

Turbulence in the stratified boundary layer under ice: observations from Lake Baikal and a new similarity model

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Abstract.

Seasonal ice cover on lakes and polar seas creates seasonally developing boundary layer at the ice base with specific features: fixed temperature at the solid boundary and stable density stratification beneath. Turbulent transport in the boundary layer determines the ice growth and melting conditions at the ice-water interface, especially in large lakes and marginal seas, where large-scale water circulation can produce highly variable mixing conditions. Since the boundary mixing under ice is difficult to measure, existing models of ice cover dynamics usually neglect or parameterize it in a very simplistic form. We present first detailed observations on mixing under ice of Lake Baikal, obtained with the help of advanced acoustic methods. The dissipation rate of the turbulent kinetic energy (TKE) was derived from correlations (structure functions) of current velocities within the boundary layer. The range of the dissipation rate variability covered 2 orders of magnitude, demonstrating strongly turbulent conditions. Intensity of mixing was closely connected to the mean speeds of the large-scale under-ice currents, ~~the latter being of geostrophic origin and having lake-wide scales~~. Mixing developed on the background of stable density (temperature) stratification, which affected the vertical structure of the boundary layer. To account for stratification effects, we propose a model of the turbulent energy budget based on the length scale incorporating the dissipation rate and the buoyancy frequency (Dougherty-Ozmidov scaling). The model agrees well with the observations and yields a scaling relationship for the ice-water heat flux as a function of the shear velocity squared. The ice-water heat fluxes in the field were the largest among all reported in lakes (up to 40 W m^{-2}) and scaled well against the proposed relationship. The ultimate result consists in a strong dependence of the water-ice heat flux on the shear velocity under ice. The result suggests large errors in the heat flux estimations, when the traditional “bulk” approach is applied to stratified boundary layers. It also implies that under-ice currents may have much stronger effect on the ice melt than estimated by traditional models.

1 Introduction

The demand on a better quantitative description of the formation, evolution, and decay of the seasonal ice has grown recently because of large-scale trends to shortening of the ice season in the Northern Hemisphere and the drastic decrease of the arctic sea ice extent. Closure of the global mass budget of the arctic seasonal ice is a complex problem, related, apart from the atmospheric and terrestrial heat sources, to the upward transport of heat stored in the under-ice water body. An important role in the heat budget of seasonal ice is played by the storage of the solar radiation in the under-ice water, which is subsequently transported to the ice base by the under-ice currents. The effect of currents on ice melt is particularly apparent in the Arctic ocean, where the loss of ice mass in spring and summer occurs mainly from the ice bottom (McPhee, 1992; Perovich et al., 2011; Carmack et al., 2015; Peterson et al., 2017). Apart from the polar oceans and seas, seasonal formation of ice cover is an essential feature of high-latitude freshwater lakes. Physics of seasonal ice cover on lakes has gained particular attention, as an essential part of climate change research (Magnuson et al., 2000; Kirillin et al., 2012). A shorter seasonal ice cover as a result of global warming may produce a positive feedback due to increase of greenhouse gas emission and changing the global carbon budget (Tranvik et al., 2009). Hence, estimation of the consequences of phenological changes on inland waters requires quantification of the physical mechanisms that control the formation and melting of ice. The heat and mass transfer at the ice-water interface is the least studied among these mechanisms (Kirillin et al., 2012; Aslamov et al., 2014a).

The seasonal ice cover on lakes, especially, on large ones, shares many basic features with the seasonal sea ice. Storage of the heat from solar radiation in the surface mixed layer (SML) and its subsequent release to the ice base is the major mechanism of the ice cover melt in lakes (Kirillin et al., 2012) as well as in the ocean (Perovich and Richter-Menge, 2009). However, in contrast to the seawater, lakes own some specific physical features affecting formation and melting of ice. Water temperatures in ice-covered freshwater lakes are below their value of maximum density. Therefore, solar heating of upper layers produces free convection, which is the major mechanism of the SML formation (Mironov et al., 2002). In addition to the storage of the heat from the short-wave radiation penetrating the ice, convective mixing in the SML entrains the warmer water from the deep layers (Kirillin et al., 2012). The convective SML is separated from the ice base by a stably stratified interfacial layer (IL) with the upward temperature drop down to the freezing point of freshwater. At low salinities, water temperatures remain always higher than that of ice; hence the IL with downward temperature increase always exists under the ice base; its thickness depends on the strength of radiative heating and the temperature of the SML underneath (Barnes and Hobbie, 1960). The strong stratification in the IL prevents convective mixing despite the negative buoyancy production by the decrease of the solar radiation with depth and reduces convective transport of heat to the ice-water interface. As a result, only a small amount of the heat is available for ice melt, despite strong convection in the SML (Kirillin et al., 2018b). The situation is akin to formation of a stably stratified layer beneath the ice base and the near-surface temperature maximum in marginal polar seas, driven by freshening of the surface waters due to ~~river runoff or the accelerated~~ runoff or accelerated sea ice melt (Jackson et al., 2010, 2012).

The ice-water interaction becomes more complex when a freshwater lake becomes essentially large compared to the Rossby radius of deformation (Gill, 1982). The ~~latter~~ condition suggests long-lasting water circulation under ice, which, similarly to the

ocean circulation, is able to produce significant velocity shear at the ice base and accelerate by this the upward heat transport (McPhee, 1992). Among such lakes, Lake Baikal—the largest lake by volume on earth—most closely resembles the Arctic Ocean with regard to the seasonal ice dynamics. Thanks to the strong winter cooling under influence of the Siberian atmospheric pressure maximum, Lake Baikal reveals a steady ice cover over the entire lake for 3-5 months of the year. Consequently, the seasonal ice regime plays a crucial role in hydrodynamics and ecosystem functioning of the lake. ~~Aslamov et al. (2014a) reported high heat fluxes from water to ice of~~ Aslamov et al. (2014a, b) reported high water-to-ice heat fluxes in Lake Baikal during the period of ice growth. The fluxes were apparently related to the surface water circulation pattern beneath the ice cover. ~~Despite the thin snow cover, the convection due to penetrating solar radiation was not able to produce such a strong boundary mixing, exceeding the fluxes~~ Zhdanov et al. (2017). The observed ice to water heat fluxes exceeded fluxes measured in small lakes (Kirillin et al., 2018b) up to an order of magnitude. Free convection due to penetrating solar radiation was not strong enough to produce the upward heat release at these rates. Hence, the high turbulence level was tentatively referred by the authors to the shear mixing produced by the water circulation under the ice surface.

