

1 S1 Minimalist Analytical Solution

2 The analytical solution for the minimalist PHM is derived by equating supply (T_s^{phm} ; Eqn. 1 of the article) and demand (T_s^{phm} ;
3 Eqns. 2-3 of the article) and solving for ψ_l^* as shown in Equation S1.

$$\begin{aligned}
 g_{sp} \cdot (\psi_s - \psi_l) &= \frac{\psi_{l,c} - \psi_l}{\psi_{l,c} - \psi_{l,o}} \cdot T_{ww} \\
 \psi_l^* &= \frac{\psi_{l,c} - \psi_{l,o}}{T_{ww} - g_{sp} \cdot (\psi_{l,c} - \psi_{l,o})} \cdot \left(\frac{T_{ww} \cdot \psi_{l,c}}{\psi_{l,c} - \psi_{l,o}} - g_{sp} \cdot \psi_s \right) \\
 \psi_l^* &= \frac{T_{ww} \cdot \psi_{l,c} - g_{sp} \cdot \psi_s \cdot (\psi_{l,c} - \psi_{l,o})}{T_{ww} - g_{sp} \cdot (\psi_{l,c} - \psi_{l,o})} \\
 \psi_l^* &= \frac{\frac{T_{ww} \cdot \psi_{l,c}}{g_{sp}} + \psi_s \cdot (\psi_{l,o} - \psi_{l,c})}{\frac{T_{ww}}{g_{sp}} + (\psi_{l,o} - \psi_{l,c})}
 \end{aligned} \tag{S1}$$

4 Substituting ψ_l^* back into Equation 1 of the article yields the analytical solution for the minimalist PHM (Eqn. S2 and Eqn.
5 4 in the article). Algebraic manipulations shows that the solution is simply T_d^{phm} with an additional dependence on the ratio of
6 atmospheric moisture demand and soil-plant conductance in the denominator.

$$\begin{aligned}
 T^{phm} &= g_{sp} \cdot (\psi_s - \psi_l^*) \\
 &= g_{sp} \cdot \left(\psi_s - \frac{\frac{T_{ww} \cdot \psi_{l,c}}{g_{sp}} + \psi_s \cdot (\psi_{l,o} - \psi_{l,c})}{\frac{T_{ww}}{g_{sp}} + (\psi_{l,o} - \psi_{l,c})} \right) \\
 &= g_{sp} \cdot \left(\frac{\psi_s \cdot \frac{T_{ww}}{g_{sp}} + \psi_s \cdot (\psi_{l,o} - \psi_{l,c})}{\frac{T_{ww}}{g_{sp}} + (\psi_{l,o} - \psi_{l,c})} - \frac{\frac{T_{ww} \cdot \psi_{l,c}}{g_{sp}} + \psi_s \cdot (\psi_{l,o} - \psi_{l,c})}{\frac{T_{ww}}{g_{sp}} + (\psi_{l,o} - \psi_{l,c})} \right) \\
 &= T_{ww} \cdot \frac{(\psi_{l,c} - \psi_s)}{(\psi_{l,c} - \psi_{l,o}) - \frac{T_{ww}}{g_{sp}}}
 \end{aligned} \tag{S2}$$

7 A key conclusion of this work relates to the nonlinearity in the PHM with respect to T_{ww} , even in the simplest case of the mini-
8 malist model. This nonlinearity can be shown formally by violating the superposition principle $T^{phm}(\psi_s, c_1 \cdot T_{ww,1} + c_2 \cdot T_{ww,2}) \neq$
9 $c_1 \cdot T^{phm}(\psi_s, T_{ww,1}) + c_2 \cdot T^{phm}(\psi_s, T_{ww,2})$. This is the fundamental difference between β and PHMs and results in the T_{ww}/g_{sp}
10 term in the denominator of Equation S2.

11 **S2 Additional LSM Results**

12 **S2.1 Soil Water Availability and Atmospheric Moisture Demand**

13 The improved performance of PHMs during midday of July-August (Figs. 4c-d in the article) are explained by looking at the
14 temporal breakdown of the well-watered transpiration (T_{ww}) and site data of soil water availability (Fig. S1). The T_{ww} is a proxy
15 for stomata-regulated atmospheric moisture demand at the site and is the greatest from 10 AM to 3 PM during the later summer
16 months. The measured volumetric water content shows water stress during the later summer months as well. Therefore, these
17 diurnal results suggest that PHMs are most important during periods of high atmospheric moisture demand and low soil water
18 availability.

19 **S2.2 Fitting β Schemes**

20 The three β transpiration downregulation schemes used in this work were ‘calibrated’ by fitting their respective parameters
21 to the outputs of the calibrated LSM that uses a PHM scheme (the calibration process is detailed more extensively in section
22 S5). The calibrated LSM outputs are relative transpiration, T/T_{ww} , for the sunlit and shaded big leaf (dots in Fig. S2e-f). We
23 decided to avoid calibrating each LSM directly with a β scheme to the site data, and instead derive the β scheme from a fitted
24 PHM scheme, because we wanted to ensure that any improvements resulting from PHM can be directly related to the ability of
25 PHM to capture downregulation more realistically.

26 The single β scheme (β_s) has a Weibull curve (Eqn. 16 in the article) fit to the combined calibrated sunlit and shaded
27 T/T_{ww} using nonlinear least squares in MATLAB. The fitted β_s parameter values for $\psi_{s,50}$ and b_s are -6.95 MPa and 2.54,
28 respectively (shown in Fig. S2a-d in light gray). The two-leaf β scheme (β_{2L}) fits a β curve to the calibrated sunlit and shaded
29 T/T_{ww} separately. The fitted β_{2L} parameter values for $\psi_{l,50}$ and b_l are -6.08 (-7.62) MPa and 2.12 (3.5) for the sunlit (shaded)
30 big leaf (shown in Fig. S2a-d in dark gray). The reader is referred to section S6 for details of how these β curves are used in
31 LSM calculations.

32 The ‘dynamic β ’ scheme (β_{dyn}) was fit to the calibrated T/T_{ww} using a two-step process. First, the T/T_{ww} values were
33 parsed into 10 bins covering the T_{ww} range for the sunlit and shaded big leaf separately and a single β curve was fit to each
34 bin (shown by the black circles in Figure S2a-d). Second, a line was fit to the parameters $\psi_{s,50}$ and b_s as a function of T_{ww}
35 shown by the red line and the corresponding linear equation in Figure S2a-d. Therefore, the parameters of β can dynamically
36 change with the atmospheric moisture demand represented by T_{ww} . This is illustrated in Figure S2e-f by the isolines of β_{dyn}
37 with respect to T_{ww} and closely match the color gradient of the calibrated T/T_{ww} values. The variation of β_{dyn} with respect to
38 T_{ww} is well described by linear functions, with the exception of a slight noise in the shaded leaf b_s value, which is likely due to
39 the clustering of T/T_{ww} values around 1 in wetter conditions in Figure S2f.

40 The β_{dyn} has great potential for parsimoniously representing the complexity of a PHM. The slope and intercepts of the β_{dyn}
41 linear parameters for sunlit and shaded leaves are very similar making separate leaf fits unnecessary. Therefore, the complexity
42 of the PHM can be represented by a ‘dynamic β ’ with 4 total parameters (2 slope and 2 intercept), which is two more than the

43 original β model as mentioned in the paper. A promising avenue of future work is to relate these four parameters to key plant
44 hydraulic traits and soil parameters.

45 **S2.3 RMSE Comparison of PHM and β Schemes**

46 The improvements of the PHM scheme to the β_s and ‘dynamic β ’ schemes are shown in terms of reduction in percent bias in
47 Figure 4e-f. These results are corroborated by the change in root mean square error as shown in Figure S3. The RMSE results
48 only differ from those based on reduction in percent bias in terms of improvements that are concentrated toward the highest T_{ww}
49 periods, since that is where the highest magnitude errors occur.

50 **S2.4 LSM Cumulative Energy and Carbon Budget Errors**

51 To aid the interpretation of the LSM case study, we have also calculated the cumulative error compared to key measured fluxes
52 at the US-Me2 site for the LSM run with five separate transpiration downregulation schemes (Table S1). Clearly, the PHM
53 and β_{dyn} schemes provide the greatest improvement to evapotranspiration (ET) and gross primary productivity (GPP) with
54 a 9% and 5% reduction in cumulative error, respectively, while differences in sensible heat flux (H), net radiation (R_n) and
55 outgoing longwave radiative flux (L_{out}) appear less significant. Although they can be outweighed by energy balance closure
56 errors in the flux tower data (up to 20%¹), these improvements in percent bias (Fig. 4e-f) and root mean square error (Fig. S3)
57 are consistent with our theoretical analysis of fundamental differences between β and a PHM under varying environmental
58 conditions. Therefore, these errors may persist and grow under longer simulations and more variable environmental conditions.

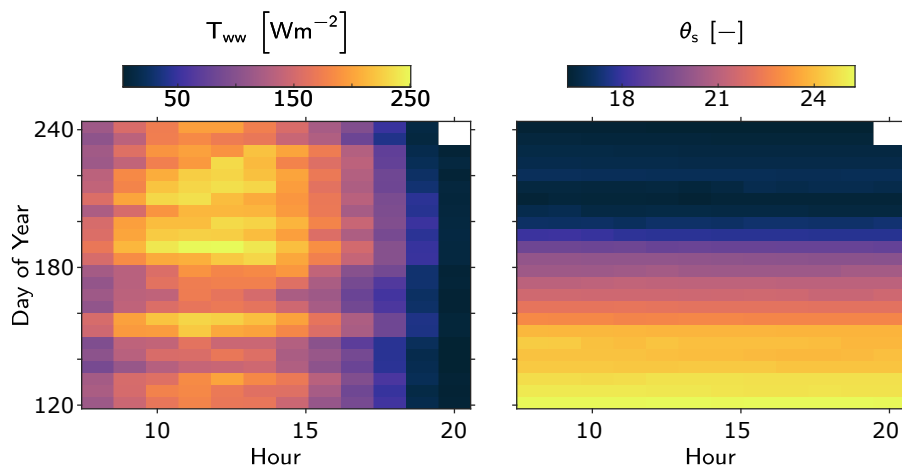


Figure S1. Left: Well-watered transpiration rate calculated from the LSM run with no transpiration downregulation. This is a proxy for the stomata-regulated atmospheric moisture demand. Right: Measured volumetric water content of soil at the US-Me2 site at 50 cm depth. The colors are the average value for the temporal bins for May-August 2013-2014.

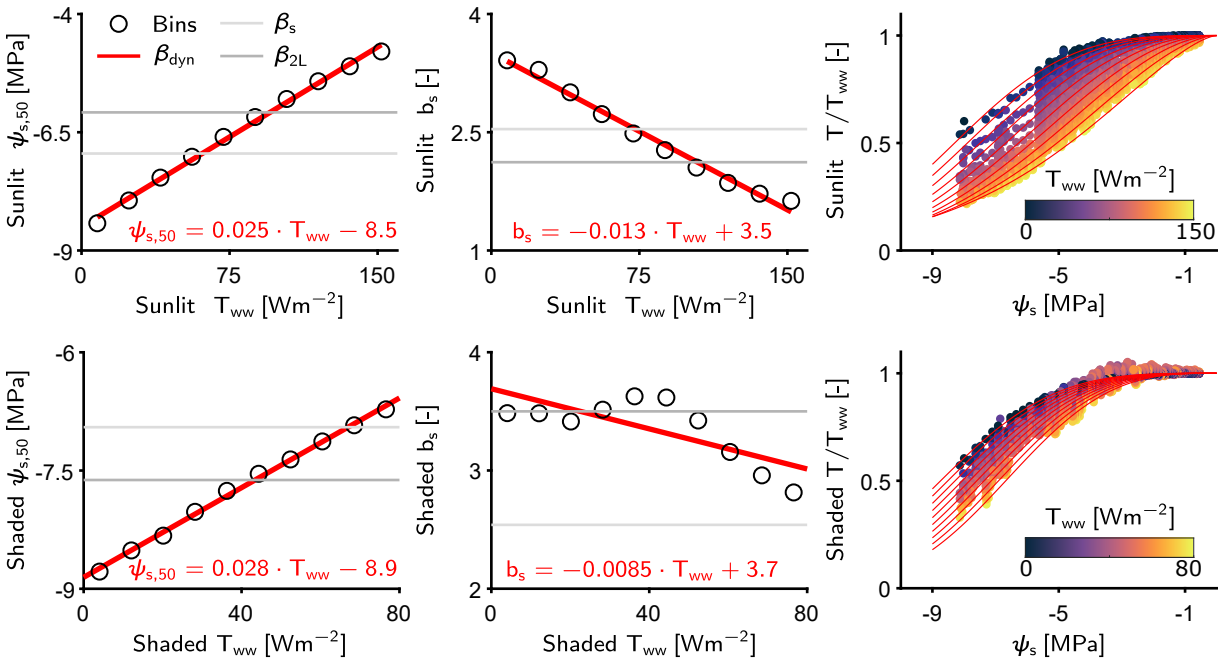


Figure S2. The ‘dynamic β ’ (β_{dyn}) fits used for the sunlit (top row) and shaded big leaf (bottom row). The first column is the dependence of the soil water potential at 50% loss of stomatal conductance on well-watered transpiration T_{ww} . The second column is the dependence of the stomatal sensitivity parameter (b_s) to T_{ww} . The black circles are parameter values fit to relative transpiration (T/T_{ww}) binned over the range of T_{ww} . The linear relationship for both parameters is shown in red. The last column shows the relative transpiration outputs from the calibrated PHM with dot colors corresponding to T_{ww} . The red lines are the β_{dyn} model isolines at 10 values of T_{ww} (Equation 16 of the main article). These isolines clearly follow the color gradient of the PHM results indicating that β_{dyn} is able to capture the complexity of a PHM.

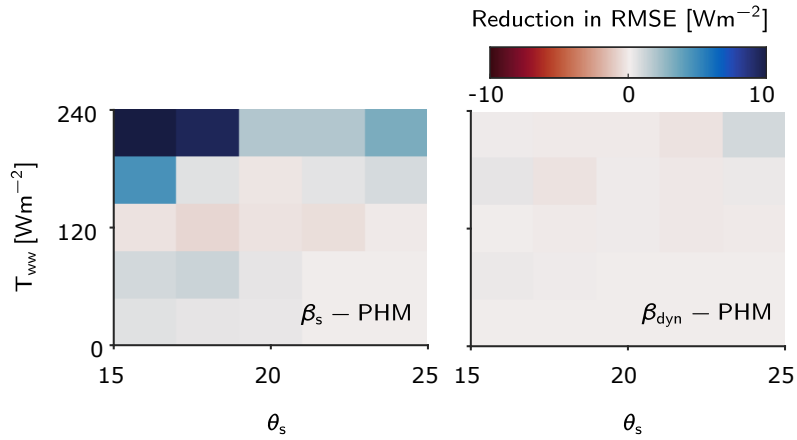


Figure S3. Analogous results to Figure 5e-f in the main text using root mean square error instead of percent bias as the performance metric. The differences in reduction of RMSE between the PHM and β_s scheme (left) and β_{dyn} scheme (right).

Table S1. Cumulative total for evapotranspiration (ET), gross primary productivity (GPP), sensible heat flux (H), net radiative flux (R_n), and longwave radiative flux (L_{out}) for Ameriflux Me2 data and 5 LSM simulations during high evaporative demand ($ET_{ww} > 150 W m^{-2}$) and low soil moisture (Volumetric soil water content < 0.2) for May-Aug 2013-2014. The surface energy fluxes are in units of cm H₂O and GPP is in units kg CO₂. The values in parentheses are the percent error compared to the observations.

	ET [cm]	GPP [kg]	H [cm]	R_n [cm]	L_{out} [cm]
WW	18.2 (34.4%)	1.7 (30.8%)	23.8 (-9.4%)	43.5 (-14.8%)	43.6 (10.4%)
PHM	13.6 (0.9%)	1.2 (-2.5%)	27.5 (4.7%)	42.6 (-16.4%)	44.5 (12.5%)
Beta	15 (10.8%)	1.4 (7%)	26.4 (0.6%)	42.9 (-15.9%)	44.2 (11.9%)
Beta2L	14.4 (6.9%)	1.3 (3.1%)	26.8 (2.2%)	42.8 (-16.1%)	44.3 (12.2%)
BetaDyn	13.6 (0.7%)	1.2 (-2.7%)	27.5 (4.8%)	42.6 (-16.4%)	44.5 (12.6%)

59 S3 Defining a Threshold for Transport-limitation

60 Quantifying the values of particular soil parameters and plant hydraulic traits that define a soil-plant system as transport-limited
61 is an important avenue of future work. Figure 2 in the article illustrates clearly that, even in the minimalist model, there is
62 a complex interplay of drivers that contribute to the differences between PHM and β and, in turn, if a system is transport-
63 or supply-limited. However, the overall soil-plant conductance in the minimalist model seems to be the main control on
64 transport-limitation and a $g_{sp} \approx 10^3 \text{ Wm}^{-2} \text{ MPa}^{-1}$ appears to yield a supply-limited system (Fig. S4). The definition of
65 transport-limitation is somewhat subjective as it depends on how much difference between PHM and β is considered acceptable.

