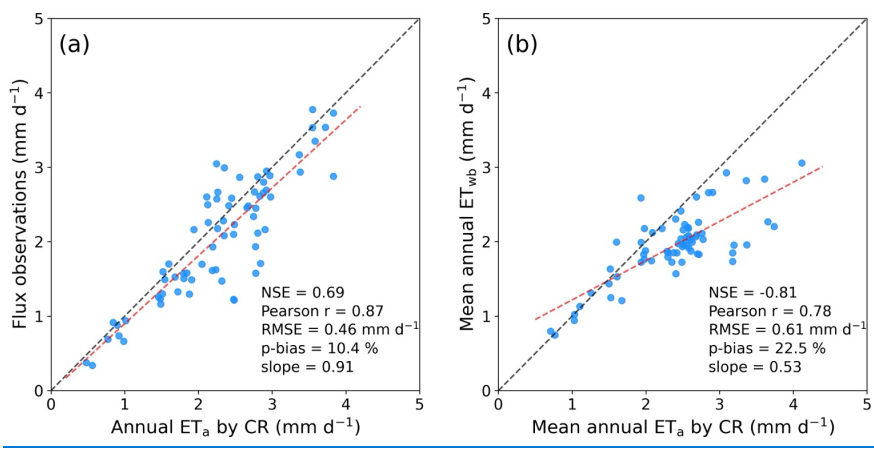
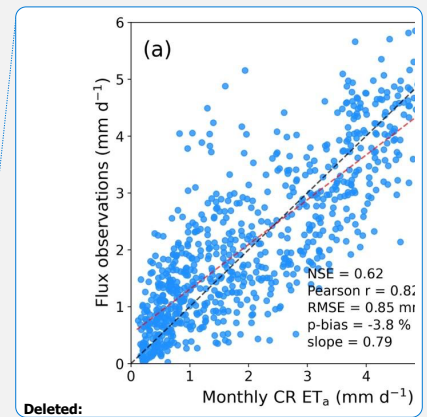
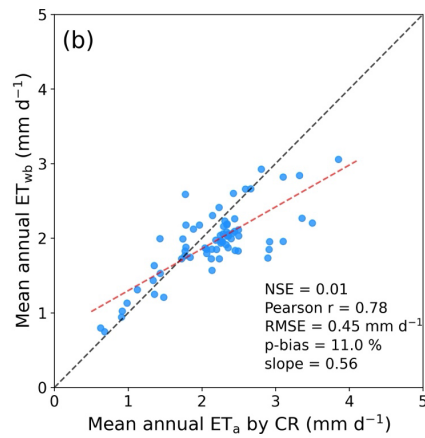
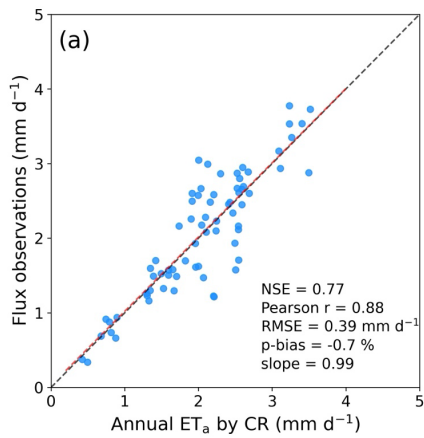


Figure 2: The 1:1 comparison between the CR ET<sub>a</sub> estimates with  $\alpha_c = 1.15$  and (a) the annual FLUXNET2015 observations and (b) the mean annual ET<sub>wb</sub> of the 71 CAMELS basins for 1998-2014.

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Figure 3: Same as Figure 2 except  $\alpha_c = 1.10$ .