The intensity of turbulence produced by velocity shear in the boundary layer and the resulting heat transport from water to the ice base may vary depending on the current velocity, ice structure, and density stratification under ice. In order to estimate the effect of under-ice circulation on the ice-water heat flux in lakes, we performed a field experiment combining temperature measurements with high temporal and vertical resolution within the ice cover and fine-scale registration of ~~current~~ currents velocities under the ice base. The temperature observations were subsequently used for estimation of the heat budget at the ice-water interface and derivation of the ice-water heat fluxes. The data on fine-scale velocity fluctuations provided information on variability of mean currents under ice as well as on the characteristics of turbulent mixing in the ice-water boundary layer in form of the dissipation rate of the turbulent kinetic energy (TKE). Below, both outcomes of the field experiment are combined to ~~reveal~~ analyze the characteristics of the turbulent boundary layer and to analyze the effect produced by turbulent mixing on the ice cover thickness. The overarching goal of the presented study ~~consists in establishing is to establish~~ the scaling relationships linking the under-ice circulation ~~and with~~ the seasonal ice cover dynamics and suitable for parameterization of the ice-water heat exchange in regional and global models of seasonal ice.

2 Heat budget of seasonal ice cover and scaling of the under-ice boundary layer

To a good approximation, the base of the lake ice can be represented as a rigid boundary on top of a fluid, i.e. the vertical heat transport at the ice-water interface is close to purely conductive on both ice and water sides, governed by molecular forces within the ice cover and within a thin “conduction” layer of water. It should be noted that the assumption holds generally true for a solid freshwater ice with low amount of impurities: Salt-water ice undergoes brine extraction, which can induce convection by mass flux at the boundary and ~~increase remarkably~~ markedly increase the heat transport. ~~Similar~~ Similarly, the increase of water flow ~~and destruction of~~ can destroy the conduction layer at the ice-water interface ~~may take place~~ in a “rotten” freshwater ice subject to internal melting, especially in presence of impurities (Bluteau et al., 2017). In the majority of freshwater lakes, the aforementioned effects are negligible during the most of the ice season. In particular, in Lake Baikal,

due to cold winters and low snow precipitation, practically 100% of the ice cover consists of clear congelation (“black”) ice, which grows at the ice-water interface, has homogeneous crystal structure and much lower amount of impurities than ~~the~~-sea ice or ~~the~~-river ice (Kirillin et al., 2012). Hence, the heat balance at the ice-water interface (IWI) can be expressed as the sum of conductive (molecular) boundary fluxes ~~Q_{cw} and Q_{ci}~~ and the heat release/consumption due to the phase change (freezing or melting) (Aslamov et al., 2014a):

$$Q_{\underline{cw}i\underline{w}} = Q_{ci} - \rho_i L_f \frac{dh_i}{dt} \quad (1)$$

where the vertical coordinate is directed downwards with the origin z_i at the IWI, $dh_i dt^{-1}$ is the rate of basal ice melting (growth); $\rho_i L_f$ is the product of the ice density and latent heat of fusion, ~~Q_{cw}~~ $Q_{i\underline{w}}$ is the conductive heat flux from/to the water,

$$10 \quad Q_{\underline{cw}i\underline{w}} = C_{pw} \rho_w \kappa_w \left. \frac{dT}{dz} \right|_{z_i+0}, \quad (2)$$

and Q_{ci} is the conductive heat flux from/to the ice,

$$Q_{ci} = C_{pi} \rho_i \kappa_i \left. \frac{dT}{dz} \right|_{z_i-0}. \quad (3)$$

Temperature at the IWI is fixed at the melting point of 0 °C that corresponds to the following thermodynamic characteristics: molecular heat diffusion coefficient for fresh water $\kappa_w \approx 1.4 \cdot 10^{-7} \text{ m}^2 \text{ s}^{-1}$, molecular heat diffusion coefficient for ice $\kappa_i \approx$
 15 $1.1 \cdot 10^{-6} \text{ m}^2 \text{ s}^{-1}$; the product of the water heat capacity and density is $C_{pw} \rho_w \approx 4.18 \cdot 10^6 \text{ J K}^{-1} \text{ m}^{-3}$, and the same product for ice is $C_{pi} \rho_i \approx 1.96 \cdot 10^6 \text{ J K}^{-1} \text{ m}^{-3}$ (see e.g. Leppäranta, 1983).

Equation (1) can be applied for reliable estimation of the ice-water heat flux Q_{wi} if the temperature profile within the ice cover and the time variations of the ice thickness $dh_i dt^{-1}$ are known. This approach was used for estimation of the heat fluxes in Lake Baikal by Aslamov et al. (2014a), who recorded the temperature profile within the ice cover and the variations of the
 20 ice thickness $dh_i dt^{-1}$ with high temporal resolution.

However, direct estimation of ~~Q_{cw}~~ $Q_{i\underline{w}}$ in the absence of detailed data on the ice cover dynamics and temperature is less straightforward. The bulk of the water column under the ice is turbulent: While ice-covered waters are isolated from the direct influence of wind, vertical heat transport remains higher than the purely molecular one, intensified by convective mixing due to solar radiation penetrating the ice cover and due to shear turbulence produced by under-ice currents. As a result, the thickness of
 25 the “diffusive” layer in the immediate vicinity of the ice base, where Eq. (2) holds true, does ~~typically not~~ not typically exceed several millimetres. The temperature gradient $dT dz^{-1}$ at $z = 0$ is barely detectable by the traditional observation methods and is varying continuously depending on the mixing and temperature conditions in the underlying water column.

The turbulent kinetic energy (TKE) under ice is supplied by the ~~decay of the convective motions~~ free convection in the underlying convectively mixed layer (Mironov et al., 2002) and/or by mean horizontal ~~circulation~~ current (McPhee, 1992;
 30 Aslamov et al., 2014a). In the latter case, the vertical turbulent transport of momentum $\tau = \langle u'w' \rangle = u_*^2$ is created by the current velocity shear $S = \partial U_{\text{mean}} \partial z^{-1}$ at the ice base. Hence, close to the IWI, the distance from the ice base z is a major

parameter, determining the turbulent mixing characteristics. Assumption of the proportionality between the mixing length scale and the distance from the solid boundary leads to the relationships for the neutral (logarithmic) boundary layer,

$$S = \frac{\partial U_{\text{mean}}}{\partial z} = \frac{u_*}{\kappa z}, \quad (4)$$

or

$$5 \quad U_{\text{mean}}(z) = \frac{u_*}{\kappa} \ln\left(\frac{z}{z_0}\right), \quad (5)$$

where z is the turbulent pulsation length scale equal to the distance from the lower boundary of the ice, z_0 is the roughness ~~parameter of length at~~ the ice bottom, $\kappa \approx 0.4$ is the empirically determined von Kármán constant, u_*^2 is the turbulent stress (shear velocity squared) produced by the vertical shear of the mean velocity S .

The logarithmic velocity distribution in the ice cover vicinity makes possible estimation of the momentum flux based on
10 mean velocity values only using a direct relationship derived from Eq. (5) (“bulk” formula),

$$u_*^2 = C_Z U_z^2, \quad (6)$$

where the bulk transfer ~~or drag~~ coefficient C_Z corresponds to the depth Z of the current speed measurements, and is defined as

$$C_Z = \left(\frac{\kappa}{\log z - \log z_0} \frac{\kappa}{\ln z - \ln z_0} \right)^2.$$

For a non-stratified steady-state turbulent boundary layer, the TKE budget tends to ~~be~~ the local balance of the largest terms
15 in the TKE transport equation, production and dissipation:

$$\varepsilon = u_*^2 S \propto \frac{u_*^3}{z}, \quad (7)$$

where ε is the TKE dissipation rate.