66 Determining a threshold of transport-limitation for the complex PHM is even less clear given the additional parameters.
67 Therefore, a sensitivity analysis was performed using the recent Variogram Analysis of the Response Surface (VARS) method²
68 implemented with the VARSTOOL package in MATLAB³. The integrated difference in β and PHM-generated relative
69 transpiration at high T_{ww} (150 Wm^{-2}) normalized by the relative soil saturation of soil water stress (M_{dif} , Eqn. S3) was used
70 as our metric to quantify the performance of each parameter set. The ranges and sensitivity scores for the 8 selected PHM
71 parameters are shown in Table S2. The VARSTOOL analysis reveals that the maximum xylem-to-leaf conductance ($g_{xl,max}$) is
72 the most sensitive parameter; thus, as maximum conductance in the plant decreases, a single β curve becomes increasingly
73 ineffective at downregulating transpiration realistically. The next most sensitive parameters are $\psi_{x,50}$, b , $\psi_{l,50}$, b_l , and a , but
74 they are of secondary importance. Lastly, the remaining two soil parameters, $\psi_{s,sat}$ and $g_{sx,max}$, were found to be the least
75 sensitive parameters because transport-limitation from soil is primarily controlled by b .

$$M_{dif} = \frac{1}{T_{ww} \cdot (\theta_o - \theta_c)} \cdot \int_{\psi_s} T^\beta(\psi_s) - T^{phm}(\psi_s, T_{ww}) d\psi_s \quad (\text{S3})$$

76 Focusing on $g_{xl,max}$, we estimate a threshold for transport-limitation similar to the minimalist model. We do so by parsing
77 the $g_{xl,max}$ range into 14 bins and sampling 5000 parameter sets from each bin (the 7 other parameters are sampled from their
78 entire range in Table S2 for this analysis). The resulting sensitivity metrics were plotted for each bin in Figure S4. As $g_{xl,max}$
79 becomes lower ($g_{xl,max} < 10^3 \text{ Wm}^{-2} \text{ MPa}^{-1}$) there is a tendency for the PHM results to diverge substantially from those of a
80 single β curve. This threshold notably coincides with that predicted by the minimalist model. The large amount of spread is
81 likely caused by the interactions amongst the other parameters. Further work must be done to create a more robust relationship
82 based on measurable plant and soil hydraulic parameters.

Table S2. VARSTOOL results for plant hydraulics model based on 35,600 parameter sets created using Progressive Latin Hypercube Sampling and 200 STAR sampling centers. The IVARS50 is an integrated metric of sensitivity that accounts for correlation of nearby parameter values in the parameter space. The sources for each parameter are how we determined a realistic range to sample from.

Parameter	Description	Range	Units	IVARS50	Sources
$g_{xl,max}$	Max xylem-to-leaf conductance	$[10^{-10}, 10^{-3}]$	$\frac{m}{sMPa}$	$1.6e^{-3}$	4,5
$\psi_{x,50}$	Xylem water potential at 50% loss of conductance	$[-0.1, -15]$	MPa	$8.0e^{-4}$	6-8
b	Soil retention curve exponent	$[2, 14]$	-	$1.5e^{-5}$	9
$\psi_{l,50}$	Leaf water potential at 50% loss of conductance	$[-0.1, -15]$	MPa	$3.5e^{-4}$	4
a	Xylem vulnerability curve shape parameter	$[0.2, 10]$	-	$2.4e^{-4}$	10
b_l	Leaf vulnerability curve shape parameter	$[0.2, 5]$	-	$1.0e^{-4}$	10
$g_{sx,max}$	Max soil-to-xylem conductance	$[10^{-2}, 10^3]$	$\frac{m}{sMPa}$	$3.0e^{-5}$	4,11
ψ_{sat}	Saturated soil water potential	$[10^{-3}, 10^{-2}]$	MPa	$3.2e^{-6}$	9

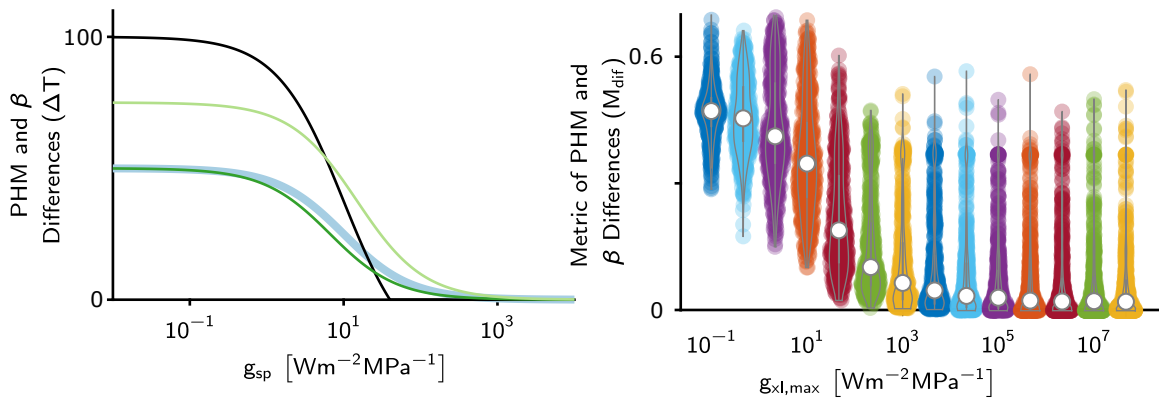


Figure S4. The control of soil-plant conductance (g_{sp}) on transport-limitation of a soil-plant system. Left: Differences in minimalist PHM and β as a function of overall soil-plant conductance. The thick blue line represents change in g_{sp} with three other drivers at baseline values (see Figure 2 in main article) while the thin lines represent 50% increase in ψ_s (black), T_{ww} (light green) and $\psi_{l,c} - \psi_{l,o}$ (dark green) compared to their baseline values. Right: Differences between a more complex formulation of PHM and β used in the LSM analysis with respect to maximum xylem-to-leaf conductance. The metric used integrates the difference between relative transpiration of β and PHM at a $T_{ww} = 150 W m^{-2}$ normalized by the range of soil water availability over which downregulation occurs (Eqn. S3).

83 S4 LSM Forcing Data

84 The LSM for the US-Me2 ponderosa pine site was forced with half-hourly atmospheric and subsurface measurements at the
85 site. This site was specifically selected for the LSM case study based on its extensive subsurface soil moisture and temperature
86 profiles as well as its separate measurements of photosynthetically active radiation (PAR) and near infrared radiation (NIR).
87 The extensive soil moisture and temperature data and detailed shortwave radiation measurements were used as forcing for the
88 LSM in lieu of one-dimensional mass and heat transfer equations and atmospheric radiation partitioning models. The main
89 focus of this work was on the scalar transport of the LSM, so use of these measurements help reduce confounding errors from
90 other model structures (although there would still be measurement errors).

91 The following atmospheric measurements from the Ameriflux US-Me2 dataset for May-August 2013-2014 were used to
92 force the LSM: friction velocity (u^*), mean streamwise velocity (\bar{u}), air temperature (T_a), water vapor pressure (e_a), atmospheric
93 pressure (P_{atm}), and CO₂ partial pressure (c_a). The radiative site measurements consisted of total incoming shortwave (S_{in})
94 and longwave radiation (L_{in}) as well as total and diffuse PAR. The LSM requires partitioning of shortwave radiation into PAR
95 and NIR as well as direct beam and diffuse quantities. The diffuse incoming PAR ($S_{in,par,d}$) was measured at the site and the
96 direct beam PAR ($S_{in,par,b}$) was calculated by the difference of total PAR ($S_{in,par}$) and diffuse PAR. Unfortunately, the site did
97 not differentiate between direct beam ($S_{in,nir,b}$) and diffuse NIR ($S_{in,nir,d}$); therefore, total NIR was partitioned using the same
98 ratio of direct and diffuse PAR at every time step. These detailed radiation measurements constrained the use of data from
99 2013-2014, as this was when they were most consistently available.

100 The subsurface moisture and temperature data was used to calculate the soil water availability of the root zone and the
101 ground heat flux G at each time step. The G used to force the model was simply the thermopile measurements at 5 cm. In
102 contrast, selecting a depth for the soil water content (θ_s) that would be representative of root-zone soil water availability
103 was more difficult given there is minimal information at the US-Me2 site about the root structure. The US-Me2 site has θ_s
104 measurements at 10, 20, 30, 50, 70, 100, 130, and 160 cm at multiple locations. To select a representative depth, we analyzed
105 GPP deviations from the mean in terms of θ_s and vapor pressure deficit (D) at each depth (Fig. S5). All GPP values were
106 studentized (i.e., mean subtracted and normalized by standard deviation) by hourly subsets for the period of May-August
107 2002-2014 to remove diurnal variation in flux magnitude and the median of these scores is plotted for each θ_s - D bin. The blue
108 (red) values indicate lower (higher) than average GPP fluxes. As expected, measurements at each depth show lower values
109 during water stress periods (low θ_s and high D). However, the ranges of θ_s experienced varies with depth, likely due to the
110 combined effects of variable soil moisture profile, soil texture heterogeneity, and sensor inaccuracy. We selected the depth of
111 50 cm to use as our soil moisture forcing for two reasons: 1) there is a clear signal of GPP downregulation covering a wider
112 range of soil moistures, and 2) a depth of 50 cm seems reasonable to represent the average moisture conditions when looking at
113 meta-analyses of temperate coniferous forest root measurements¹².

114 A crucial consequence of using the subsurface inputs as model forcings is that it allowed the model time steps to be run in
115 parallel. Typically, the model must be run sequentially since the subsurface models are partial differential equations in space

116 (soil column) and time, and each time step relies on previous energy stored in the subsurface. The observations codify this
117 temporal information, thereby allowing the LSM to run steady-state energy partitioning on top of the temporal dynamics of soil
118 moisture and heat. Additionally, the LSM simulation was run only for time steps 24 hours after precipitation, since the model
119 was not coded to handle canopy precipitation interception. Lastly, atmospheric stability effects were ignored for simplicity, as
120 they add an additional layer of complexity to the solution scheme¹³.

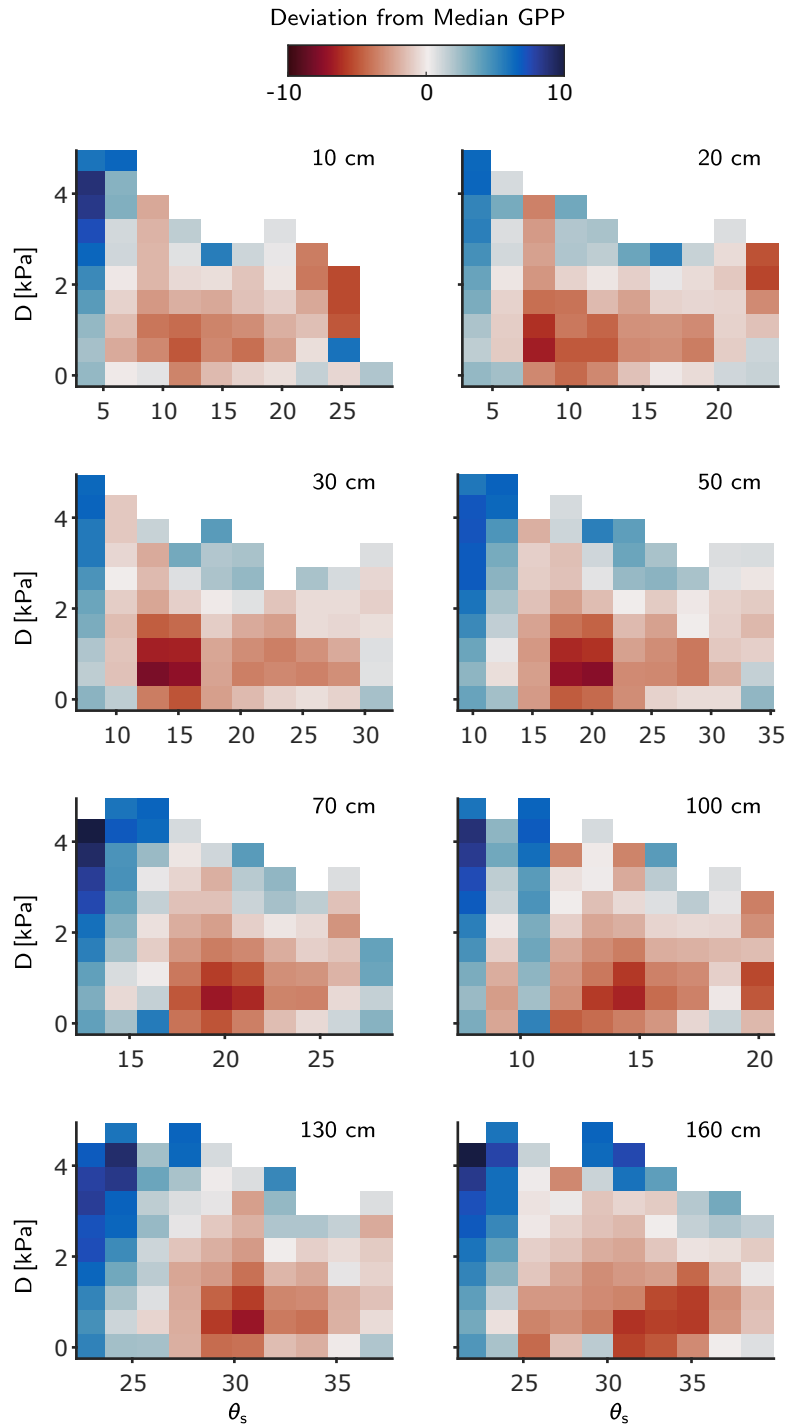


Figure S5. Median scores of the studentized gross primary productivity (GPP) measurements at the US-Me2 flux site for differing depth soil water content θ_s measurements. The θ_s and vapor pressure deficit (D) measurements help identify water stress periods. The GPP data used are from May-August 2002-2014 and are studentized by their hourly subset to remove diurnal variations. Blue (red) in the plots is an indicator of decreased (increased) GPP from the mean value.

S5 LSM Calibration

The LSM was calibrated using a grid search of 13,600 parameter sets with 15 soil, plant and radiative parameters shown in Table S3. The parameters sets were created using Progressive Latin Hypercube Sampling available in the VARSTOOL package³ in MATLAB. The 13,600 simulations were run for May-August 2013-2014 for the hours of 8 AM to 8 PM, excluding 24 hours following any precipitation event.

The best simulation was selected based on a performance metric evaluating evapotranspiration (ET), sensible heat flux (H), gross primary productivity (GPP) and net radiation (R_n) predictions. The performance metric (M_{cal} ; Eqn. S4) consists of Taylor diagram statistics¹⁴: 1) the correlation coefficient (R), 2) the centered root mean square error ($CRMSE$), and 3) the variance (σ). The percent bias (P_b) was also added to the metric to account for the mean difference in simulation and observation. Each index i in the summation of Equation S4 represents a different flux which are combined to form a single metric.

$$M_{cal} = \sum_i^n \frac{R_i}{\max(R)} - CRMSE_i - P_{b,i} - \frac{\Delta\sigma_i}{\max(\Delta\sigma)} \quad (S4)$$

The three best fits from the VARSTOOL grid search (VT1-VT3) were selected based on the metric. The best VARSTOOL parameter set (VT1) was further adjusted by reducing $g_{xl,max}$ by 60% to reduce biases found in representing ET . This is the calibrated parameter set used in the article and is labeled ‘Best Fit’ in the following figures. The metric value of all LSM runs are shown in Figure S6 in terms of R and $CRMSE$. The selected best parameter sets trade-off improvements between ET , H , GPP , and R_n . The outgoing longwave (L_{out}) and shortwave radiation (S_{out}) were ignored as including them in the metric had minimal effect. The ‘Best Fit’ (red x) provides clear improvement to the R and $CRMSE$ compared to VT1-VT3 for ET . The median diurnal fluxes for the observations and best fit LSM runs during May-June (Fig. S7) and July-August (Fig. S8) reveal the largest performance differences between parameter sets are for ET and GPP predictions. The over-prediction of ET during soil water stress (Fig. S8) informed our decision to create the ‘Best Fit’ parameter set by reducing $g_{xl,max}$ to correct the bias. This manual adjustment also provides slight performance increases to the second order statistics of all fluxes illustrated in the Taylor diagram¹⁴ in Figure S9.