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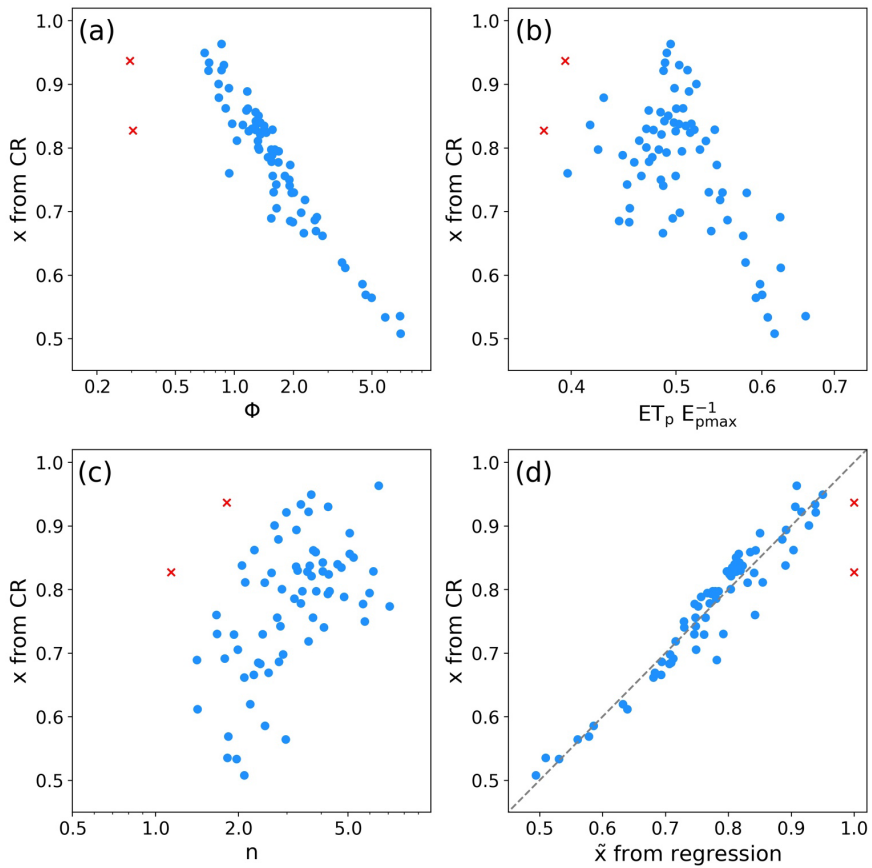


Figure 4: The scatter plots between the  $x$  estimated by CR with  $ET_{wb}$  for 1981-1997 and the corresponding (a)  $\Phi$ , (b)  $ET_p/E_{pmax}$ , and (c)  $n$  values, and (d) the 1:1 plot between the  $x$  from CR and the  $\tilde{x}$  predicted by Eq. (15). The red  $x$  symbols are the outliers excluded from the regression analysis.

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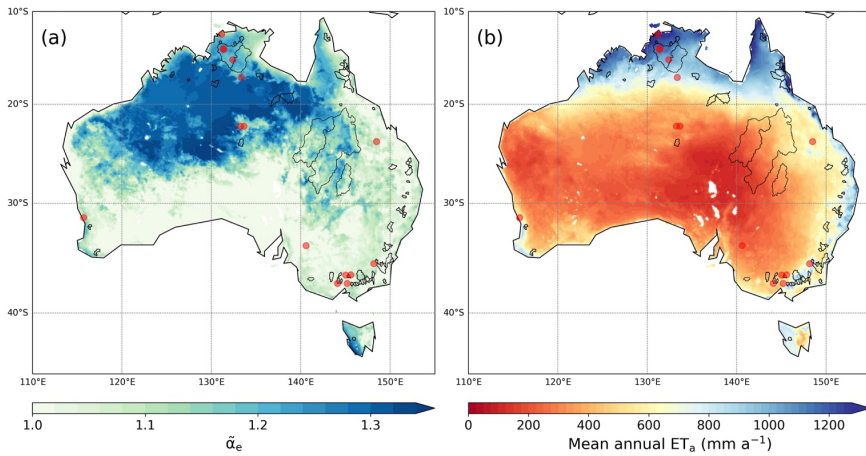
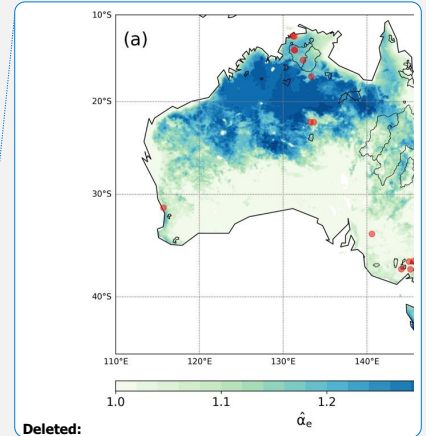