The second factor influencing the buoyancy flux at the ~~IWI—the IWI is the~~ destabilizing buoyancy flux B_R due to volumetric
20 absorption of solar radiation $I(z)$ within the convectively mixed water column of thickness h_S ~~—is—~~. ~~The buoyancy flux is~~
derived from the heat transport equation with radiation term as

$$B_R = \beta \left(I(0) + I(h_S) - 2h_S^{-1} \int_0^{h_S} I(z) dz \right). \quad (8)$$

Here, the assumption of height-constant warming rate within the convective layer was used (Mironov et al., 2002). $\beta = g\alpha_T(T)$
is the buoyancy parameter, α_T is the thermal expansion coefficient. The latter is generally not constant in freshwater due
to non-linearity of the equation of state at temperatures close to the maximum density value $T_{md} \approx 3.98$ ~~$T_{md} \approx 3.98$~~ °C
25 [$\alpha_T(T) \approx 0.825 \cdot 10^{-5} (T - T_{md}) \alpha_T(T) \approx 0.825 \cdot 10^{-5} (T - T_{md})$ K⁻¹, see e.g. Farmer and Carmack, 1981].

The ice-water boundary layer in freshwater lakes is rarely neutrally stratified: A distinctive feature of the layer is the fixed
temperature at the IWI. As a result, the water adjacent to the IWI in fresh or brackish environments is always subject to stable

stratification, with deeper waters being warmer and thereby heavier. Stratification may alter the turbulent length scale, affecting Eqs. (4)-(7). Stratification counteracts the shear production of turbulence, and in the asymptotic case of a strongly stratified layer, is the sole mechanism of turbulence damping. This effect can be accounted for by introduction an additional length scale apart from the distance to the boundary z , as expressed by a simple formula following from dimensional analysis,

$$5 \quad S = \frac{u_*}{\kappa z} \left(1 + C_x \frac{z}{L_x} \right), \quad (9)$$

where L_x is the stratification length scale and C_x is an empirical coefficient. Eq. (9) replaces Eq. (4) in stratified conditions with corresponding changes in Eqs. (5)-(7). For conditions dominated by the stabilizing buoyancy flux at the boundary $B_s = \beta Q_{cw} (c_{pw} \rho_w)^{-1} B_s = \beta Q_{iw} (c_{pw} \rho_w)^{-1}$, the stratification length scale turns to the well-known Monin-Obukhov length scale $L_{MO} = u_*^3 B_s^{-1}$, with empirically determined coefficient $C_x \approx 5$ (Stull, 2012), building the basis for the Monin-Obukhov

10 similarity theory (MOST).

If stratification is created outside the boundary layer, its effect on boundary mixing is independent on the surface buoyancy flux. A characteristic length scale for turbulence in stratified media was independently proposed by Dougherty (1961) and Ozmidov (1965), as

$$L_N = \varepsilon^{1/2} N^{-3/2}, \quad (10)$$

15 where

$$N = \left(\frac{g}{\rho} \frac{\partial \rho}{\partial z} \right)^{1/2} \approx \left(\beta \frac{\partial T}{\partial z} \right)^{1/2}$$

is the buoyancy frequency, ρ is water density in assumption of negligible salinity effects. Replacement of L_x by L_N in Eq. (9) yields in this case

$$S = \frac{u_*}{\kappa z} \left(1 + C_N \frac{z N^{3/2}}{\varepsilon^{1/2}} \right), \quad (11)$$

20 with the corresponding expression for the TKE production rate P (cf. Eq. 7)

$$P = u_*^2 S = \frac{u_*^3}{\kappa} \left(\frac{1}{z} + \frac{C_N}{L_N} \right). \quad (12)$$

Close to the boundary, $z \ll L_N$, Eq. (11) approaches the neutral scaling relationship (4). At large distances from the boundary, $z \gg L_N$, Eq. (11) turns to a “z-less” scaling

$$S = C_N \frac{u_*}{\kappa L_N}, \quad (13)$$

25 which in turn yields the N -scaling for P

$$P \propto u_*^3 L_N^{-1}.$$

In stably stratified conditions, production of TKE is ~~balanced~~ counteracted by two major loss processes, dissipation ε and work against the stability forces B_{St} . The latter can be expressed in form $B_{St} = K_\rho N^2$, where K_ρ is the diapycnal diffusivity. From

analysis of dimensions, the turbulent diffusivity can be scaled as $K_\rho \propto u_*^2 N^{-1}$ (see e.g. Monin and Ozmidov, 1981). Then, the TKE budget can be approximated as

$$\frac{u_*^3}{\kappa} \left(\frac{1}{z} + \frac{C_N}{L_N} \right) = C_B u_*^2 N + \underline{\epsilon} \epsilon, \quad (14)$$

with coefficients C_N, C_B subject to estimation from empirical data, or

$$5 \quad u_*^2 N \left(Ri^{-1/2} - C_B \right) = \epsilon, \quad (15)$$

where Ri is the gradient Richardson number,

$$Ri = \frac{N^2}{S^2}, \quad (16)$$

expressing the relative importance of stratification and velocity shear for the vertical transport. Its critical value $Ri_{cr} \approx 1/4$ (Turner, 1979) marks the boundary between turbulent conditions, at which the shear can destroy the stratification and the quiet conditions, at which strong N ultimately suppresses any turbulent motions. Hence, for turbulence to exist at weakly supercritical Ri , it is required $0 < C_B < 2$ in Eq. (15). In the following we tentatively assume $C_B \approx 1$. Another scaling relationship relevant to the turbulent mixing on the background of stable stratification is the buoyancy Reynolds number,

$$Re_b = \frac{\epsilon}{\nu N^2}, \quad (17)$$

where $\nu = \mathcal{O}(10^{-6}) \text{ m}^2 \text{ s}^{-1}$ is the water viscosity. Re_b refers to the work of turbulence against stratification and viscosity, which becomes important at distances from the solid boundary shorter than characteristic length scales of turbulence; its critical value is reported to be $\mathcal{O}(10^1)$ (Gargett et al., 1984).

In neutral conditions, the coefficient of turbulent heat transfer $K_Z \propto u_* z$ (assuming the turbulent Prandtl number is approximately 1), and the corresponding relationship for the ice-water heat flux Q_{iw} (Eq. 42) can be written as:

$$Q_{iw} \propto u_* \Delta T \propto \frac{u_* h_T}{g \alpha_T} N^2, \quad (18)$$

where ΔT is the temperature difference across the layer h_T beneath the ice base, often assumed in models as water temperature at the vertical grid point closest to the ice. The expression is sometimes used in form of the ‘‘bulk’’ formulation, assuming direct relationship between the friction velocity and the main current speed $u_* \propto U \propto u_* \propto U_{mean}$ (cf. Eq. 6):

$$Q_{cw iw} = C_Q U_{mean} \Delta T, \quad (19)$$

where ΔT is the temperature difference across h_T , and C_Q is an empirical ‘‘bulk’’ heat transfer coefficient. The values of C_Q were reported in the range $[0.8 \pm 0.3] \cdot 10^{-3}$ (Hamblin and Carmack, 1990; Nan et al., 2016); stratification effects on C_Q were not investigated.