The ‘Best Fit’ parameters set fit the ET observations well, but, as illustrated in Figure 4e of the article, the PHM downregulation scheme is not perfect as there are $T_{ww}-\theta_s$ bins where the β_s scheme performs the best. Looking at the P_b statistic for β_s , PHM and well-watered LSM runs in Figure S10, we see there are two primary reasons for β_s having the best performance for particular bins (highlighted in red): 1) the well-watered scheme is nearly unbiased so any downregulation will bias the result, and 2) the PHM over-regulates for these lower T_{ww} bins. For both situations, β_s downregulates less because it is fit to the mean behavior of the PHM, resulting in a better result. Therefore, the β_s is really just causing less bias in areas where the PHM performs poorly due to our fitting assumptions. The results in Figure S10 indicate that greater attention in the calibration process should be paid to the lower T_{ww} time steps to help correct these errors.

Table S3. The calibration parameters for the LSM with PHM downregulation scheme. The parameter ranges was used to create 13,600 parameters sets that were each run in the LSM. The calibrated value was selected based a performance metric (Eqn. S4) and additional adjustment in the above text. These values were used to run the LSM with PHM downregulation in the main article.

Parameter	Description	Range	Units	Calibrated Value
K_{sap}	Sapwood hydraulic conductivity	$[10^{-7}, 10^{-2}]$	$\text{kg m}^{-1} \text{s}^{-1} \text{MPa}^{-1}$	9.3e^{-4}
$\psi_{x,50}$	Xylem water potential at 50% loss of conductance	$[-0.1, -15]$	MPa	-2.3
a	Xylem vulnerability curve shape parameter	$[0.2, 10]$	-	0.3
$\psi_{l,50}$	Leaf water potential at 50% loss of conductance	$[-0.1, -15]$	MPa	-9.9
b_l	Leaf vulnerability curve shape parameter	$[0.2, 5]$	-	3.4
b	Soil retention curve exponent	$[2, 14]$	-	5.1
$\psi_{s,sat}$	Saturated soil water potential	$[10^{-3}, 10^{-2}]$	MPa	9.9e^{-3}
$K_{s,sat}$	Saturated soil hydraulic conductivity	$[0.01, 20]$	m d^{-1}	10
θ_i	Incipient soil water content for restricting bare soil evaporation	$[0, 0.57]$	-	0.57
g_1	Medlyn Slope Parameter	$[0.5, 5]$	$\text{kPa}^{0.5}$	0.9
$V_{max,25}$	Max Rubisco-limited carboxylation rate	$[5, 200]$	$\mu\text{mol CO}_2 \text{m}^{-2} \text{s}^{-1}$	122
LAI	Leaf area index	$[1.5, 4]$	$\text{m}^{-2} \text{LA m}^{-2} \text{GA}$	3.2
$\alpha_{l,par}$	Leaf reflectance to PAR	$[0.5, 1]$	-	0.74
$\alpha_{l,nir}$	Leaf reflectance to NIR	$[0, 0.6]$	-	0.43
χ_l	Leaf angle distribution parameter	$[-0.4, 0.6]$	-	0.11

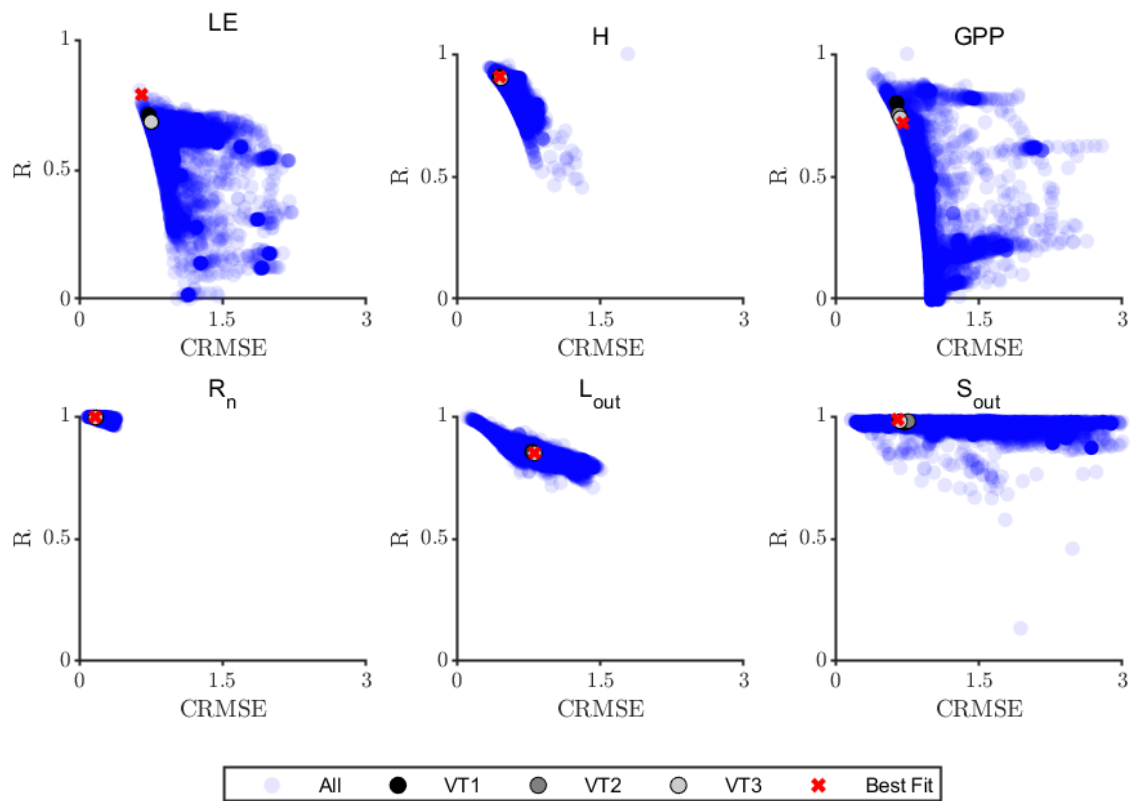


Figure S6. Centered root mean square error (CRMSE) and correlation (R) statistics for the LSM with PHM downregulation scheme for 13,600 parameters sets. An ideal score would be $R = 1$ and $CRMSE = 0$. The best fits from the VARSTOOL-created parameters sets, labelled VT1-VT3, are based on the metric in Equation S4, while a manual adjustment to VT1 was used to create the overall Best Fit used in the main article.

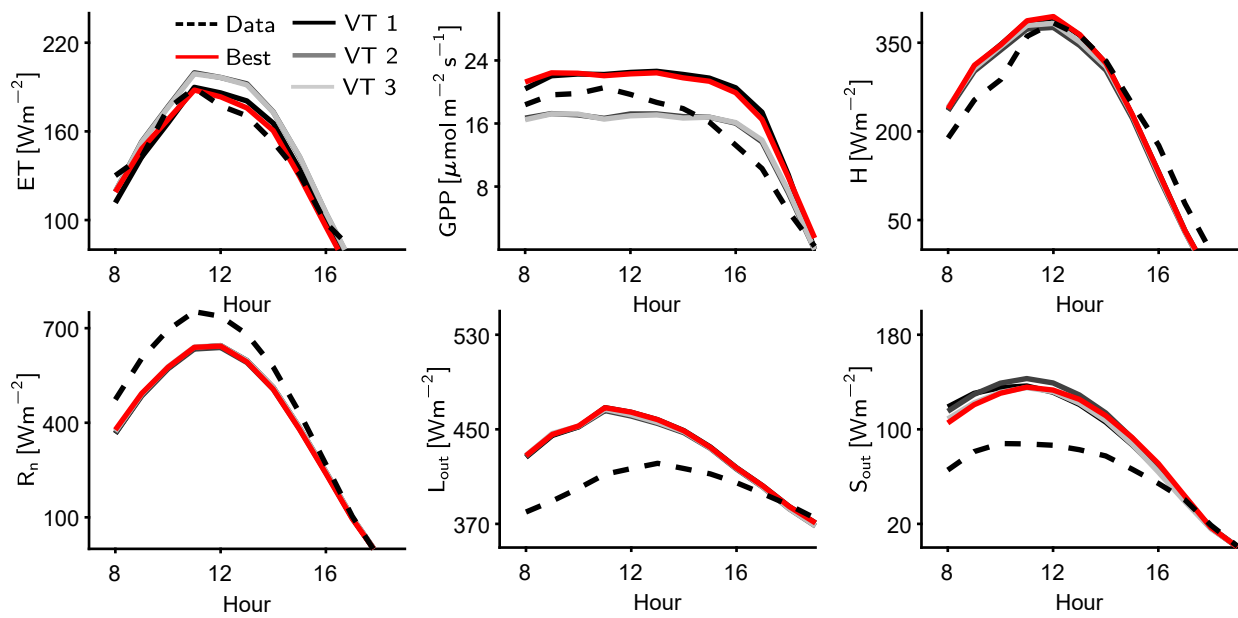


Figure S7. The median diurnal fluxes for May-June 2013-2014 for the three best VARSTOOL parameter sets (VT1-VT3) and the best overall calibrated fit (red) compared to the US-Me2 flux data for evapotranspiration (ET), gross primary productivity (GPP), sensible heat flux (H), net radiation (R_n), outgoing longwave radiation (L_{out}) and outgoing shortwave radiation (S_{out}).

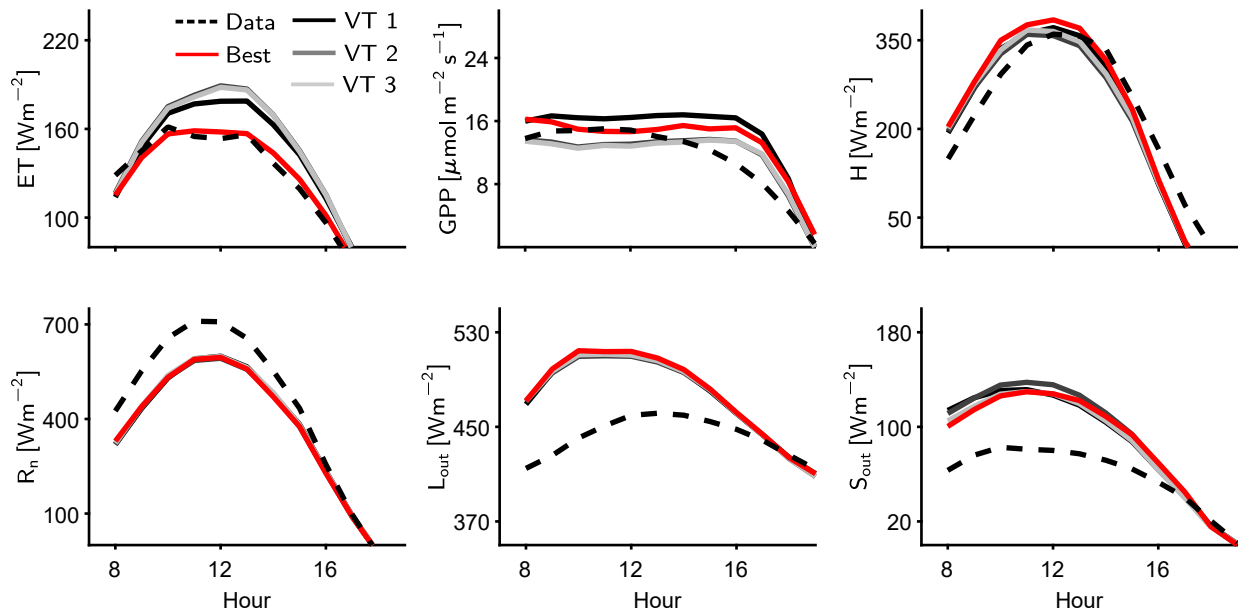


Figure S8. Same as Figure S7 for July-August 2013-2014 where there is soil water stress.

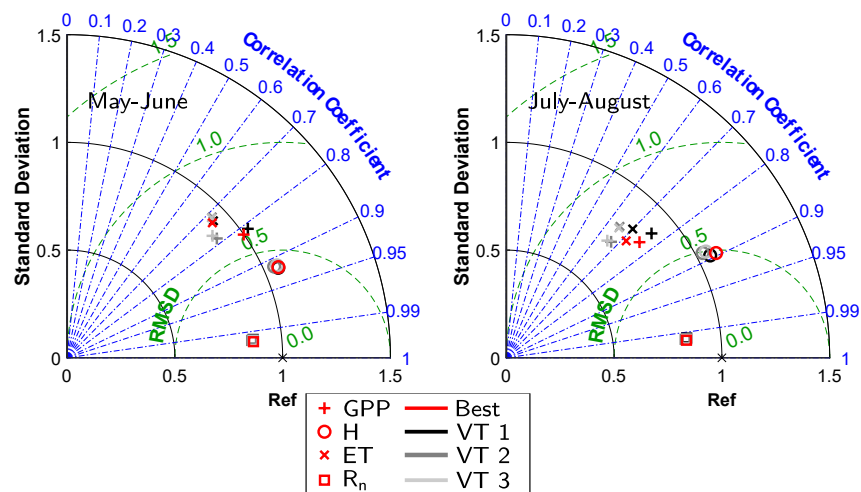


Figure S9. The Taylor diagrams for May-June 2013-2014 (left) and July-August 2013-2014 (right) for the three best VARSTOOL parameter sets (VT1-VT3) and the best overall calibrated fit (red) compared to the US-Me2 flux data for Evapotranspiration (ET), gross primary productivity (GPP), sensible heat flux (H), net radiation (R_n), outgoing longwave radiation (L_{out}) and outgoing shortwave radiation (S_{out}).

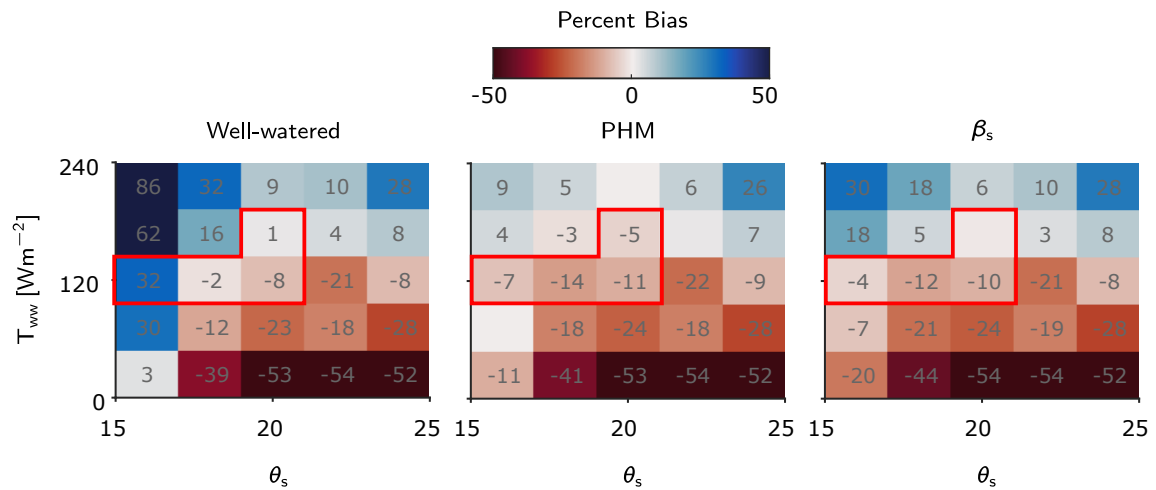


Figure S10. The percent bias (P_b) of the LSM with well-watered, PHM and β_s downregulation schemes compared to ET observations at the US-Me2 flux tower site. The P_b is broken down by well-watered transpiration T_{wvw} and volumetric soil water content (θ_s) as in Figure 4e-f of the main article. The gray numbers give the exact P_b value for each bin while the red outline highlights the primary bins where β_s appears to outperform the PHM in Figure 4e of the main article. See text for explanation.