Figure 5: Distributions of (a) the  $\alpha_e$  values from Eq. (15), and (b) the mean annual  $ET_a$  for 1998-2014 by the CR method and the  $\alpha_e$  values.



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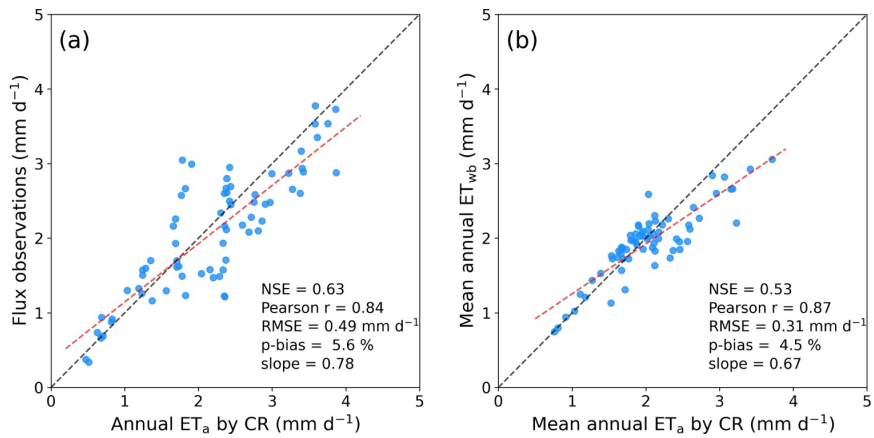
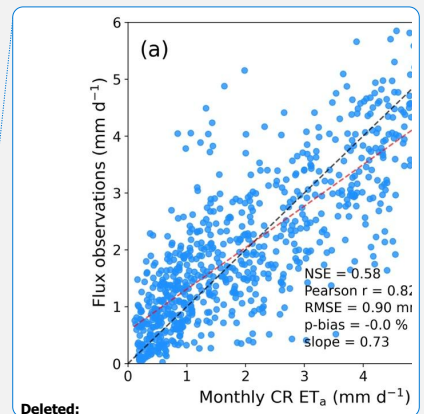


Figure 6: Same as Figure 2, except that the  $\sigma_e$  values from Eq. (15) were used for CR ET<sub>a</sub>.



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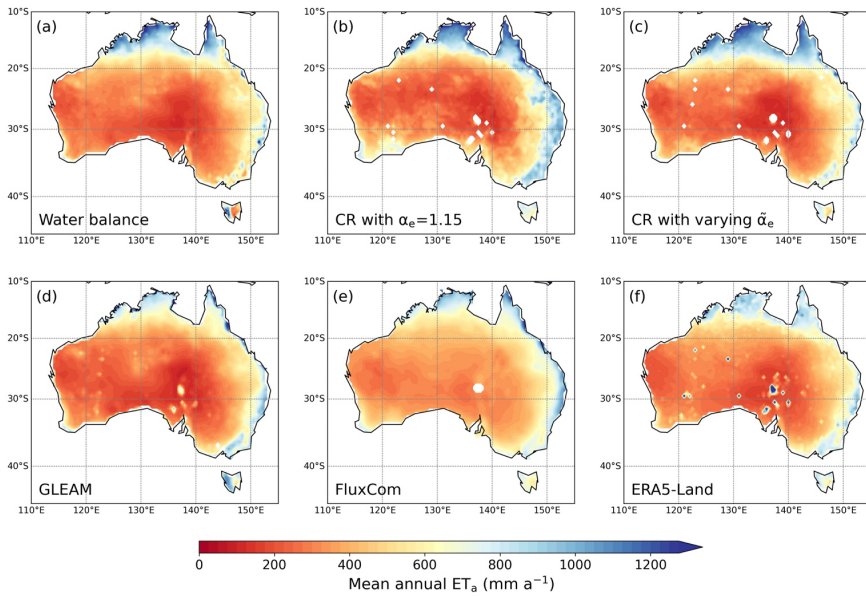