Adopting the same scaling considerations as in Eqs. (14)-(15), the heat flux at the IWI Q_{iw} in strongly stratified conditions may be assumed to depend on the work of turbulence against the stability,

$$Q_{iw} = K_\rho \frac{\partial T}{\partial z} \propto u_*^2 N^{-1} \Delta T h_T^{-1},$$

or, in terms of buoyancy flux B_{iw} ,

$$B_{iw} = g\alpha_T Q_{iw} \propto u_*^2 N. \quad (20)$$

Herewith, a strongly stratified case is characterized by the flux dependence on the shear velocity squared and a weaker dependence on the stratification, expressed by 1/2 exponent at the vertical density gradient (as revealed by the direct proportionality to the buoyancy frequency N).

Summarizing the considerations above, validation of the Dougherty-Ozmidov scaling (Eqs. 10-14) for the ice boundary layer, and the ice-water flux parameterization (Eqs. 18-20) is possible when field data are available on both the TKE dissipation rates and the mean fields of governing forces (production of convective instability by radiative heating and/or mean horizontal flow due to under-ice currents).

3 Study site and field methods

The field study was performed in February-March 2017 in the southern part of Lake Baikal. Two custom-made autonomous stations were installed in the vicinity of a quasi-stationary longshore current ~~-, which has~~ (Fig. 1). Under-ice currents with characteristic velocity and spatial scales of the observed under-ice currents are (10^{-2} - 10^{-1} m s $^{-1}$) and (10^4 - 10^5 m), respectively, have been regularly observed in this region during ice cover period (Fig. 1)(Aslamov et al., 2014a, 2017). The scales of the flow suggest that away from boundaries it is balanced mainly by the Coriolos force, while the forcing may be refered to density gradients created by in- and outflows, and topographic effects.

Station S1 was installed 4.5 km away from the lake shore ($51^\circ 51.923'N$, $105^\circ 4.779'E$) in the area of quasi-stationary jet-like alongshore current. Station S2 was located 3.5 km to the south from station S1. The total water depth in the vicinity of both stations amounted at ~~≈ 1600~~ ≈ 1400 m. Each station registered temperature at 30 vertical levels distributed within the ice cover, the water boundary layer, and the air above the ice. The distance between temperature sensors was 5 cm within the ice and in the under-ice boundary layer, increasing up to 10-50 cm at larger distances from the ice boundary in the water and in the air. Three ~~sensors of the~~ short-wave solar radiation ~~registered the vertical radiation decay within~~ sensors were deployed vertically to measure the decay of solar radiative fluxes across the air-ice-water system. Ice thickness was measured by a 330-kHz echo-sounder, deployed upward-looking at a fixed distance from the ice surface. The resolution of the system was 0.002 °C for temperature, 0.1 W m $^{-2}$ for solar radiation, and 0.1 mm for ice thickness (the operation range of 0.2–2.8 m). The system collected data with a period of 2 min, logging them internally, and sending data several times a day via cellular network to a remote Internet-server (see Aslamov et al., 2017, for detailed description of the ice station configuration). Two-dimensional electromagnetic current meters “INFINITY-EM” (JFE Advantech Co., Ltd.) were used to measure the current velocities: (velocity range ± 5 m s $^{-1}$, resolution 0.02 cm s $^{-1}$, accuracy ± 1 cm s $^{-1}$). The current meters were positioned at a distance of 1 m from the surface of the ice cover. Three additional current meters were deployed at Station S1 at distances 0.6, 0.8 and 1.4 m from the ice surface.

Characteristics of turbulent mixing in the under-ice boundary layer were obtained with the help of the high-resolution Doppler current velocity profiler HR Aquadopp (Nortek AS, Norway). The profiler was deployed for 48 hours successively at

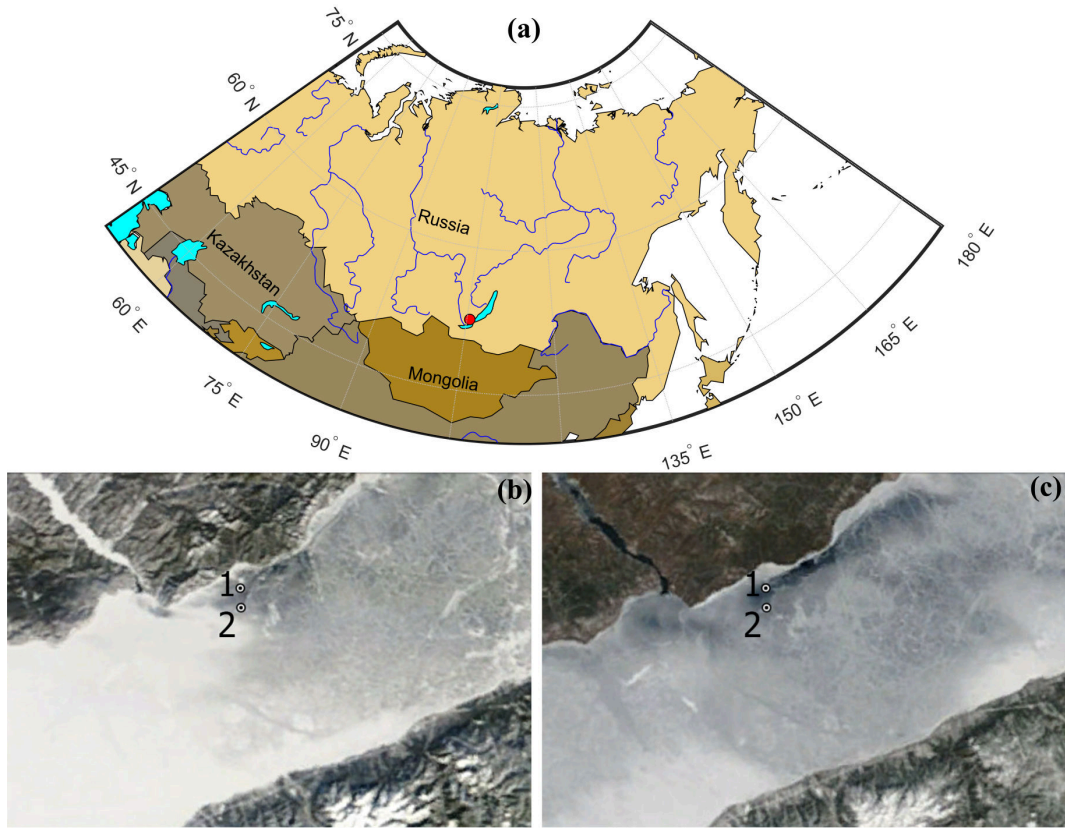


Figure 1. (a) [Geographical location of the study site](#). Ice conditions in Southern Baikal on (b) 9 February and (c) 12 April 2017 (b) and locations of the autonomous measurement stations. The satellite images are from the Irkutsk Center of Remote Sensing (sputnik.irk.ru, 2019). Note the stronger ice melt in the area of the jet current around Station S1, visible as a dark area in Panel [bc](#).

each of the two stations, on 05-07 March 2017 on Station S1, on 08-10 March on Station S2. The profiler was frozen into the ice downlooking, with the acoustic head at 2 cm beneath the ice base (verified with an ROV camera). Three components of current velocity were registered with a time interval of 2 s and a spatial resolution of 15 mm in the pulse-to-pulse coherence (high-resolution) mode.