150 S6 LSM Description

151 This section lays out the land surface model (LSM) coded in MATLAB used for the analysis of the US-Me2 ponderosa pine
152 Ameriflux site. The model is a two big leaf, dual source model¹⁵ following closely the formulation laid out in the Community
153 Land Model version 5¹⁶ with key modifications. The general model structure for scalar transport is shown in Figure S11 with
154 the main modules highlighted. Here, module refers to a smaller model within the overall LSM, e.g., the Plant Hydraulics Model
155 (PHM). The purpose of this LSM is to compare the scalar transport (temperature, water vapor, and carbon transport) scheme
156 using PHM and empirical (β) transpiration downregulation schemes; therefore, the model is simplified to be forced at the
157 boundaries by incoming radiation, air scalar concentrations as well as soil water availability and heat flux. We are exploring the
158 LSM component only during the growing season, so nutrient cycling, plant demographics, snow dynamics, and phenology
159 components—common in terrestrial biosphere models like CLM—are ignored. This section is organized by the energy balance,
160 radiative transfer, scalar transport, transpiration downregulation, and solution schemes.

161 We adopt a slight modification in terminology within this LSM description section. In the main text and other sections of
162 this supplement, the transpiration flux is represented by the variable T ; however, temperature is very prevalent in the LSM
163 equations and is traditionally represented by T . To avoid confusion and be consistent with the conservation of energy in the
164 LSM, we elect to represent transpiration in energy flux units (Wm^{-2}) and represent it with the variable for latent heat flux from
165 the canopy (LE_l), where the subscript represents the big leaf approximation. Similarly the bare soil evaporation is represented
166 (LE_g), where the subscript represents the ground. Thus, the latent heat flux (LE) is the sum of canopy and ground latent heat
167 fluxes, which is simply evapotranspiration (ET) in energy units. The notation frees up the variable T to represent temperature.

168 S6.1 LSM Energy Balance

169 The energy balance of the soil-plant-atmosphere for the two big leaf, dual source LSM is shown by Equation S5. The net
170 radiation (R_n) of the soil plant system is the difference of the incoming and outgoing shortwave (S_{in} and S_{out} , respectively) and
171 longwave (L_{in} and L_{out} , respectively) radiation, i.e., the radiation absorbed by the soil-plant system. This absorbed radiation
172 is available for sensible (H), latent (LE), ground heat flux (G) and storage (not included in this formulation). We assume
173 here one-dimensional (vertical), steady-state energy transport (no energy storage) common in many LSMs. The dynamics
174 in model outputs are controlled by the change in the forcing data. The steady-state simplification turns the solution from a
175 numerical integration of a partial differential equation to numerical solutions of nonlinear equations and allows parallelization
176 in computation.

$$R_n = S_{in} - S_{out} + L_{in} - L_{out} = H + LE + G \quad (S5)$$

177 The ‘dual source’ and ‘two big leaf’ descriptors indicate how the overall energy balance is broken up into smaller
178 components. The dual source LSM structure means the surface is partitioned into plant canopy and ground components as
179 sources of scalars (illustrated in Fig. S11). Additionally, we elect the two-layer form of the dual source structure, similar to

180 CLM v5¹⁶, where both canopy and soil interact with a canopy airspace (Fig. S11), which, in turn, interacts with the atmosphere
 181 above the canopy. The two big leaf approximation further partitions the canopy component into a sunlit and shaded big leaf
 182 approximation representing the integrated fluxes of all sunlit and shaded leaves.

183 Before diving further into the energy balance of these LSM components, it is important to define some notation rules for the
 184 equations in this section that will clearly delineate the model structure. The notation rules for this LSM structure are as follows:
 185 1) a subscript of ‘*l*’ or ‘*g*’ indicates canopy/big leaf or ground fluxes, respectively, 2) an additional subscript following ‘*l*’ such
 186 as ‘*sl*’ or ‘*sh*’ indicates the sunlit or shaded big leaf component of the canopy flux, 3) the index ‘*k*’ in lieu of ‘*sl*’ or ‘*sh*’ refers
 187 to both sunlit and shaded big leaves, 4) shortwave radiation terms have an additional subscript ‘*par*’ or ‘*nir*’, identifying the
 188 specific radiation band, i.e., whether it is photosynthetically active radiation (PAR) or near infrared radiation (NIR), and 5) a
 189 single subscript ‘*Λ*’ in lieu of ‘*par*’ or ‘*nir*’ refers to both bands.

190 Using the above conventions, Equation S5 can then be further broken down into three smaller balances for the sunlit big
 191 leaf (Eqn. S6), shaded big leaf (Eqn. S7), and the soil or ground (Eqn. S8). Balancing each of these equations separately is
 192 equivalent to balancing the overall energy budget in Equation S5. Furthermore, each total flux (Eqn. S5) requires consistency
 193 between model components as shown in Equations S9-S12.

$$R_{n,l,sl} = S_{l,sl,par} + S_{l,sl,nir} + L_{l,sl} = H_{l,sl} + LE_{l,sl} \quad (S6)$$

$$R_{n,l,sh} = S_{l,sh,par} + S_{l,sh,nir} + L_{l,sh} = H_{l,sh} + LE_{l,sh} \quad (S7)$$

$$R_{n,g} = S_{g,par} + S_{g,nir} + L_g = H_g + LE_g + G_g \quad (S8)$$

$$R_n = R_{n,l,sl} + R_{n,l,sh} + R_{n,g} = R_{n,l,k} + R_{n,g} \quad (S9)$$

$$H = H_{l,sl} + H_{l,sh} + H_g = H_{l,k} + H_g \quad (S10)$$

$$LE = LE_{l,sl} + LE_{l,sh} + LE_g = LE_{l,k} + LE_g \quad (S11)$$

$$G = G_g \quad (S12)$$

194 S6.2 Radiative Transfer

195 The radiative transfer model was forced with incoming PAR, NIR and longwave radiation based on site measurements (see
 196 section S4 for details). Here we discuss the separate shortwave and longwave radiative transfer models.

197 S6.2.1 Shortwave Radiative Transfer

198 We use the Goudriaan and van Laar (GvL) model¹⁷ to estimate shortwave radiative transfer in lieu of the two-stream
 199 approximation^{16,18} used in CLM v5. Both approaches are two-stream models that focus on the upward and downward net
 200 fluxes of diffuse radiation with single scattering¹³. However, the GvL model yields simpler analytical forms and is used in other
 201 TBMs such as CABLE¹⁹. The reader is referred to^{13,17} for detailed derivation of the model. Shortwave radiation is partitioned

202 into direct beam, scattered beam, and diffuse components of PAR and NIR. The two big leaf assumption also requires assuming
 203 that diffuse leaves only receive scattered beam and diffuse radiation, while sunlit leaves receive the same as well as direct beam
 204 radiation^{13,16,20}.

205 The total canopy shortwave radiation absorption $S_{l,\Lambda}$ is given by Equation S13. This values must be partitioned appropriately
 206 between the sunlit and shaded big leaf. For ease of calculation and completeness, the sunlit leaf shortwave radiation absorption
 207 ($S_{l,sl,\Lambda}$, Eqn. S14) is partitioned into direct beam ($S_{l,sl,\Lambda,b}$, Eqn. S15), diffuse ($S_{l,sl,\Lambda,d}$, Eqn. S16), and scattered direct beam
 208 ($S_{l,sl,\Lambda,sb}$, Eqn. S17) components following²⁰. The shaded leaf shortwave absorption ($S_{l,sh,\Lambda}$, Eqn. S18) is simply the difference
 209 of total canopy absorption and sunlit leaf absorption, although analogous forms of the sunlit equations (Eqns. S15-S17) can
 210 also be used²⁰.

$$S_{l,\Lambda} = (1 - \rho'_{l,\Lambda,b}) S_{in,\Lambda,b} (1 - \exp[-K'_{b,\Lambda} \cdot LAI]) + (1 - \rho'_{l,\Lambda,d}) S_{in,\Lambda,d} (1 - \exp[-K'_{d,\Lambda} \cdot LAI]) \quad (S13)$$

$$S_{l,sl,\Lambda} = S_{l,sl,\Lambda,b} + S_{l,sl,\Lambda,d} + S_{l,sl,\Lambda,sb} \quad (S14)$$

$$S_{l,sl,\Lambda,b} = S_{in,\Lambda,b} \cdot \alpha_{l,\Lambda} \cdot (1 - \exp[-K_b \cdot LAI]) \quad (S15)$$

$$S_{l,sl,\Lambda,d} = S_{in,\Lambda,d} \cdot (1 - \rho'_{l,\Lambda,d}) \cdot (1 - \exp[-(K'_{d,\Lambda} + K_b) \cdot LAI]) \cdot \frac{K'_{d,\Lambda}}{K'_{d,\Lambda} + K_b} \quad (S16)$$

$$S_{l,sl,\Lambda,sb} = S_{in,\Lambda,b} \cdot ((1 - \rho'_{l,\Lambda,b}) \cdot (1 - \exp[-(K'_{b,\Lambda} + K_b) \cdot LAI]) \cdot \frac{K'_{b,\Lambda}}{K'_{b,\Lambda} + K_b} + \alpha_{l,\Lambda} \cdot (1 - \exp[-2K_b \cdot LAI])/2) \quad (S17)$$

$$S_{l,sh,\Lambda} = S_{l,\Lambda} - S_{l,sl,\Lambda} \quad (S18)$$

211 These radiative transfer equations rely on four essential parameters in the GvL model for shortwave radiative transfer: the
 212 direct (K_b) and diffuse extinction coefficients (K_d) and the direct ($\rho'_{l,b}$) and diffuse canopy reflectance coefficients ($\rho'_{l,d}$). The
 213 K_b value is calculated by dividing the mean leaf angle ($G(Z)$) by the projection of sunlight onto a horizontal surface (Eqn.
 214 S19), where Z is the sun zenith angle. The K_b value will change throughout the day as the sun moves across the sky since the
 215 angle of incidence with respect to leaf angles will vary. The function $G(Z)$ is known as the ‘Ross-Goudriaan’ function (Eqns.
 216 S20-S22), which depends on a parameter, χ_l , that describes the leaf angle distribution’s deviation from a spherical (i.e., random)
 217 distribution. As mentioned in section S4, we calibrated χ_l to vary between -0.4 and 0.6.

$$K_b = \frac{G(Z)}{\cos(Z)} \quad (S19)$$

$$G(Z) = \phi_1 + \phi_2 \cos(Z) \quad (S20)$$

$$\phi_1 = 0.5 - 0.633\chi_l - 0.33\chi_l^2 \quad (S21)$$

$$\phi_2 = 0.877(1 - 2\phi_1) \quad (S22)$$

218 The diffuse radiation extinction coefficient, K_d , is calculated by integrating the direct beam transmissivity ($\tau_{l,b}$ shown in
 219 Eqn. S23) over every solid angle of the hemisphere (Eqn. S24) and then inverting the transmissivity law (Eqn. S25). The
 220 transmissivity describes the attenuation of the percent of radiation that makes it through the canopy to the soil as a function of
 221 distance from the canopy, through an exponential function.

$$\tau_{l,b} = \exp(-K_b \cdot LAI) \quad (S23)$$

$$\tau_{l,d} = \int_0^{\pi/2} \tau_{l,b} \cdot \cos Z \cdot \sin Z dZ \quad (S24)$$

$$K_d = \frac{-\ln \tau_{l,d}}{LAI} \quad (S25)$$

222 The GvL model has fewer equations than the CLM v5 two-stream approximation due to several simplifying assumptions.
 223 First, the single scattering of radiation can be accounted for in the extinction coefficients (K_b and K_d) simply by multiplying by
 224 the square root of leaf absorption (α_l)¹⁷. The extinction coefficients accounting for single-scattering are shown in Equations
 225 S26-S27. Second, leaf transmissivity and reflectance are assumed identical—a reasonable assumption for green canopies¹⁷—
 226 allowing derivation of simplified relationships for direct beam ($\rho_{l,b}$, Eqn. S28) and diffuse canopy reflectance ($\rho_{l,d}$, Eqn. S29)
 227 based on idealized reflectance of horizontal leaves ($\rho_{l,h}$, Eqn. S30). Readers are referred to^{17,21} for further details on these
 228 assumptions.

$$K'_{b,\Lambda} = \sqrt{\alpha_{l,\Lambda}} \cdot K_b \quad (S26)$$

$$K'_{d,\Lambda} = \sqrt{\alpha_{l,\Lambda}} \cdot K_d \quad (S27)$$

$$\rho_{l,b} = \frac{2K_b}{K_b + K_d} \rho_{l,h} \quad (S28)$$

$$\rho_{l,d} = \int_0^{\pi/2} 2 \cdot \rho_{l,b} \cdot \cos Z \cdot \sin Z dZ \quad (S29)$$

$$\rho_{l,h} = \frac{1 - \sqrt{\alpha_{l,\Lambda}}}{1 + \sqrt{\alpha_{l,\Lambda}}} \quad (S30)$$

229 The above canopy reflectance equations were derived for infinitely deep canopies. To account for the ground reflectance
 230 (ρ_g), the approximations in Equations S31-S32 are used. These approximations assume radiation travels through the canopy,
 231 reflects off the soil according to ρ_g , and travels back up through the canopy (hence the factor of 2).

$$\rho'_{l,\Lambda,b} = \rho_{l,b} + (\rho_g - \rho_{l,b}) \exp(-2K'_b LAI) \quad (S31)$$

$$\rho'_{l,\Lambda,d} = \rho_{l,d} + (\rho_g - \rho_{l,b}) \exp(-2K'_d LAI) \quad (S32)$$

232 **S6.2.2 Longwave Radiative Transfer**

233 The longwave radiative transfer model follows the method laid out in²², which is derived assuming exponential extinction of
234 longwave radiation through the plant canopy. The net absorbed longwave radiation ($L_{l,k}$) is given by Equation S33, which
235 depends on the sunlit and shaded leaf temperature ($T_{l,k}$), ground temperature (T_g), fraction of longwave radiation absorbed
236 by the canopy (δ_l , Eqn. S34), the sunlit and shaded leaf fraction (F_k), and the Stefan-Boltzmann constant (σ). As mentioned
237 previously, k is used to indicate that the equations are identical for sunlit or shaded big leaves.

$$L_{l,k} = (L_{in} - 2\sigma T_{l,k}^4 + \sigma T_g^4) \cdot \delta_l \cdot F_k \quad (\text{S33})$$

$$\delta_l = 1 - \exp(-LAI) \quad (\text{S34})$$

$$F_{k=1} = F_{sl} = \frac{1 - \exp(-K_b \cdot LAI)}{K_b \cdot LAI} \quad (\text{S35})$$

$$F_{k=2} = F_{sh} = 1 - F_{sl} \quad (\text{S36})$$

238 **S6.3 Scalar Transport**

239 Scalar transport for this LSM consists of prognostic equations for latent heat flux (LE), sensible heat flux (H) and gross primary
240 productivity (GPP). The conserved quantities are mass of H_2O and CO_2 as well as enthalpy ($c_p \cdot T$). The states of the soil-plant
241 system are given by partial pressure of H_2O (e), partial pressure of CO_2 (c) and temperature (T). First, we will describe the
242 latent and sensible heat fluxes occurring between the canopy, ground, canopy airspace, and atmosphere. Then, we will elaborate
243 on the coupled water vapor and CO_2 transport controlled by stomatal response to varying environmental conditions.

244 The two layer approach¹³ used in this LSM splits the transport equations into canopy, ground, and atmospheric fluxes that
245 are coupled via the canopy airspace (shown in Figure S11). In effect, there are four transport pathways: 1) sunlit canopy to
246 canopy airspace, 2) shaded canopy to canopy airspace, 3) ground to canopy airspace, and 4) canopy airspace to atmosphere
247 above canopy. The first three pathways must balance with the last pathway under the steady-state conditions. All transport
248 equations use integrated flux-gradient relationships (also known as bulk transfer relations or an analogy to Ohm's law) to
249 calculate fluxes as the difference in potentials between two points in space multiplied by a conductance (inverse of resistance).
250 As previously mentioned, the index k in an equation represents either the sunlit or shaded big leaf; the forms of the equations
251 are identical, but the states experienced by the each big leaf and its respective fluxes will differ.