Figure 7: Distributions of (a) the mean annual water-balance  $ET_{wb}$  for 1998-2016, and the predictions by (b) CR with  $\alpha_e = 1.15$ , (c) CR with spatially varying  $\alpha_e$ , (d) GLEAM, (e) FluxCom, and (f) ERA5-Land.

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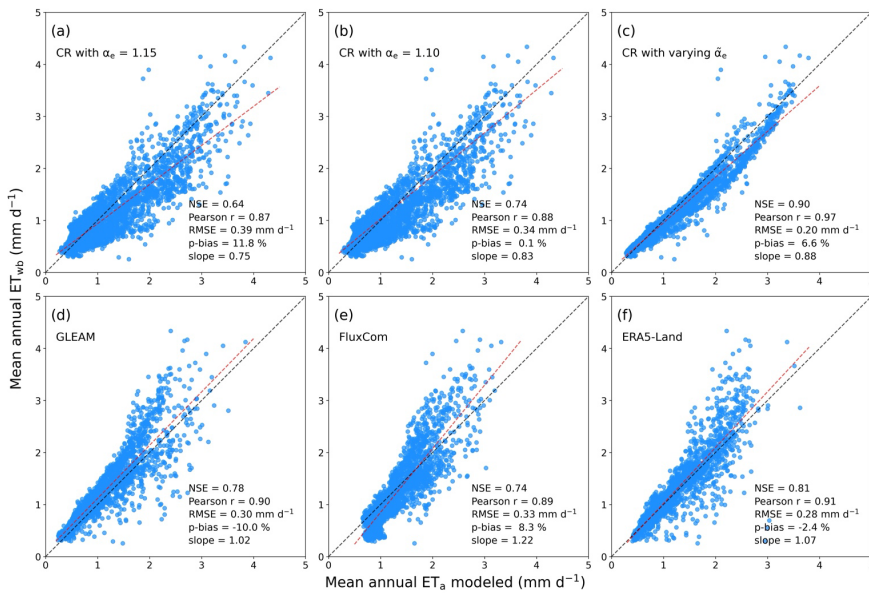
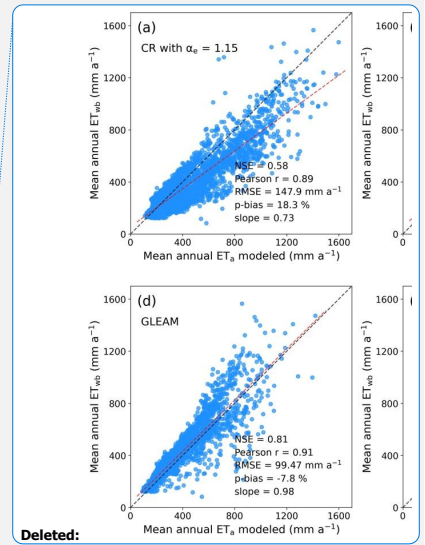


Figure 8: Scatter plots between the mean annual  $ET_{wb}$  for 1998-2016 at  $0.5^\circ \times 0.5^\circ$  and the predictions by (a) CR with  $\alpha_e = 1.15$ , (b) CR with  $\alpha_e = 1.10$ , (c) CR with spatially varying  $\alpha_e$ , (d) GLEAM, (e) FluxCom, and (f) ERA5-Land.



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