- 5 The values of short-period fluctuations of the flow velocity were used to calculate dissipation rates of the kinetic energy of turbulence (TKE) based on the Kolmogorov's 1941 hypothesis on the self-similarity of the velocity structure functions using the method described by Wiles et al. (2006). ε was derived as a coefficient in the semi-empirical equation for the velocity structure function $D_i(r)$ along the i -th acoustic beam,

$$D_i(r) = \text{Noise} + C_v \varepsilon^{2/3} r^{2/3}, \quad (21)$$

- 10 which includes noise estimation Noise due to instrumental noise and non-turbulent velocity fluctuations. Here, the constant $C_v = 3^{1/3}$ (see e.g. Lien and D'Asaro, 2002). The velocity structure function was calculated from the measured along-beam

velocities $v_i(z)$ at the distance z from the instrument's head as

$$D_i(r) = \langle (v_i(z) - v_i(z+r))^2 \rangle \quad (22)$$

Quality check was performed based on values of Noise in Eq. (21); the ε values from three beams were compared for similarity and averaged. The detailed procedure of data post-processing and quality check is described in (Kirillin et al., 2018a; Volkov et al., 2018).

4 Results

4.1 Atmospheric conditions and ice cover thickness

The ice cover formed on Lake Baikal during the second half of January 2017 with several periods of ice break-up and re-freeze. The autonomous stations were installed on 1 Feb 2017 and provided background information on the major forces driving the ice cover development. The temperatures of the ice surface remained below the freezing point of water during the entire observations period, varying between -14° C and -2° C with a slight warming trend (Fig. 2a). The initial ice thicknesses were nearly the same at both stations: 23 cm at Station S1 and 24 cm (a day later) at Station S2. During the first 2 weeks of February the ice grew at a nearly constant rate of 1.2-1.3 cm day⁻¹ (Fig. 2b). During this period, the ice surface temperatures at both stations were nearly equal and followed closely the air temperatures at the 1.5 m height. This quasi-neutral stratification in the air-ice boundary layer lasted until the end of February, caused apparently by convective heat flux from the ice surface due to release of the latent heat of ice formation. Later, the ice thickness at Station S1 (the one with strong under-ice currents) remained nearly constant, while basal ice at Station S2 continued to grow at a slow rate of ≈ 0.3 cm day⁻¹ ~~at Station S2~~ (Fig. 2**eb**). In mid-March, a stable stratification developed in the air above the ice with air temperatures dropping down to -16° C. Whereas the temperature at the ice surface of Station S2 also decreased following the air temperature trend, the ice surface at Station S1 remained relatively warm, suggesting, together with the nearly constant ice thickness, a balance between the heat release to the atmosphere and the heat supply from the water column.

4.2 Mean currents, temperatures and stratification

The thermal structure was similar at both stations S1 and S2 (Fig. 3): Under ice, the water temperatures slightly increased with depth ~~(Fig. 3): the~~. The mean vertical gradient of 0.6° C over the upper 10 m of the water column was about an order of magnitude weaker than those typically observed in shallow ice-covered lakes (Kirillin et al., 2018a). Below 10 m depth the water column was well-mixed vertically down to 30 m. Closer to the ice base, two horizontal layers could be distinguished: a < 0.5 m thin layer adjacent to the ice with the temperature difference of $\approx 0.3^\circ$ C across it. Underneath, a layer with nearly linear temperature increase of $\approx 0.03^\circ$ C m⁻¹ spread down to the 10 m depth.

In terms of stability, the two-layered thermal structure can be described by two nearly constant buoyancy frequencies $N_\delta \approx 2 \cdot 10^{-2}$ s⁻¹ in the layer $0 \leq z < \delta$ and $N_S \approx 4 \cdot 10^{-3}$ s⁻¹ in the layer ~~$\delta \leq z < z_S$~~ $\delta \leq z < h_S$, where the thickness of the sub-ice layer $\delta \approx 0.4$ m and the lower boundary of the stratified ~~interfacial layer $z_S \approx 10$ m~~ IL $h_S \approx 10$ m.

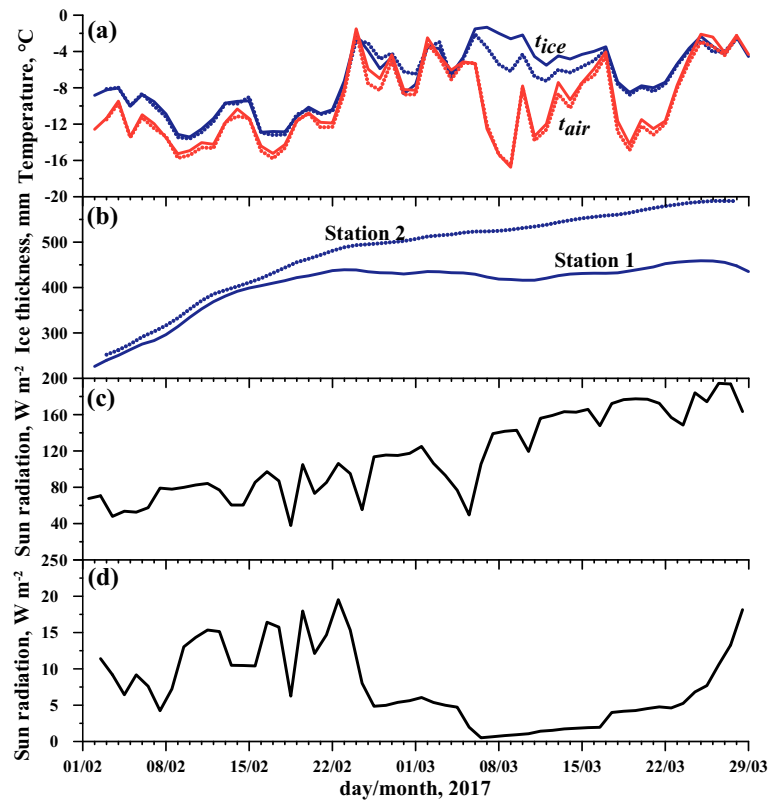


Figure 2. Background data on the Lake Baikal ice regime during the observations: (a) Daily averaged temperatures of ice surface (t_{ice}) and air temperatures at 1.5 m above the ice (t_{air}), (b) ice thickness, (c) incoming and (d) penetrated solar radiation. In Panels (a) and (b) solid lines correspond to Station S1 and dotted lines are for Station S2.