252 **S6.3.1 Latent and Sensible Heat Fluxes**

253 The transport of water vapor from the canopy to the canopy air space (transpiration) consists of two steps: 1) transport from
254 the leaf mesophyll cells through the stomatal opening ($LE_{l,k}$, Eqn. S37) and 2) transport through the laminar boundary layer
255 at the leaf surface to the canopy air space (Eqn. S38). The transpiration through the stomata apertures is driven by potential
256 differences in the stomatal cavity vapor pressure ($e_{i,k}$) and the vapor pressure at the surface of the leaf ($e_{s,j}$) and mediated by the
257 stomatal aperture controlled by stomatal conductance $g_{s,k}$. Likewise, the transport from the leaf surface to the canopy air space

258 is driven by the difference in $e_{s,j}$ and vapor pressure in the canopy air space (e_{ca}) and mediated by the laminar boundary layer
 259 conductance to water vapor (g_{bv}). Since we assume steady state and use Ohm's analogy to represent transport, we can treat
 260 these two pathways as two resistors in series and calculate the overall transpiration from the canopy in a single equation (Eqn.
 261 S39). Note that scaling from the individual leaf to the big leaf approximation is done simply by multiplying by sunlit or shaded
 262 leaf area index (LAI_k). This assumes that all sunlit leaves have the same stomatal conductance and internal vapor pressure as do
 263 the shaded leaves. Additionally, we apply a mass-to-energy unit conversion (C_e) consisting of the latent heat of vaporization
 264 (\mathcal{L}_v), density of air (ρ_a), ratio of molar mass of water to molar mass of air (ϵ), and atmospheric pressure (P_{atm}). For simplicity,
 265 we have assumed a constant air density and have not modified it based on water vapor concentration or temperature. The LE
 266 equation is written assuming stomata on one side of the leaf as is common practice¹³. If a plant has stomata on both sides, it is
 267 usually accounted for in the stomatal conductance measurement and parameters.

$$LE_{l,k} = LAI_k \cdot C_e \cdot g_{s,k} \cdot (e_{i,k} - e_{s,k}) \quad (S37)$$

$$LE_{l,k} = LAI_k \cdot C_e \cdot g_{bv} \cdot (e_{s,k} - e_{ca}) \quad (S38)$$

$$LE_{l,k} = LAI_k \cdot C_e \cdot \frac{g_{s,k} \cdot g_{bv}}{g_{s,k} + g_{bv}} \cdot (e_{i,k} - e_{ca}) \quad (S39)$$

$$C_e = \frac{\mathcal{L}_v \cdot \rho_a \cdot \epsilon}{P_{atm}} \quad (S40)$$

268 The description of sensible heat flux from the canopy is simpler than that of latent heat flux, as we assume no temperature
 269 gradient within a leaf. Therefore, heat transport is driven by temperature difference between the leaf ($T_{l,k}$) and canopy airspace
 270 (T_{ca}) only and mediated by the laminar boundary layer conductance to heat (g_{bh}). The result is scaled from a single leaf to the
 271 entire canopy by multiplying by the sunlit or shaded LAI as shown in Equation S41. The underlying assumption here is that all
 272 sunlit leaves have one temperature and all shaded leaves have another at each timestep. Furthermore, a conversion factor (C_h ,
 273 Eqn. S41) consisting of ρ_a and specific heat at constant pressure (c_p) is required to make the transport in terms of enthalpy
 274 which is the conserved quantity (not temperature). The factor of 2 in Equation S41 represents transport from both sides of the
 275 leaf.

$$H_{l,k} = 2LAI_k \cdot C_h \cdot g_{bh} \cdot (T_{l,k} - T_{ca}) \quad (S41)$$

$$C_h = \rho_a \cdot c_p \quad (S42)$$

276 There are four unknown conductances that must be calculated. The stomatal conductance g_s will be covered in the next
 277 section as it is coupled to carbon assimilation. The laminar boundary layer conductances for water vapor and heat are assumed
 278 identical based on Reynold's analogy¹³ and are calculated using equations derived from heat transfer experiments on rigid steel

279 leaves (Eqn. S43). The calculation requires a turbulent transfer coefficient (C_l), a characteristic leaf dimension (d_l) and the
 280 friction velocity (u_*) measured at the flux tower.

$$g_{bv} = g_{bh} = \frac{C_l \cdot u_*}{d_l} \quad (\text{S43})$$

281 Next, the transport of water and heat from the ground to the canopy airspace is shown in Equations S44-S45. Much like
 282 $LE_{l,k}$, latent heat flux from the ground (LE_g) consists of two conductances in series driven by the vapor pressure difference in
 283 ground (e_g) and canopy airspace (e_{ca}). The conductances represent vapor transport through the tortuous soil pores when soil is
 284 not saturated (g_{sv}) and the subsequent transport from the soil surface to the canopy airspace through a laminar boundary layer
 285 (g'_{av}). The sensible heat flux from the ground to canopy airspace H_g is driven by the difference in ground temperature T_g and
 286 T_{ca} mediated by conductance of heat between soil surface and canopy airspace (g'_{ah}).

$$LE_g = C_e \cdot \frac{g_{sv} \cdot g'_{av}}{g_{sv} + g'_{av}} \cdot (e_g - e_{ca}) \quad (\text{S44})$$

$$H_g = C_h \cdot g'_{ah} \cdot (T_g - T_{ca}) \quad (\text{S45})$$

287 The conductance for both heat and water vapor from the soil are again assumed equivalent by Reynold's analogy and is
 288 calculated using a turbulent transfer coefficient (C_g) and u_* as assumed in¹⁶ (Eqn. S46). The turbulent transfer coefficient
 289 is balanced between bare soil and dense canopy values using Equations S47-S49. The reader is referred to¹⁶ and references
 290 therein for justification of these parametrizations.

$$g'_{av} = g'_{ah} = C_g \cdot u_* \quad (\text{S46})$$

$$C_g = W \cdot C_{g,bare} + (1 - W) \cdot C_{g,dense} \quad (\text{S47})$$

$$W = \exp(-LAI - SAI) \quad (\text{S48})$$

$$C_{g,bare} = \frac{k}{0.13} \cdot \left(\frac{z_{0m,g} \cdot u_*}{\nu} \right)^{-0.45} \quad (\text{S49})$$

291 The additional conductance accounted for in unsaturated soils, g_{sv} , is calculated with Equation S50 using an estimate of the
 292 dry soil layer (DSL), the water vapor diffusivity (D_v) and a shape factor describing the tortuosity of the soil pores (τ). The value
 293 of g_{sv} approaches ∞ as the soil becomes saturated to an incipient level (θ_i) which was calibrated in our analysis. If g_{sv} is infinite,
 294 the conductance in Equation S44 simplifies to g'_{av} . The reader is again referred to¹⁶ and references therein for justification of
 295 these parametrizations.

$$g_{sv} = \frac{D_v \cdot \tau}{DSL} \quad (S50)$$

$$D_v = 2.12 \times 10^{-5} \cdot \left(\frac{T_g + 273.15}{273.15} \right)^{1.75} \quad (S51)$$

$$DSL = D_{max} \cdot \frac{\theta_i - \theta_s}{\theta_i - \theta_{air}} \quad (S52)$$

$$\tau = \phi_{air}^2 \cdot \left(\frac{\theta_{sat} - \theta_{air}}{\theta_{sat}} \right)^{3/b} \quad (S53)$$

296 Lastly, the latent and sensible heat fluxes from the canopy airspace to the atmosphere at the measurement point z are
 297 described in Equations S54-S55. The potential differences are between vapor pressure and temperature in the canopy airspace
 298 (T_{ca} and e_{ca}) and the atmosphere at the flux tower measurement height (T_a and e_a). The conductance from the canopy airspace to
 299 the atmosphere is again the same for heat (g_{ah}) and vapor (g_{av}) by Reynold's analogy shown in Equation S56. The conductance
 300 is based on the Monin-Obukhov similarity theory (MOST)²³ also known as the 'log-law'. The momentum roughness length
 301 (z_{om}), heat/vapor roughness length (z_{oh}), and zero-plane displacement height (d_o) are empirical parameters. The z_{om} was
 302 determined from literature while the other two parameters are calculated using practical relationships²⁴ (Eqns. S57-S58).
 303 For this study, we neglected the impact of atmospheric stability on the atmospheric conductance term. These effects are
 304 usually handled by correction factors accounting for how density stratifications in the atmosphere enhance or suppress turbulent
 305 transport. However, the stability corrections add another level of complexity to the numerical scheme, as they are dependent on
 306 H and LE , and are not important to the overall question of this research.

$$LE = C_e \cdot g_{av} \cdot (e_{ca} - e_a) \quad (S54)$$

$$H = C_h \cdot g_{ah} \cdot (T_{ca} - T_a) \quad (S55)$$

$$g_{ah} = g_{av} = \frac{\bar{u} \cdot k^2}{\ln\left(\frac{z_m - d_o}{z_{om}}\right) \cdot \ln\left(\frac{z_m - d_o}{z_{oh}}\right)} \quad (S56)$$

$$z_{oh} = 0.1z_{om} \quad (S57)$$

$$d_o = 0.7h \quad (S58)$$

307 In summary, Equations S37-S58 contain five prognostic variables: $T_{l,sl}$, $T_{l,sh}$, T_g , $g_{s,sl}$, and $g_{s,sh}$. An important assumption
 308 for scalar transport is that the vapor pressures $e_{i,k}$ and e_g are assumed to be dependent on $T_{l,k}$ and T_g via the Clausius-Clapeyron
 309 relationship¹³. Furthermore, the states of the canopy airspace, e_{ca} and T_{ca} , are completely determined by the states and
 310 conductances of the canopy, ground, and atmosphere. Substituting Equations S39, S41, S44, S45, S54 and S55 into Equations
 311 S10-S11 and solving for e_{ca} and T_{ca} yields weighted averages of the other conductances and states (Eqns. S59-S60). All other

312 terms in the scalar transport equations are either forcing data, parameters, or constants. Therefore, we have at least five variables
 313 thus far that must be solved for.

$$e_{ca} = \frac{g_{av} \cdot e_a + g_{l,sl} \cdot e_{i,sl} + g_{l,sh} \cdot e_{i,sh} + g_{av,g} \cdot e_g}{g_{av} + g_{l,sl} + g_{l,sh} + g_{av,g}} \quad (S59)$$

$$T_{ca} = \frac{g_{ah} \cdot T_a + g_{bh} \cdot T_{l,sl} + g_{bh} \cdot T_{l,sh} + g_{ah,g} \cdot T_g}{g_{ah} + 2g_{bh} + g_{ah,g}} \quad (S60)$$

$$g_{l,k} = \frac{LAI_k \cdot g_{s,k} \cdot g_{bv}}{g_{s,k} + g_{bv}} \quad (S61)$$

314 **S6.4 Stomatal Conductance and CO₂ Assimilation**

315 Stomatal conductance is intrinsically tied to CO₂ assimilation as stomatal aperture and CO₂ gradient controls photosynthetic
 316 carbon fixation. We utilize a steady state, coupled stomatal conductance-photosynthesis scheme similar to²⁵ that balances CO₂
 317 assimilation with CO₂ diffusion into the leaf. Specifically, we utilize the Medlyn stomatal conductance model²⁶ to represent
 318 stomatal conductance responses to atmospheric conditions coupled with the Farquhar, von Caemmerer, and Berry (1980) C3
 319 photosynthesis model²⁷ (hereafter, referred to as FvCB model).

320 **S6.4.1 Medlyn Stomatal Conductance Model**

321 The Medlyn stomatal conductance model (Eqn. S62) is derived assuming plants optimize the ratio of carbon gain to water lost
 322 at each instant²⁶. The solution of a resulting calculus of variations problem yields a relation where stomata close under higher
 323 vapor pressure deficit ($D_k = e_{i,k} - e_{s,k}$) and leaf surface CO₂ concentration (c_s), and open with higher CO₂ assimilation ($A_{n,k}$).
 324 This model provided a unifying framework for previously successful empirical methods²⁸. The model is parametrized by the
 325 minimum stomatal conductance (g_o) and a species-specific slope parameter (g_1) related to the marginal carbon gain to water
 326 loss.

$$g_{s,k} = g_o + \left(1 + \frac{g_1}{\sqrt{D^k}}\right) \frac{1.6A_{n,k}}{c_s/P_{atm}} \quad (S62)$$

327 The stomatal conductance model is coupled to the photosynthesis model via the term $A_{n,k}$ and the CO₂ diffusion equation
 328 (Eqn. S63). The transport of CO₂ into the leaf via diffusion is nearly identical to that of water vapor (Eqn. S39), with increases
 329 to stomatal and laminar boundary layer conductances of 1.6 and 1.4, respectively, to account for the differing diffusivities of
 330 CO₂ compared to H₂O.

$$A_{n,k}^d = \frac{g_{s,k} \cdot g_{bv}}{1.4g_{s,k} + 1.6g_{bv}} \cdot \frac{(c_{i,k} - c_{ca})}{P_{atm}} \cdot 10^6 \quad (S63)$$

331 **S6.4.2 FvCB C3 Photosynthesis Model**

332 The FcVB model²⁷ represents the three limiting mechanisms of the Calvin Cycle for steady-state carbon assimilation from
 333 atmospheric CO₂: 1) the enzyme kinetics of Ribulose 1,5 bisphosphate carboxylase-oxygenase (Rubisco), 2) the rate of
 334 Ribulose 1,5 bisphosphate (RuBP) regeneration rate governed by ATP and NADPH created in the electron transport chain of the
 335 light reactions, and 3) the amount of triose phosphates (starches) a plant can use. The equations here are for C3 photosynthesis
 336 only following¹⁶.

337 Rubisco-limitation is represented using Michaelis-Menten (MM) kinetics that describe uptake velocity of a fixed amount
 338 of Rubisco when RuBP is saturated at an internal concentration of CO₂ (Eqn. S64). The equation determines the amount
 339 of CO₂ assimilated or released depending on whether Rubisco combines RuBP with CO₂ (carboxylation) or RuBP with O₂
 340 (oxygenation). Thus, the equation requires values for partial pressure of oxygen in the leaf (o_i , Eqn. S65), MM constant for
 341 CO₂ (K_c , Eqn. S66), MM constant for O₂ (K_o , Eqn. S67), and the CO₂ compensation point (Γ , Eqn. S68).

$$A_{c,k} = V_{max25} \frac{c_{i,k} - \Gamma}{c_{i,k} + K_c (1 + o_i / K_o)} \quad (S64)$$

$$o_i = 0.209 \cdot P_{atm} \quad (S65)$$

$$K_c = 404.9 \times 10^{-6} \cdot P_{atm} \quad (S66)$$

$$K_o = 278.4 \times 10^{-3} \cdot P_{atm} \quad (S67)$$

$$\Gamma = 42.75 \times 10^{-6} \cdot P_{atm} \quad (S68)$$

342 The RuBP-limited assimilation rate (A_j , Eqn. S69), also known as the light-limited rate, describes conditions where the
 343 RuBP is limiting due to shortages in NADPH and ATP from the electron transport chain in the thylakoid of the mesophyll
 344 cells. A balance of the number of electrons required to create the required NADPH for RuBP regeneration yields Equation S69
 345 where the rate of electron transport (J) is a key quantity. The electron transport rate is itself co-limited between a maximum
 346 rate (J_{max25}) and the efficiency of photosystem II at delivering electrons (I_{PSII} , Eqn. S70) from the absorbed PAR by the leaf
 347 ($S_{l,k,par}$). The factor of 4.6 in Equation S70 represents unit conversion to quanta of energy²⁹. The quantum efficiency of
 348 photosystem II (Φ_{PSII}) is usually taken to be 0.7¹⁶.

$$A_{j,k} = J \frac{c_{i,k} - \Gamma}{4c_{i,k} + 8\Gamma} \quad (S69)$$

$$I_{PSII,k} = 0.5\Phi_{PSII} (4.6 \cdot S_{l,k,par}) \quad (S70)$$

$$\Phi_{PSII} \cdot J^2 - (I_{PSII,k} + J_{max25}) \cdot J + I_{PSII,k} \cdot J_{max25} = 0 \quad (S71)$$

349 The product-limited assimilation rate (A_p , Eqn. S72) represents the upper limit on assimilation based on the plant's need for
 350 the sugars. See¹⁶ and sources within for justifications of the relationship with V_{max25} .

$$A_p = V_{max25}/6 \quad (S72)$$

351 Altogether, we want to calculate what the co-limitation of these three controls on the CO₂ assimilation of a plant. To do
 352 this, we use quadratic equations to estimate the co-limitation as laid out in³⁰ to allow a gradual transition across the three
 353 mechanisms and to account for additional limitations, rather than explicitly calculating any of the mechanisms separately. The
 354 Θ_{cj} and Θ_{ip} are empirical curvature factors that control for this gradual transition, given in¹⁶. The overall CO₂ assimilation A_k
 355 is given by the root of Equations S73 and S74. Lastly, we must remove from A_k the amount of CO₂ that is released through dark
 356 respiration R_d to get the overall net assimilation $A_{n,k}$ (Eqn. S75). $A_{n,k}$ is the amount of CO₂ assimilated from the atmosphere,
 357 balanced with leaf diffusion (Eqn. S60) for a big leaf as mediated by stomatal conductance.

$$\Theta_{cj} \cdot A_{i,k}^2 - (A_{c,k} + A_{j,k}) \cdot A_{i,k} + A_{c,k} \cdot A_{j,k} = 0 \quad (S73)$$

$$\Theta_{ip} \cdot A_k^2 - (A_{i,k} + A_{p,k}) \cdot A_k + A_{i,k} \cdot A_{p,k} = 0 \quad (S74)$$

$$A_{n,k} = A_k - R_d \quad (S75)$$

$$R_d = 0.015 \cdot V_{max25} \quad (S76)$$

358 For simplicity, we have omitted the temperature dependence of the photosynthetic parameters V_{max25} , J_{max25} , R_d , K_c , K_o ,
 359 and Γ and simply use the values at 25°C^{31–33}. These dependencies are typically handled with an Arrhenius functions¹³ to
 360 account for the breakdown or acceleration of various metabolic processes at high and low temperatures. Since the goal of this
 361 paper was to test the transpiration downregulation schemes, we omitted the temperature dependence due to the need for many
 362 more parameters to properly use the Arrhenius functions.