The two-layer structure is less distinct and the mixed layer temperature was slightly higher at S2 (Fig. 3b); S1 had in turn a stronger vertical gradient close to the ice base. These temperature differences between the two stations suggested a stronger upward heat transport at S1 due to stronger vertical mixing caused by water flow. Current speeds were indeed almost twice as high at S1 than at S2 (Figs. 3, 5). The currents in the upper 20 m of the water column had uniform WSW direction aligned with the shoreline (see velocity vectors in Fig. 3). A weak, 10-15° anticlockwise rotation of the current vector was detectable within the 1-2 m thin layer adjacent to the ice base, suggesting some effect of the Coriolis force on the boundary layer currents. The diurnal and synoptic variations of the ice and water temperatures were similar to those observed in the previous years (Aslamov et al., 2014a, 2017). The diurnal temperature oscillations driven by the solar radiation cycle were apparent in both water column and ice cover, with amplitudes decaying towards the ice-water interface. The begin of the melt phase after 26 Mar 2017 was indicated by homogenizing of the ice cover temperature at the melting point of 0° C. Earlier, occasional increases of the air temperature, e.g. on 25 Feb, provoked deceleration of the ice growth or short-term melting periods on both stations (Fig. 4). Relevant to the matter of the present study, a remarkable increase of the ice temperatures was observed on both stations

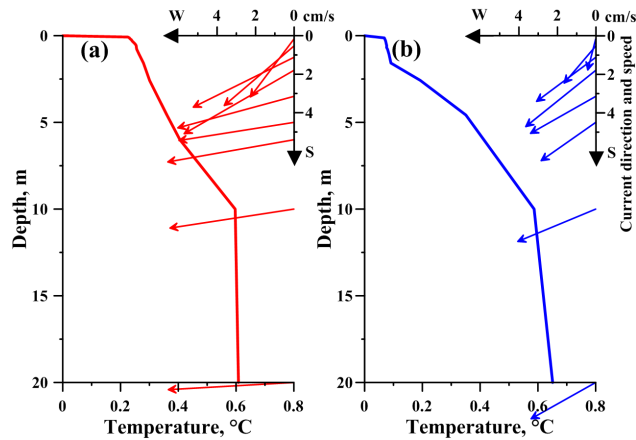


Figure 3. Temperature and currents velocity vector profiles measured from 14:00 to 16:00 06.03.2017 (local time) on (a) Station S1, (b) Station S2.

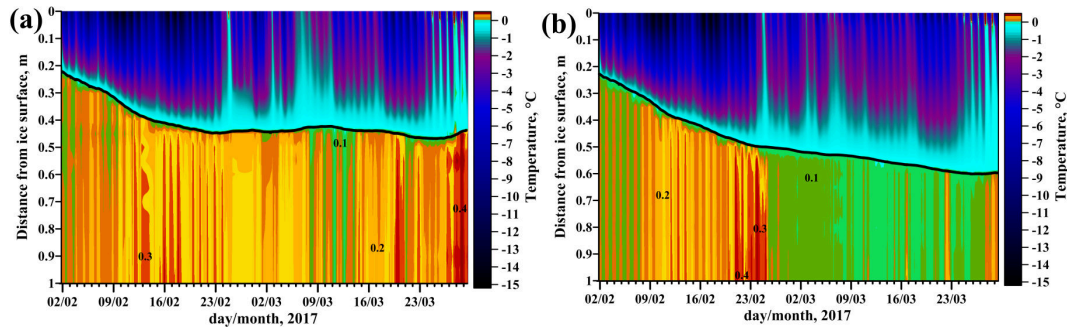


Figure 4. Temperature maps in the ice boundary layer and within the ice cover at Station S1 (a) and Station S2 (b). Note the different color scales for ice and water.

during the period of turbulence measurements on 06-12 Mar. The warming was not correlated with the air temperature: the latter dropped significantly during the same time (Fig. 2). At S1, the warming was strong enough to produce decrease of the ice thickness (Fig. 4(a)), while the effect at S2 was too weak to cause any ice melt (Fig. 4(b)).

The mean currents obtained with the acoustic Doppler profiler at a time interval of 2 s and a spatial resolution of 15 mm (Fig. 5a,b) agreed remarkably well with the records from the 5 electromagnetic velocity loggers at coarser temporal and spatial resolution (Fig. 5a,b). The result allowed later extension of the boundary layer turbulence analysis on the whole period of electromagnetic velocity measurements, after a relationship between the mean flow characteristics and the turbulent energy production was established from the short-term acoustic profiling.

The mean current velocities from the two neighboring stations demonstrated different water flow patterns. At Station S1, current velocities in the upper 1 m of the water column had mean values $\geq 5 \cdot 10^{-2} \text{ m s}^{-1}$. The magnitudes underwent variations on synoptic time scales, changing at 1 m under the ice from $\approx 10 \cdot 10^{-2} \text{ m s}^{-1}$ to $\approx 3 \cdot 10^{-2} \text{ m s}^{-1}$ within 48 hours (Fig. 5a).

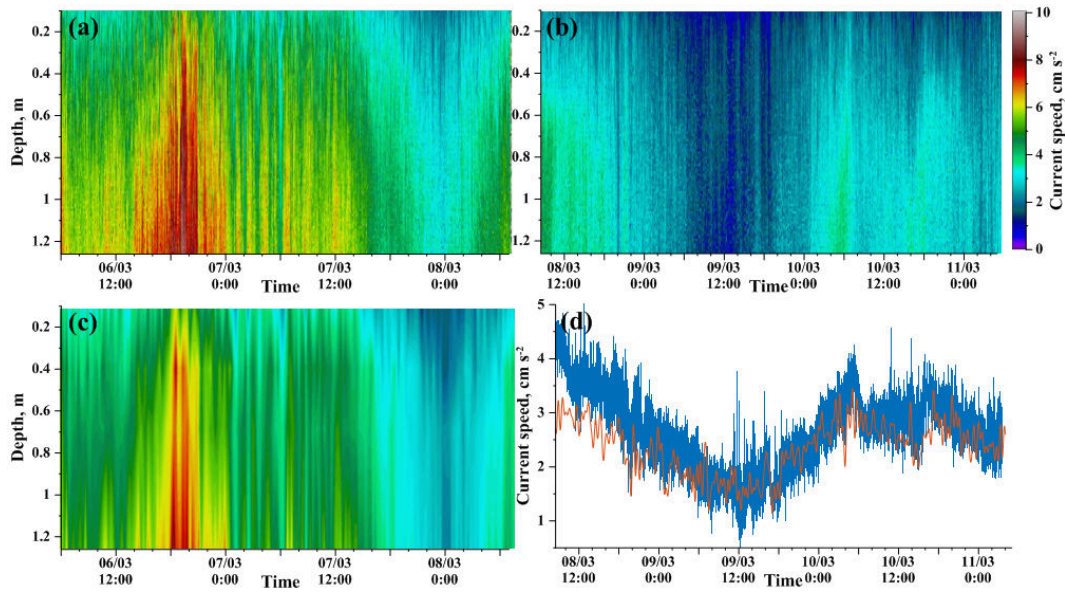


Figure 5. Horizontal current speeds at Station S1 (a, c) and Station S2 (b, d) measured by the acoustic Doppler profiler Aquadopp (a, b) and the electromagnetic loggers INFINITY (c). Panels (a-c) are the time-depth maps, Panel (d) shows the horizontal flow speed measured by a single INFINITY logger at 1 m under ice (thin red line) and the Aquadopp velocity record from the same depth (thick blue line).

The event coincided with melting of the ice cover (Fig. 4(a-a)) suggesting the upward heat transport by the currents to be the mechanism of ice heating in this case despite the low air temperatures (Fig. 2a). Farther from the lake shore, at Station S2, the currents, as measured during the next 2 days, revealed a lower variability with time, and had lower magnitudes of 1 to $4 \cdot 10^{-2} \text{ m s}^{-1}$.

5 4.3 Solar radiation

The solar radiation flux at the ice surface doubled within the 2 months of observations (Fig. 2c) contributing to the deceleration of the ice growth. The light conditions under ice were estimated from continuous measurements of photosynthetically active radiation (PAR) $I(z)$ at the single depth $z = 1.5 \text{ m}$ under the ice surface assuming one-band exponential decay of radiation flux (Beer Law) $I(z) = I_0 \exp(-\gamma z)$. We estimated the decay rate of radiation within the water column (light attenuation coefficient) γ using PAR profiles collected in previous studies in 2011, using the evidence that year-to-year variations of water transparency of Lake Baikal are small (Hampton et al., 2008). The light attenuation coefficient was estimated from 9 PAR profiles as $\gamma \approx 0.17 \pm 0.01 \text{ m}^{-1}$. The radiation flux at the ice bottom amounted at $\approx 8-18 \approx 1-18\%$ of the surface radiation, and varied depending on the snow conditions at the ice surface (Fig. 2d). The mean daily under-ice short-wave radiation was $I_0 = 9.7 \text{ W m}^{-2}$ with maximum values of up to 23.5 W m^{-2} . The drop of the under-ice solar radiation after 23 Feb (Fig. 2d) was caused by a (relatively light, <0.5 cm) snowfall, which prevented the light penetration through the otherwise transparent

congelation ice. Variations in the under-ice radiation could have affected the temperature distribution under ice by slowing down or even canceling the warming in the convectively mixed layer at depths below 10 m.

With γ and I_0 known, we estimated the theoretical thickness of the stratified interfacial layer δ_R . For a single-band exponential decay of the short-wave solar radiation within the water column $I(z) = I(0) \exp(-\gamma z)$, the steady-state solution of radiation-conduction balance in a layer of thickness δ_R can be written as (Barnes and Hobbie, 1960)

$$\gamma \kappa T_m + I_0([1 + \gamma \delta_R] e^{-\gamma \delta_R} - 1) = 0, \quad (23)$$

which represents a transcendental equation with respect to δ_R . When substituted in Eq (23), the values of I_0 and γ yield $\delta_R \approx 0.2-0.4$ m, adopting the temperature of the well-mixed layer of 0.6 °C for T_m . The estimate coincides well with the observed thickness of the ice-adjacent gradient layer δ (Fig.3). The non-zero vertical temperature gradient beneath this layer is contrasting to the typical picture of convection in ice-covered lakes and suggests that the part of the convectively mixed layer $\delta < z < h_S$ is altered by the turbulent shear due to under-ice currents. Based on this suggestion, the under-ice radiation values were used to estimate the destabilizing buoyancy flux from Eq. (8) across the linearly stratified layer $\delta < z < h_S$ as $B_R = g\alpha I_R = g\alpha \left(I(\delta) + I(h_S) - 2h_S^{-1} \int_{\delta}^{h_S} I(z) dz \right)$. The resulting estimations are $I_R \approx 2$ W m⁻², and $B_R \approx 2.5 \cdot 10^{-10}$ m² s⁻³. The characteristic scale of convective velocities (Deardorff, 1970) $w_* = (B_R h_S)^{1/3} \approx 1.3$ mm s⁻¹, which value agrees well with previous reports on radiative convection under lake ice (Mironov et al., 2002; Kirillin et al., 2018b; Volkov et al., 2018).

These estimates of the convective velocities suggest that radiation was of minor importance for the mixing conditions in the boundary layer compared with the mean flow (cf. Fig. 5), especially after 23 Feb, when the subsurface radiation level dropped significantly. Therefore, shear and stratification in the boundary layer appeared to be the major factors determining water-ice heat transport.

20 4.4 Turbulence intensities in the ice-water boundary layer

Fluctuations of current velocities around their means were characteristic of developed turbulence: The structure functions (22) scaled well as the distance in 2/3 degree, clearly demonstrating existence of the inertial interval in the wavenumber domain (Fig. 6a). According to the 2/3-scaling, the upper boundary of the inertial interval reached $0.1 - 0.3$ m, which can be treated as a characteristic size of turbulent eddies. In low turbulent conditions $\varepsilon < 10^{-9}$ m² s⁻³, the TKE dissipation rates were at their minimum at the depth of ≈ 0.8 m and increased towards the ice base (asterisks in Fig. 6b) supporting the scaling $\varepsilon \propto z^{-1}$ (Eq. 7). The scatter of ε around the straight line $\varepsilon^{-1} \propto z$ increased with the distance from the ice z , starting from the $z \approx L_e$. During periods of high turbulence ($\varepsilon > 10^{-8}$ m² s⁻³), the reciprocal of the TKE dissipation rate ε^{-1} increased with depth more homogeneously. Nevertheless, a small local extreme in the line $\varepsilon^{-1}(z)$ and a slight change of the slope were recognizable at the same critical distance $z \approx L_e \approx 0.8$ m from the ice (circles in Fig. 6b).

30 In the area with weak under-ice currents at S2, the TKE dissipation rates ε varied around a value of 10^{-9} m² s⁻³, close to the a threshold between turbulent and laminar conditions. In turn, the average ε in the vicinity of the jet-like under-ice current at S1 was two orders of magnitude higher (Fig. 7). In contrast to the under-ice water temperatures, neither TKE dissipation rates nor friction velocities demonstrated any diurnal variations, suggesting minor effect of the radiation-driven convection on the

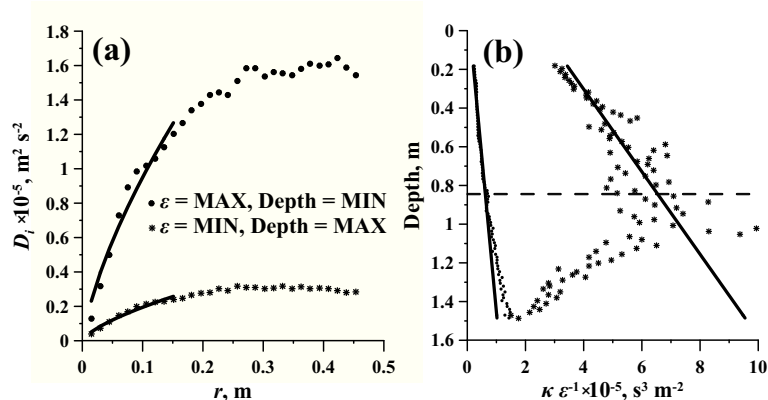


Figure 6. Turbulence-related characteristics of the boundary layer: (a) Velocity structure functions for high (circles) and low (asterisks) levels of TKE dissipation rates. Solid lines are the approximations by Eq. (22). (b) Vertical profiles of the reciprocal TKE dissipation rates $\kappa \varepsilon^{-1}$ for high (circles) and low (asterisks) turbulent conditions. Solid lines are the data approximations by Eq. (7). Horizontal dashed line in Panel b marks the depth equal to the mean Ozmidov length $L_N \approx 0.8 \text{ m}$