363 **S6.4.3 Scale Correction of Photosynthetic Parameters**

364 The maximum carboxylation rate of the Rubisco enzyme (V_{max25}) and the maximum electron transport rate (J_{max25}) are
 365 dependent on nitrogen availability in the leaf. Nitrogen has been found to exponentially decay with relative cumulative
 366 leaf area in the canopy³⁴; therefore, both V_{max25} and J_{max25} vary nonlinearly with plant height. For simplicity, we follow²⁰ for
 367 scaling V_{max25} and²² for scaling J_{max25} to account for this nonlinear nitrogen profile. These methods differ from the optimality
 368 principles used in CLM v5¹⁶.

369 The overall Rubisco carboxylation capacity of the canopy ($V_{l,max25}$) factoring in leaf nitrogen is give in Equation S77, where
 370 K_n is the extinction coefficient for leaf nitrogen content. The two big leaf model requires separate consideration of the sunlit
 371 and shaded big leaf³⁵ shown in Equations S77-S79. The maximum electron transport rate of the canopy ($J_{l,max25}$) factoring in

372 leaf nitrogen is give in Equation S80, while the sunlit and shaded big leaf values are shown in Equations S81-S82. The values
 373 of $V_{l,k,max25}$ and $J_{l,k,max25}$ are used in place of the V_{max25} and J_{max25} parameters for the FvCB model described in the previous
 374 section.

$$V_{l,max25} = LAI \cdot V_{max25} \cdot [1 - \exp(-K_n)] \quad (S77)$$

$$V_{l,sl,max25} = LAI \cdot V_{max} \cdot \frac{1 - \exp(-K_n - K_b \cdot LAI)}{K_n + K_b \cdot LAI} \quad (S78)$$

$$V_{l,sh,max25} = V_{l,max25} - V_{l,sl,max25} \quad (S79)$$

$$J_{l,max25} = J_{max25} \cdot \frac{1 - \exp(-K'_d \cdot LAI)}{K'_d} \quad (S80)$$

$$J_{l,sl,max25} = J_{max25} \cdot \frac{1 - \exp(-[K'_d + K_b] \cdot LAI)}{K'_d + K_b} \quad (S81)$$

$$J_{l,sh,max25} = J_{l,max25} - J_{l,sl,max25} \quad (S82)$$

375 S6.5 Transpiration Downregulation

376 the transpiration downregulation schemes used in the paper are the empirical β and Plant Hydraulic Model schemes (PHM).
 377 We will discuss how each is implemented to suppress transpiration under soil water stress. The reader is referred to the main
 378 article for detailed discussion on the theoretical justification for the two methods.

379 S6.5.1 Well-Watered Transpiration

380 Before discussing the transpiration downregulation schemes, we must first clarify the terminology ‘well-watered’. As stated
 381 in the main article, well-watered refers to soil water conditions that do not cause any limitation to transpiration through
 382 stomatal closure via low leaf water potential. In other words, the transpiration meets the stomata-regulated atmospheric
 383 moisture demand—determined by the Medlyn model (Eqn. S62) and the vapor pressure deficit (D). This definition becomes
 384 slightly more ambiguous as we introduce a dual source, two big leaf model structure, as the states (vapor pressure and
 385 temperature) experienced by the hypothetical big leaves at a time step adjust to downregulation. Therefore, for clarity, the
 386 well-watered transpiration rate is the value corresponding to the states when there is no stomatal closure from soil moisture
 387 effects. Computationally, this is simple, since the well-watered rate is the LSM output when downregulation is turned off. This
 388 approach differs from CLM v5¹⁶, which considers well-watered transpiration to be the rate under the downregulated states.
 389 This distinction between the two definitions of the well-watered rate will become important shortly, as the well-watered rate is a
 390 key variable in the transpiration downregulation schemes. Also, note that the well-watered rate is different between sunlit and
 391 shaded big leaf as they encounter differing temperatures, light, and vapor pressures.

392 **S6.5.2 β Downregulation Schemes**

393 As mentioned in the main article and in section S2, the LSM utilizes a Weibull function to represent the empirical β curve (Eqn.
394 16 in the main article). There are three variants of this method used: 1) a single β , 2) a 2-leaf β , and 3) a dynamic β . Since the
395 method is empirical, there is not firm guidance on where within the plant to apply this downregulation, as some models apply it
396 directly to well-watered stomatal conductance and other apply it to photosynthetic parameters like V_{max25} . Here, we apply β to
397 the well-watered transpiration rate of the sunlit and shaded big leaf to maintain consistency with our minimalist analysis. The
398 solution scheme section will discuss in greater detail how β is applied.

399 **S6.5.3 PHM Downregulation Scheme**

400 We will elaborate here on the PHM laid out in the main article and extend its formulation to the two-big leaf approach of the
401 LSM. The PHM is similar to that in CLM v5^{16,36}, with simplification to soil-to-xylem, xylem-to-leaf, and leaf-to atmosphere
402 segmentation. For readability the equations shown in the main article are repeated here. Each segment has a conductance curve
403 that downregulates from the maximum conductance values based on water potentials through the segment. The conductivity
404 equations follow closely the work of^{4,37} and references therein. All parameter values and units for the following equations can
405 be found in Table S7.

406 The soil-to-xylem conductance (g_{sx} , Eqn. S83) consists of the well-known unsaturated hydraulic conductivity curve for soil⁹
407 and a maximum conductance value ($g_{sx,max}$, Eqn. S84). The downregulation function is parametrized by saturated soil water
408 potential (ψ_{sat}), soil water retention exponent (b), unsaturated hydraulic conductivity exponent ($c = 2b + 3$), and a correction
409 factor ($d = 4$) to account for roots' ability to reach water³⁸. The $g_{sx,max}$ value is calculated using the saturated hydraulic
410 conductivity ($k_{s,sat}$), specific weight of water ($\rho_w \cdot g$) and a length scale based on root area index (RAI), fine root diameter (d_r)
411 and effective rooting depth (Z_r) to convert to conductance.

$$g_{sx}(\psi) = g_{sx,max} \cdot \left(\frac{\psi_{sat}}{\psi} \right)^{\frac{c-d}{b}} \quad (S83)$$

$$g_{sx,max} = \frac{k_{s,sat}}{\rho_w g} \sqrt{\frac{RAI}{d_r Z_r}} \cdot 10^{-6} \quad (S84)$$

412 The xylem-to-leaf conductance, g_{xl} (Eq. S85), is the maximum xylem-to-leaf conductance ($g_{xl,max}$, Eqn. S85) downregulated
413 by a sigmoidal function³⁹ parametrized by the vulnerability exponent a and the xylem water potential (ψ_x) at 50% loss of
414 conductance ($\psi_{x,50}$). The $g_{xl,max}$ is estimated using sapwood hydraulic conductivity (K_{sap}), and the height of vegetation (h_v).

$$g_{xl}(\psi) = g_{xl,max} \cdot \left[1 - \frac{1}{e^{a(\psi - \psi_{x,50})}} \right] \quad (S85)$$

$$g_{xl,max} = \frac{K_{sap}}{h_v \cdot \rho_w} \quad (S86)$$

415 The leaf-to-atmosphere conductance (Eq. S87) is the stomatal conductance for the sunlit and shaded leaf, $g_{s,k}$, downregulated
 416 from its well-watered value ($g_{s,ww,k}$) using a Weibull function parametrized by a shape factor (b_l) describing stomatal sensitivity
 417 and the leaf water potential at 50% loss of conductance ($\psi_{l,50}$)⁸. The $g_{s,ww,k}$ value is calculated using the Medlyn model
 418 previously discussed in Equation S59²⁶. The values for stomatal conductance are defined for both sunlit and shaded leaf by
 419 index k as they will almost always differ.

$$g_{s,k} = g_{s,ww,k} \cdot e^{-\left(\frac{\psi_{l,k}}{\psi_{l,50}}\right)^{b_l}} \quad (\text{S87})$$

420 In order to calculate the water flux through each segment, we must utilize a Kirchhoff transform (Eqn. S88) to account
 421 for the the varying potential (and conductance) along each segment⁴⁰. The transform is only performed on the soil-to-xylem
 422 and xylem-to-leaf segments as the distance traveled through the leaf to stomata is assumed negligible. The total flux potential
 423 for soil-to-xylem ($\Phi_{sx}(\psi)$, Eqn. S89) and xylem-to-leaf ($\Phi_{xl}(\psi)$, Eqn. S90) give an upper limit on the water that could be
 424 extracted from a segment based on the potential. Using this linearized flow theory, the flux through each segment is simply
 425 calculated by taking the difference in total flux potential between the end points of each segment.

$$\Phi(\psi) = \int_{-\infty}^{\psi} K(\psi') d\psi' \quad (\text{S88})$$

$$\Phi_{sx}(\psi) = \frac{b \cdot g_{sx,max} \cdot \psi}{b - c + d} \cdot \left(\frac{\psi_{sat}}{\psi}\right)^{\frac{c-d}{b}} \quad (\text{S89})$$

$$\Phi_{xl}(\psi) = g_{xl,max} \cdot \left[\frac{\ln(e^{-a\psi} + e^{-a\psi_{50}})}{a} + \psi \right] \quad (\text{S90})$$

426 The two-big leaf configuration of this model requires five total segments: soil-to-xylem, xylem-to-sunlit leaf, xylem-to-
 427 shaded leaf, sunlit leaf-to-atmosphere, and shaded leaf-to-atmosphere. The underlying assumption is that the transport from
 428 xylem to the sunlit and shaded leaf is completely independent. The transport in each segment is shown below in Equations
 429 S91-S93. Note these equations are the same as Equations 15-17 in the main article except adapted for the two big leaf
 430 configuration.

$$LE_{sx} = [\Phi_{sx}(\psi_s) - \Phi_{sx}(\psi_x)] \cdot \rho_w \cdot \mathcal{L}_v \quad (\text{S91})$$

$$LE_{xl,k} = [\Phi_{xl}(\psi_x) - \Phi_{xl}(\psi_{l,k})] \cdot \rho_w \cdot \mathcal{L}_v \quad (\text{S92})$$

$$LE_{la,k} = LAI_k \cdot g_{s,k} \cdot (e_{i,k} - e_{s,k}) \cdot C_e \quad (\text{S93})$$

431 We assume a steady-state solution where the supply through the soil-plant system equals the atmospheric moisture demand.
 432 This problem can be solved using a Newton-Raphson method as done in CLM v5¹⁶. However, this method was found to be

433 unstable under certain conditions; therefore, we opted to use nonlinear least squares in MATLAB (*lsqnonlin*) to solve the
 434 problem. We used the trust-region-reflective method which is a quasi-Newton method that handles bounded constraints on the
 435 decision variable. The optimization problem is laid out in Equations S94-S96. The decision variables (ψ , Eqn. S96) are xylem,
 436 sunlit leaf, and shaded leaf water potentials that balances the flow through all segments while the residuals (R , Eqn. S95) ensure
 437 this balance with values of 0. The constraints in the optimization problem keep the water potentials between $\psi_{s,sat}$ and a value
 438 of -30 MPa.

$$\psi^* = \min_{\psi} \|R\|^2 \quad (S94)$$

$$\text{s.t. } \psi \in (-30, \psi_{s,sat})$$

$$R = \begin{bmatrix} LE_{sx} - \sum_{k=1}^2 LE_{xl,k} \\ LE_{xl,sl} - LE_{la,sl} \\ LE_{xl,sh} - LE_{la,sh} \end{bmatrix} \quad (S95)$$

$$\psi = \begin{bmatrix} \psi_{l,sl} \\ \psi_{l,sh} \\ \psi_x \end{bmatrix} \quad (S96)$$

439 S6.6 LSM Solution Scheme

440 There are numerous ways to solve the steady-state dual source scheme depending on how the equations and unknowns have
 441 been defined. Here, we have created our own method, similar to CLM v5. There are two overall computational schemes or
 442 solvers: a well-watered solver and a transpiration downregulation solver. In the well-watered scheme, there are two levels of
 443 computation: the surface energy budget solver (outer solver) and the scalar transport solver (inner solver). For the transpiration
 444 downregulation scheme, well-watered solutions are adjusted in a separate solver based on soil moisture availability. Our
 445 solvers use optimization routines rather than the linearized, Newton-Raphson methods used in CLM v5 for several reasons:
 446 1) numerical derivatives are required for both methods, 2) the optimization routine guards against solution divergence, 3) the
 447 optimization routine is simple to set up, and 4) speed between the two methods at our scale is essentially the same.

448 S6.6.1 Well-Watered Solver

449 The well-watered solver is the primary solution scheme of the LSM, which is run for every simulation with and without
 450 transpiration downregulation. The solver consists of two nested least squares optimization problems, which have been referred
 451 to as the outer and inner solvers for simplicity. There are six overall state variables that must be adjusted to balance the surface
 452 energy budget (Eqn. S5) for this steady-state problem: $T_{l,k}$, T_g , $c_{i,k}$ and e_{ca} . The outer solver is concerned with balancing the

453 surface energy budget by finding the correct leaf ($T_{l,k}$) and ground (T_g) temperatures, whereas the inner solver is focused on
 454 finding the correct internal leaf carbon concentrations ($c_{i,k}$) and canopy water vapor pressure (e_{ca}) that balance the LE and H
 455 leaving the ground and canopy with the transport from the canopy airspace to atmosphere.

456 The outer solver is a three dimensional nonlinear least squares problem shown in Equations S97-S99. The residuals being
 457 minimized (R^o) are the sunlit big leaf, shaded big leaf, and ground energy balances in Equations S6-S8, while the decision
 458 variables (T) are the temperatures of these three respective compartments. The outer solver is illustrated in (Fig. S12) as it
 459 begins by gathering all the environmental forcing data for a particular time step (section S4). The outer solver then initiates a
 460 guess for the three temperatures based on the air temperature. The next step is to solve the GvL radiative transfer model to
 461 obtain the net radiation R_n for the three compartments and their breakdown into PAR, NIR, and longwave components. At this
 462 point, the temperatures are sent to the inner solver to determine the scalar fluxes from the ground, canopy, and canopy airspace
 463 under these fixed temperatures and states. Once the inner solver finds the $c_{i,k}$ and e_{ca} that balances Equations S10-S11, the
 464 scalar fluxes for all compartments are calculated. The outer solver then checks to see if the net radiation in each compartment
 465 equals the scalar fluxes. If not, the temperatures are adjusted based on the optimization routine and the process is repeated until
 466 convergence.

$$T^* = \min_T \|R^o\|^2 \quad (S97)$$

s.t. $T \in (0, 40)$

$$R^o = \begin{bmatrix} S_{l,sl,par} + S_{l,sl,nir} + L_{l,sl} - H_{l,sl} - LE_{l,sl} \\ S_{l,sh,par} + S_{l,sh,nir} + L_{l,sh} - H_{l,sh} - LE_{l,sh} \\ S_{g,par} + S_{g,nir} + L_g - H_g - LE_g - G_g \end{bmatrix} \quad (S98)$$

$$T = \begin{bmatrix} T_{l,sl} \\ T_{l,sh} \\ T_g \end{bmatrix} \quad (S99)$$

467 The inner solver is also a three dimensional nonlinear least squares problem within the outer solver shown in Equations
 468 S100-S102. The inner solver is given temperatures and states of the two big leaves, ground, and air and must find the internal
 469 CO₂ concentrations that balance plant carbon synthesis with leaf diffusion, as well as the canopy water vapor pressure that
 470 balances transport from ground and plants with that to the atmosphere. The inner solver is shown in Figure S12 as the light gray
 471 indented panels. First, values of $c_{i,k}$ and e_{ca} are guessed based on atmospheric conditions. Then the FvCB model is solved to
 472 calculate the net CO₂ assimilation of each leaf ($A_{n,k}$), which must be matched by the Medlyn stomatal conductance model and
 473 leaf diffusion. A neat trick introduced in CLM v5¹⁶ is to substitute the diffusion equation into the Medlyn equation to obtain a

474 quadratic equation whose larger root is the solution for $g_{s,k}$ (Eqns. S103-S105). Using $g_{s,k}$, the internal carbon concentration
 475 from leaf diffusion ($c_{i,k}^+$) is calculated and checked against the assumed value of the solver $c_{i,k}$. Once $g_{s,k}$ has been determined,
 476 we can use Equation S59 to calculate a check on the canopy airspace water vapor pressure (e_{ca}^+). These values are adjusted by
 477 the optimization routine until convergence criteria is met. The results are then sent back out to the outer solver.

$$x^* = \min_x \|R^i\|^2 \quad (\text{S100})$$

s.t. $x \in (0, 40)$

$$R^i = \begin{bmatrix} c_{i,sl}^+ - c_{i,sl} \\ c_{i,sh}^+ - c_{i,sh} \\ e_{ca}^+ - e_{ca} \end{bmatrix} \quad (\text{S101})$$

$$x = \begin{bmatrix} c_{i,sl} \\ c_{i,sh} \\ e_{ca} \end{bmatrix} \quad (\text{S102})$$

$$g_{s,j}^2 - \left[2 \cdot g_o + 2 \cdot C_{1,j} + \frac{C_{1,j}^2 \cdot g_1^2}{g_{bv} \cdot C_{2,j}} \right] g_{s,j} + \left[g_o^2 + 2 \cdot C_{1,j} \cdot g_o + C_{1,j}^2 \left(1 - \frac{g_1^2}{C_{2,j}} \right) \right] = 0 \quad (\text{S103})$$

$$C_{1,j} = \frac{1.6 \cdot A_{n,j} \cdot P_{atm}}{c_{s,j} \cdot 10^6} \quad (\text{S104})$$

$$C_{2,j} = \frac{e_{i,j} - e_{ca}}{1000} \quad (\text{S105})$$

478 **S6.6.2 Transpiration Downregulation Solver**

479 The transpiration downregulation solver is an additional solver used after the well-watered solver to account for the effect of
 480 soil water stress on stomatal conductance and, in turn, on the scalar fluxes and plant microclimate. The solver scheme is a single
 481 least squares problem (Eqn. S106) in five dimensions of leaf temperatures and conductances as well as ground temperature
 482 (Eqn. S108). As in the well-watered solver, the first three residuals are the surface energy balance for the big leaves and
 483 ground (Eqn. S107). The final two residuals (Eqn. S109) ensure that the transpiration from the canopy calculated by the scalar
 484 transport module match the value calculated by the selected downregulation method; either β or PHM. For the β method, the
 485 downregulated transpiration is simply β multiplied by the well-watered transpiration rate $LE_{l,k,ww}$. For the the PHM method,
 486 the downregulated transpiration rate ($LE_{l,k,phm}$) is the solution to the PHM that balances supply and demand (Eqns. S94-S96).

487 The solver scheme is laid out in Figure S13 where it initializes the five decision variables from the well-watered solution.
 488 For the set temperatures and conductances we are able to re-calculate the longwave radiation, carbon assimilation, scalar fluxes

489 and states. At this point, we can calculate the the surface energy budget residuals in Equation S107. Now there is a choice
 490 to make whether to select the β model or the PHM. The β model is less computationally expensive as we simply multiply β
 491 by the already calculated $LE_{l,k,ww}$. Any of the three β methods (β_s , β_{2L} , and β_{dyn}) can be applied at this point as there is no
 492 real computational difference between the three, just different β values are multiplied by the well-watered rates. The PHM
 493 scheme is slightly more complex as we must solve the three-dimensional least squares problem to balance supply and demand.
 494 However, both schemes then check the last two residuals (Eqn. S109) to ensure the transpiration from the scalar transport
 495 module (Eqn. S39) match the downregulation scheme transpiration. If the residual does not converge the solver adjusts the
 496 decision variable and repeats.

497 This transpiration downregulation scheme is different than that proposed by CLM v5¹⁶ not only numerically but also in how
 498 the well-watered transpiration is defined. As seen in our scheme (Fig. S13), our well-watered transpiration is fixed during the
 499 scheme. We opted for this because the states of the plant microclimate under well-watered conditions are different then under
 500 downregulation. Therefore, under the same atmospheric forcings our method is consistent with what we would expect to see if
 501 the soil was saturated compared to when it is dry. The method in CLM v5 continually updates the well-watered transpiration
 502 during the downregulation solver. Essentially, as the microclimate states change during downregulation, CLM v5 re-calculates
 503 the well-watered stomatal conductance according to the Medlyn model and uses that in the downregulation schemes. This
 504 creates a positive feedback that increases transpiration suppression compared to our method. Also, the well-watered transpiration
 505 rate calculated in this method is the value that would be experienced in a certain plant microclimate and not necessarily under
 506 the atmospheric forcings. It is difficult to determine which method is most realistic, but they give very different values for
 507 downregulation. We think our definition of well-watered transpiration is more appropriate to defining the stomata-regulated
 508 atmospheric moisture demand so that is what was used in this analysis.

$$x^* = \min_x \|R^t\|^2 \quad (S106)$$

s.t. $x \in (0, 40)$

$$R^t = \begin{bmatrix} S_{l,sl,par} + S_{l,sl,nir} + L_{l,sl} - H_{l,sl} - LE_{l,sl} \\ S_{l,sh,par} + S_{l,sh,nir} + L_{l,sh} - H_{l,sh} - LE_{l,sh} \\ S_{g,par} + S_{g,nir} + L_g - H_g - LE_g - G_g \\ R_{(4)}^t \\ R_{(5)}^t \end{bmatrix} \quad (S107)$$

$$x = \begin{bmatrix} T_{l,sl} \\ T_{l,sh} \\ T_g \\ g_{l,sl} \\ g_{l,sh} \end{bmatrix} \quad (\text{S108})$$

$$R'_{(4,5)} = \begin{cases} LE_{l,k,phm} - LE_{l,sl} & \text{if PHM scheme} \\ LE_{l,k} - \beta \cdot LE_{l,k,ww} & \text{if } \beta \text{ scheme} \end{cases} \quad (\text{S109})$$

Surface Energy Budget

$$R_n = S_{in} - S_{out} + L_{in} - L_{out} = H + LE + G$$

Radiation Model: Goudriaan and Van Laar (1994)

Photosynthesis Model: Farquhar (1980)

Plant Hydraulics Model: Similar to Manzoni (2013)

Model Structure: Dual source, 2-Big Leaf Approximation

Model Forced With

- Vapor Pressure Deficit
- Air Temperature
- Soil Moisture
- Ground Heat Flux
- Streamwise Velocity
- Incoming Shortwave Radiation
- Incoming Longwave Radiation
- CO₂ Concentration
- Air Pressure

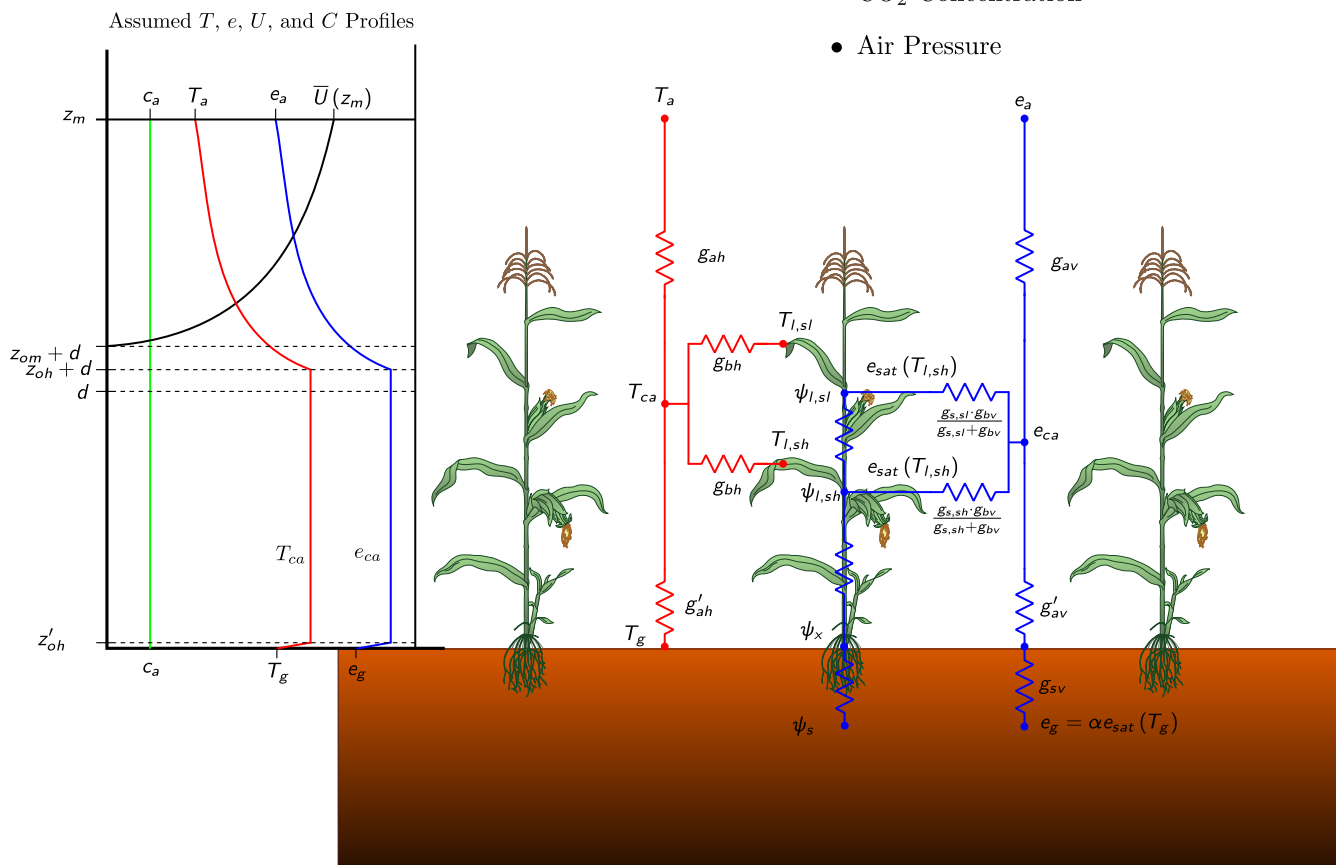


Figure S11. Schematic of our two big leaf, dual-source land surface model. The potentials and resistors indicate the scalar transport between the sunlit and shaded big leaf approximations, ground, canopy airspace and atmosphere. To the left are the assumed profiles of water vapor pressure deficit e , temperature (T), CO₂ partial pressure (c), and streamwise mean velocity (\bar{U}). The main modules used are laid out in text as well as the environmental forcings used from the US-Me2 Ameriflux site for our simulations.

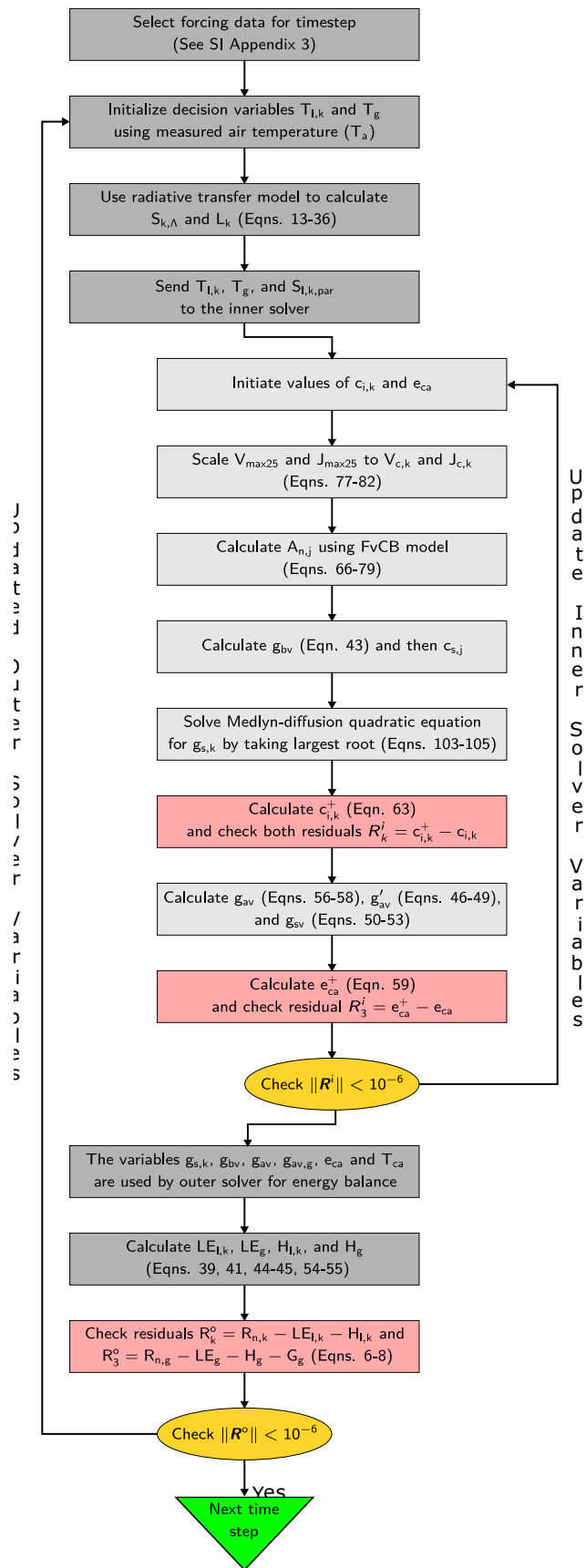


Figure S12. The well-watered solver solution scheme representing the outer solver (dark gray) and inner solver (light gray). Light red panels indicate a step where a residual to the nonlinear least squares problem is calculated and yellow indicates checking values of the residuals.

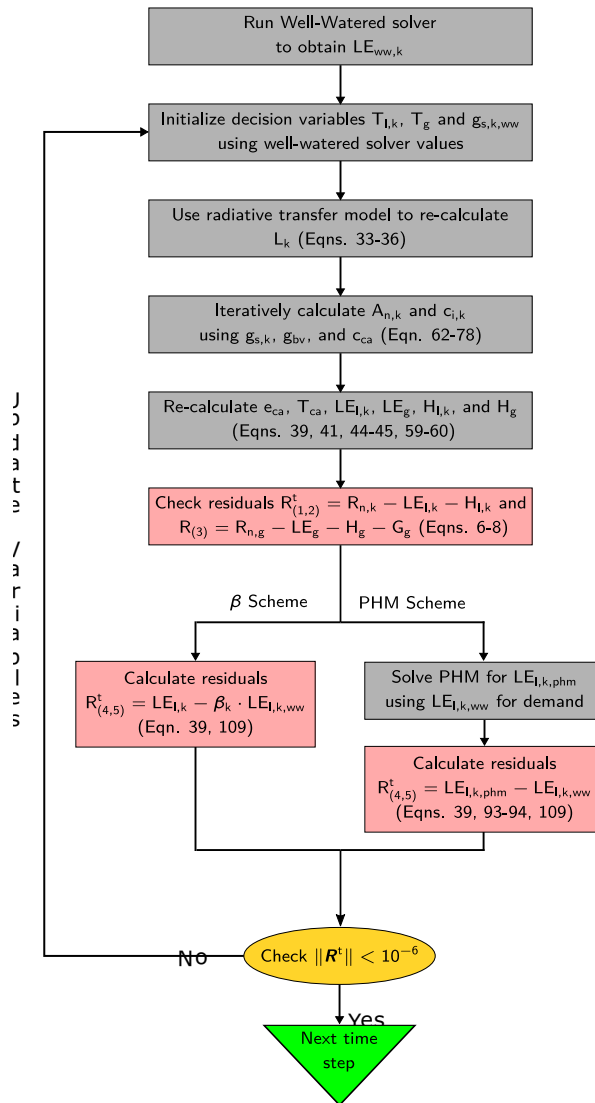


Figure S13. The transpiration downregulation scheme that is used after the well-watered solver to re-calculate fluxes and states as plants reduce transpiration from soil water stress. Light red panels indicate a step where a residual to the nonlinear least squares problem is calculated and yellow indicates checking values of the residuals. There are two separate choices for downregulation: the β model and the Plant Hydraulic Model (PHM). See text for more details.

509 **S7 LSM Variables, Parameters, and Forcings**

510 The sheer volume of equations and data discussed in this supplemental materials make it necessary to provide a comprehensive
511 table of variables, parameters, and constants with sources where necessary. This table has been split up based on the sections
512 describing the LSM: radiative transfer (Table S4), scalar transport (Table S5), coupled stomatal conductance and photosynthesis
513 (Table S6), transpiration downregulation (Table S7), and constants (Table S8). For each table, except the constants, the table is
514 broken down into subscripts, fluxes and states, forcing data, and parameters. The ‘subscripts’ section is used to cut down on
515 table entries as many subscripts are used on fluxes and parameters to describe their position in the the dual source, two big leaf
516 framework; the variable names are shown in the specific sections. The ‘fluxes and states’ section shows the main fluxes and
517 states used in the section without all the positional subscripts. The ‘forcing data’ section highlights the US-Me2 site data used
518 to force the model discussed in section S4. The ‘parameters’ section contains all the functional and constant parameters used
519 along with values and sources if constant.

Table S4. The main fluxes, states and parameters used by the radiative transfer module of the LSM.

Name	Description	Value	Units	Sources
<u>Subscript</u>				
<i>l</i>	Plant canopy	-		
<i>sl</i>	Sunlit big leaf	-		
<i>sh</i>	Shaded big leaf	-		
<i>k</i>	Sunlit or shaded big leaf	-		
<i>par</i>	Photosynthetically active radiation (PAR)	-		
<i>nir</i>	Near infrared radiation (NIR)	-		
Λ	PAR or NIR	-		
<i>b</i>	Direct beam radiation	-		
<i>d</i>	Diffuse radiation	-		
<i>sb</i>	Scattered beam radiation	-		
<i>in</i>	Incoming radiation	-		
<i>out</i>	Outgoing radiation	-		
<u>Fluxes and States</u>				
<i>S</i>	Absorbed shortwave radiation	-	$\text{W}\cdot\text{m}^{-2}$	
<i>L</i>	Absorbed longwave radiation	-	$\text{W}\cdot\text{m}^{-2}$	
<i>T</i>	Temperature	-	$^{\circ}\text{C}$	
<u>Forcing Data</u>				
<i>S_{in}</i>	Incoming shortwave radiation	-	$\text{W}\cdot\text{m}^{-2}$	41
<i>S_{in,par}</i>	Incoming PAR	-	$\text{W}\cdot\text{m}^{-2}$	
<i>S_{in,par,d}</i>	Diffuse incoming PAR	-	$\text{W}\cdot\text{m}^{-2}$	
<i>L_{in}</i>	Incoming longwave radiation	-	$\text{W}\cdot\text{m}^{-2}$	
<i>T_a</i>	Air temperature at measurement height	-	$^{\circ}\text{C}$	
<u>Parameters</u>				
<i>K</i>	Extinction coefficient	-	-	
<i>K'</i>	Extinction coefficient corrected for single-scattering	-	-	
$\alpha_{l,par}$	PAR leaf absorption coefficient	0.74	-	Calibrated
$\alpha_{l,nir}$	NIR leaf absorption coefficient	0.43	-	Calibrated
<i>LAI</i>	Leaf area index	3.2	$\text{m}^2 \text{ leaf area}\cdot\text{m}^{-2} \text{ ground area}$	Calibrated
τ	Transmissivity	-	-	
<i>G(Z)</i>	Mean leaf angle	-	radians	
<i>Z</i>	Solar zenith angle	-	radians	
χ_l	Leaf angle distribution parameter	0.11	-	Calibrated
$\rho_{l,h}$	Leaf reflectance for infinite horizontal canopy	-	-	
ρ_l	Plant canopy reflectance for infinite canopy	-	-	
ρ'_l	Plant canopy reflectance accounting for ground reflectance	-	-	
$\rho_{g,par}$	PAR ground reflectance	0.1	-	29
$\rho_{g,nir}$	NIR ground reflectance	0.2	-	29
δ_l	Fraction of longwave radiation absorbed by canopy	-	-	
<i>F_k</i>	Fraction of sunlit or shaded leaf area index	-	-	

Table S5. The main fluxes, states and parameters used by the scalar transport module of the LSM.

Name	Description	Value	Units	Sources
<u>Subscript</u>				
<i>l</i>	Plant canopy	-		
<i>sl</i>	Sunlit big leaf	-		
<i>sh</i>	Shaded big leaf	-		
<i>k</i>	Sunlit or shaded big leaf	-		
<i>i</i>	Inside the stomatal cavity of the leaf	-		
<i>s</i>	On the leaf surface	-		
<i>g</i>	Ground/soil	-		
<i>ca</i>	Canopy airspace	-		
<i>a</i>	Atmosphere above canopy at measurement height	-		
<u>Fluxes and States</u>				
<i>LE</i>	Latent heat flux	-	W·m ⁻²	
<i>H</i>	Sensible heat flux	-	W·m ⁻²	
<i>e</i>	Water vapor pressure	-	Pa	
<i>T</i>	Temperature	-	°C	
<i>c</i>	CO ₂ partial pressure	-	Pa	
<u>Forcing Data</u>				
<i>ū</i>	Mean streamwise velocity	-	m·s ⁻¹	41
<i>u*</i>	Friction velocity	-	m·s ⁻¹	
<i>θ_s</i>	Soil water content at 50 cm depth	-	m ³ water·m ⁻³ soil	
<i>e_a</i>	Water vapor pressure at measurement height	-	Pa	
<i>T_a</i>	Water vapor pressure at measurement height	-	°C	
<i>G</i>	Ground heat flux	-	W·m ⁻²	
<u>Parameters</u>				
<i>g_s</i>	Stomatal conductance	-	mol H ₂ O·m ⁻² ·s ⁻¹ or m·s ⁻¹	
<i>g_{bv}</i> or <i>g_{bh}</i>	Leaf laminar boundary layer water vapor/heat conductance	-	m·s ⁻¹	
<i>g_{av}</i> or <i>g_{ah}</i>	Atmospheric water vapor/heat conductance	-	m·s ⁻¹	
<i>g_{sv}</i>	Soil pore to soil surface water vapor conductance	-	m·s ⁻¹	
<i>g'_{av}</i> or <i>g'_{ah}</i>	Soil to canopy airspace water vapor/heat conductance	-	m·s ⁻¹	
<i>LAI</i>	Leaf area index	3.2	m ² leaf area·m ⁻² ground area	Calibrated
<i>SAI</i>	Stem area index	0.5	m ² stem area·m ⁻² ground area	10
<i>C_l</i>	Leaf turbulent transfer coefficient	0.01	m·s ⁻¹	16
<i>d_l</i>	Characteristic leaf dimension	0.04	m	16
<i>C_g</i>	Ground turbulent transfer coefficient	-	m·s ⁻¹	
<i>C_{g,bare}</i>	Bare ground turbulent transfer coefficient	-	m·s ⁻¹	
<i>C_{g,dense}</i>	Dense canopy ground turbulent transfer coefficient	0.004	m·s ⁻¹	16
<i>z_{om}</i>	Atmospheric momentum roughness length	1	m	24
<i>d_o</i>	Zero-plane displacement	-	m	
<i>z_{ov}</i> or <i>z_{ov}</i>	Atmospheric water vapor/heat roughness length	0.1	m	24
<i>z_{om,g}</i>	Ground momentum roughness length	0.01	m	16
<i>D_v</i>	Water vapor diffusivity	-	m ² ·s ⁻¹	
<i>DSL</i>	Depth of dry soil layer	-	m	
<i>D_{max}</i>	Maximum dry layer thickness	0.015	m	
<i>θ_{sat}</i>	Saturated soil water content (porosity)	0.57	m ³ water·m ⁻³ soil	42
<i>θ_i</i>	Soil water content where <i>g_{sv}</i> begins	0.57	m ³ water·m ⁻³ soil	Calibrated
<i>θ_{air}</i>	Volumetric air content in soil pores	-	m ³ air·m ⁻³ soil	
<i>φ_{air}</i>	Air filled pore space	-	m ³ air·m ⁻³ pres	
<i>τ</i>	Soil pore tortuosity	-	-	
<i>b</i>	Brooks-Corey soil retention curve exponent	5.05	-	Calibrated
<i>z</i>	Measurement height	32	m	41
<i>h_v</i>	Vegetation height	18	m	41

Table S6. The main fluxes, states and parameters used by the coupled stomatal conductance-photosynthesis module of the LSM.

Name	Description	Value	Units	Sources
<u>Subscript</u>				
<i>l</i>	Plant canopy	-		
<i>sl</i>	Sunlit big leaf	-		
<i>sh</i>	Shaded big leaf	-		
<i>k</i>	Sunlit or shaded big leaf	-		
<i>i</i>	Inside the stomatal cavity of the leaf	-		
<i>s</i>	On the leaf surface	-		
<i>g</i>	Ground/soil	-		
<i>ca</i>	Canopy airspace	-		
<u>Fluxes and States</u>				
A_n	Net CO ₂ assimilation rate	-	mol CO ₂ ·m ⁻² ·s ⁻¹	
A_c	Rubisco-limited CO ₂ assimilation rate	-	mol CO ₂ ·m ⁻² ·s ⁻¹	
A_j	Light-limited CO ₂ assimilation rate	-	mol CO ₂ ·m ⁻² ·s ⁻¹	
A_p	Product-limited CO ₂ assimilation rate	-	mol CO ₂ ·m ⁻² ·s ⁻¹	
A	CO ₂ assimilation rate	-	mol CO ₂ ·m ⁻² ·s ⁻¹	
c	CO ₂ partial pressure		Pa	
<u>Forcing Data</u>				
P_{atm}	Atmospheric Pressure	-	Pa	41
<u>Parameters</u>				
g_s	Stomatal conductance	-	mol H ₂ O·m ⁻² ·s ⁻¹ or m·s ⁻¹	
g_1	Medlyn slope parameter	0.88	kPa ^{0.5}	Calibrated
g_o	Minimal stomatal conductance	10e-4	mol H ₂ O·m ⁻² ·s ⁻¹ or m·s ⁻¹	16
g_{bv}	Leaf laminar boundary layer water vapor	-	m·s ⁻¹	
V_{max25}	Max Rubisco assimilation rate at 25°C	122	mol CO ₂ ·m ⁻² ·s ⁻¹	Calibrated
J_{max25}	Max assimilation rate based on electron transport at 25°C	256	mol photons·m ⁻² ·s ⁻¹	2.1· V_{max25}
Γ	CO ₂ compensation point	-	mol CO ₂ ·m ⁻² ·s ⁻¹	
o_i	O ₂ partial pressure	-	Pa	
K_c	Rubisco Michaelis-Menten rate constant for carboxylation	-	Pa	
K_o	Rubisco Michaelis-Menten rate constant for oxidation	-	Pa	
K_n	Nitrogen extinction coefficient	0.7	-	16
I_{PSII}	Efficiency of photosystem II to deliver electrons	-	mol photons ·m ⁻² ·s ⁻¹	
Φ_{PSII}	Quantum efficiency of photosystem II	0.7	-	16
Θ_{PSII}	Curvature factor J_{max25} and I_{PSII} co-limitation	0.85	-	16
Θ_{c_j}	Curvature factor A_c and A_j co-limitation	0.98	-	16
A_i	CO ₂ assimilation rate co-limited by A_c and A_j	-	mol CO ₂ ·m ⁻² ·s ⁻¹	
Θ_{ip}	Curvature factor A_i and A_p co-limitation	0.95	-	16
R_d	Dark respiration rate		mol CO ₂ ·m ⁻² ·s ⁻¹	

Table S7. The main fluxes, states and parameters used by the transpiration downregulation module of the LSM.

Name	Description	Value	Units	Sources
<u>Subscript</u>				
<i>l</i>	Plant canopy	-		
<i>sl</i>	Sunlit big leaf	-		
<i>sh</i>	Shaded big leaf	-		
<i>k</i>	Sunlit or shaded big leaf	-		
<i>i</i>	Inside the stomatal cavity of the leaf	-		
<i>s</i>	On the leaf surface	-		
<i>g</i>	Ground/soil	-		
<i>ca</i>	Canopy airspace	-		
<i>sx</i>	Soil-to-xylem	-		
<i>xl</i>	Xylem-to-leaf	-		
<i>la</i>	Leaf-to-atmosphere	-		
<i>ww</i>	Well-watered rate	-		
<i>max</i>	Maximum value	-		
<u>Fluxes and States</u>				
<i>LE</i>	Latent heat flux	-	mol CO ₂ ·m ⁻² ·s ⁻¹	
<i>ψ</i>	Water potential	-	MPa	
<i>e</i>	Water vapor pressure	-	Pa	
<u>Forcing Data</u>				
<i>θ_s</i>	Soil water content at 50 cm depth	-	m ³ water·m ⁻³ soil	41
<u>Parameters</u>				
<i>g_s</i>	Stomatal conductance	-	mol H ₂ O·m ⁻² ·s ⁻¹ or m·s ⁻¹	
<i>g</i>	Segment-specific conductance	-	m·s ⁻¹ ·MPa ⁻¹	
<i>ψ_{s,sat}</i>	Saturated soil water potential	-1e-3	MPa	Calibrated
<i>b</i>	Brooks-Corey soil retention curve exponent	5.05	-	Calibrated
<i>c</i>	Brooks-Corey hydraulic conductivity exponent	-	-	
<i>d</i>	Adjusting factor for roots in soil conductance	4	-	37
<i>K_{s,sat}</i>	Saturated soil hydraulic conductivity	10	m·d ⁻¹	Calibrated
<i>RAI</i>	Root area index	11	m ² root area·m ⁻² ground area	11
<i>d_r</i>	Fine root diameter	5e-04	m	11
<i>Z_r</i>	Effective rooting depth	0.1	m	42
<i>ψ_{x,50}</i>	Xylem water potential at 50% loss of conductance	-9.9	MPa	Calibrated
<i>a</i>	Xylem vulnerability curve shape parameter	0.3	-	Calibrated
<i>K_{sap}</i>	Sapwood hydraulic conductivity	9.30E-04	kg·m ⁻¹ ·s ⁻¹ ·MPa ⁻¹	Calibrated
<i>h_v</i>	Vegetation height	18	m	41
<i>ψ_{l,50}</i>	Leaf water potential at 50% loss of conductance	-9.9	MPa	Calibrated
<i>b_l</i>	Leaf vulnerability curve shape parameter	0.3	-	Calibrated
<i>LAI</i>	Leaf area index	3.2	m ² leaf area·m ⁻² ground area	Calibrated
<i>Φ</i>	Flux potential from Kirchoff transform	-	kg·s ⁻¹	

Table S8. The main physical constants used in the LSM.

Name	Description	Value	Units
ρ_w	Density of water	1000	$\text{kg}\cdot\text{m}^{-3}$
ρ_a	Density of air	1.2	$\text{kg}\cdot\text{m}^{-3}$
ε	Molar ratio of water to air	0.622	-
\mathcal{L}_v	Latent heat of vaporization	2.50E+06	$\text{J}\cdot\text{kg}^{-1}$
k	von Karmen constant	0.4	-
ν	Kinematic viscosity	1.50E-05	$\text{m}^2\cdot\text{s}^{-1}$
R_g	Universal gas constant	8314	$\text{J}\cdot\text{K}^{-1}\cdot\text{mol}^{-1}$
c_p	Specific heat of air at constant pressure	1004	$\text{J}\cdot\text{kg}^{-1}\cdot\text{K}^{-1}$

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