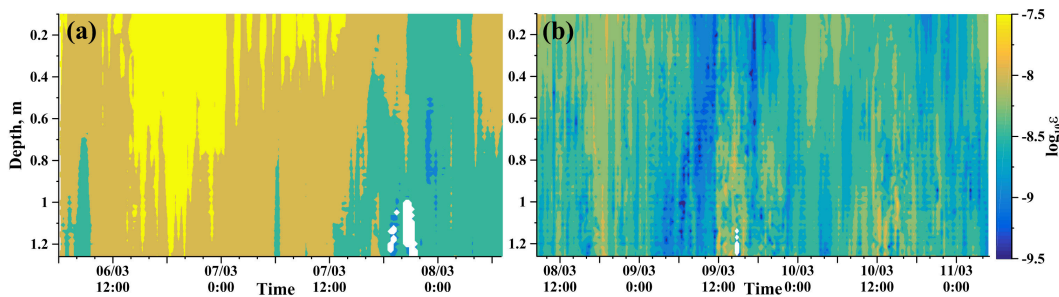


Figure 7. The TKE dissipation rates in the area of the jet stream (station S1, Panel a) and in the region of weak currents (station S2, Panel b).

turbulence generation. Instead, an apparent correlation existed between the turbulence intensity ε and the temporal variations of the mean flow velocities (Fig. 7): the highest TKE dissipation rates of $\mathcal{O}(10^{-7}) \text{ m}^2 \text{ s}^{-3}$ were observed during the currents intensification up to $\mathcal{O}(10^{-1}) \text{ m s}^{-1}$ at S1.

The maximum values of the Dougherty-Ozmidov length scale (Eq. 10), averaged over the period of observations, decreased with the distance from the boundary from $L_N \approx 1.5 \text{ m}$ at $z \approx 0.2 \text{ m}$ to $L_N \approx 0.9 \text{ m}$ at $z \approx 0.9 \text{ m}$. The decrease in L_N followed the decrease of ε . Here, the mean N_S in the layer with quasi-linear stratification of 0.5-10 m beneath the ice base was used as a characteristic value of the buoyancy frequency in the Dougherty-Ozmidov (D-O) scaling. At larger distances from the ice base, L_N remained nearly constant, varying between 0.8-0.9 m. Hence, the value $z_{crit} = L_N = 0.85$ can be treated as a boundary between the “quasi-neutral” and strongly stratified “z-less” layers, with the turbulent length scale defined by the distance z closer to the ice base, and by L_N at farther distances.

An important insight into the mechanisms of turbulence generation under ice is provided by comparison of the stratification-based turbulence scaling (Eqs. 10-13) against the quasi-neutral Law-of-the-wall (LOW) relationship (Eqs. 5-7-5). At $z < z_{crit}$, both “neutral” relationships (4) and (7) produced similar estimations of the friction velocities u_* with 20-30% higher values produced by u_* estimations from ε (Eq. 7). While the value of the von Kármán constant in the neutral LOW scaling $\kappa \approx 0.4$ is relatively well known from tunnel experiments and numerical simulations, and is supported by field data, the proportionality constant in Eq. (13) is not well established. Therefore, for “z-less” D-O scaling (13), the values of u_* were calculated from Eq. (13) assuming a unity coefficient of proportionality C_N . In average, the “stratified” scaling produced generally lower values of u_* at farther distances from the ice bottom and vice versa. The D-O scaling with $C_N = 1$ and LOW demonstrated nearly perfect agreement at $z = z_{crit} = L_N$ (Fig. 8a). The fact justified the balance between the shear producing at the boundary and the stratified production of turbulence at this distance from the wall, as well as supported the choice of the unity constant in the D-O scaling. Accordingly, $C_N = 1$ was adopted for later application of the combined log-linear scaling (Eqs. 12 and 14). The combined log-linear scaling (Eq. 12) with $C_N = 1$ produced u_* close to the neutral value in the vicinity of z_{crit} and decreasing towards both IWI and the open water column (Fig. 8b). Like the TKE dissipation rates, the mean current speeds demonstrated behavior characteristic of the stationary boundary layer, i.e. fitted well to logarithmic profiles at distances from the IWI less than z_{crit} (Fig. 8c). Farther from the ice base the mean velocity profiles were nearly linear, with the slope close to zL_N^{-1} .

The integral balance between the TKE loss terms $u_*^2 N + \varepsilon$ and the turbulent energy production $u_*^2 S$ (Eq. 7) held true within the 1.5 m thick layer covered by measurements of ε : The mean difference between the two terms integrated over the entire layer did not exceed 0.2%; temporal variations of the dissipation rate followed closely those of the shear velocity (Fig. 9ba). The balance was disturbed only at current speeds $< 0.02 \text{ m s}^{-1}$, with a corresponding drop of ε down to $< 10^{-9} \text{ m}^2 \text{ s}^{-3}$ at station S2, March 9, 2017 (not shown). During this same period, the vertical flow profiles showed a significant deviation from the logarithmic form, indicating laminarization of the boundary layer under these conditions. The boundary value of the friction velocity for the transition to turbulent regime was $u_* \approx 1.0 \text{ mm s}^{-1}$. The mean balance between the production loss terms in Eq. (14) varied however with the distance from the ice base (Fig. 9ab): close to the ice-water interface the production significantly exceeded dissipation, while below the depth of $\approx 0.8 \text{ m}$ the dissipation prevailed above the production. Remarkably, this transition depth agreed well with the thickness of the layer where $\varepsilon \propto z^{-1}$.

The good agreement of the measured velocity profiles with the logarithmic approximation at $z < z_{crit}$ allowed estimation of the roughness of the ice bottom surface z_0 from Eq. (5). The mean z_0 amounted at 1.00 mm with maximum of 3.5 mm and minimum of 0.2 mm. The roughness had a significant (the Pearson coefficient of -0.52 , $p \ll 0.01$) negative correlation with the mean velocity as $z_0 \approx 1.2 \cdot 10^{-4} U_{mean}^{-1}$.

Our estimations of z_0 and u_* yielded the following parameters for the bulk formula Eq. (6) : $C_{1m} \approx 3.4 \cdot 10^{-3}$ and $C_Z = C_{1m} Z^{-1}$, where C_{1m} is the bulk transfer coefficient for the momentum flux at 1 m depth. The independent measurements of current velocities at 4 depths made by the single-point 2-D horizontal current loggers INFINITY demonstrated good agreement with Eq. (6) when scaled against u_* from high-resolution AQUADOPP measurements (Fig. 10). The simple result has a large potential for application in modeling of the ice-water boundary layer at strong under-ice currents with a minimum of input

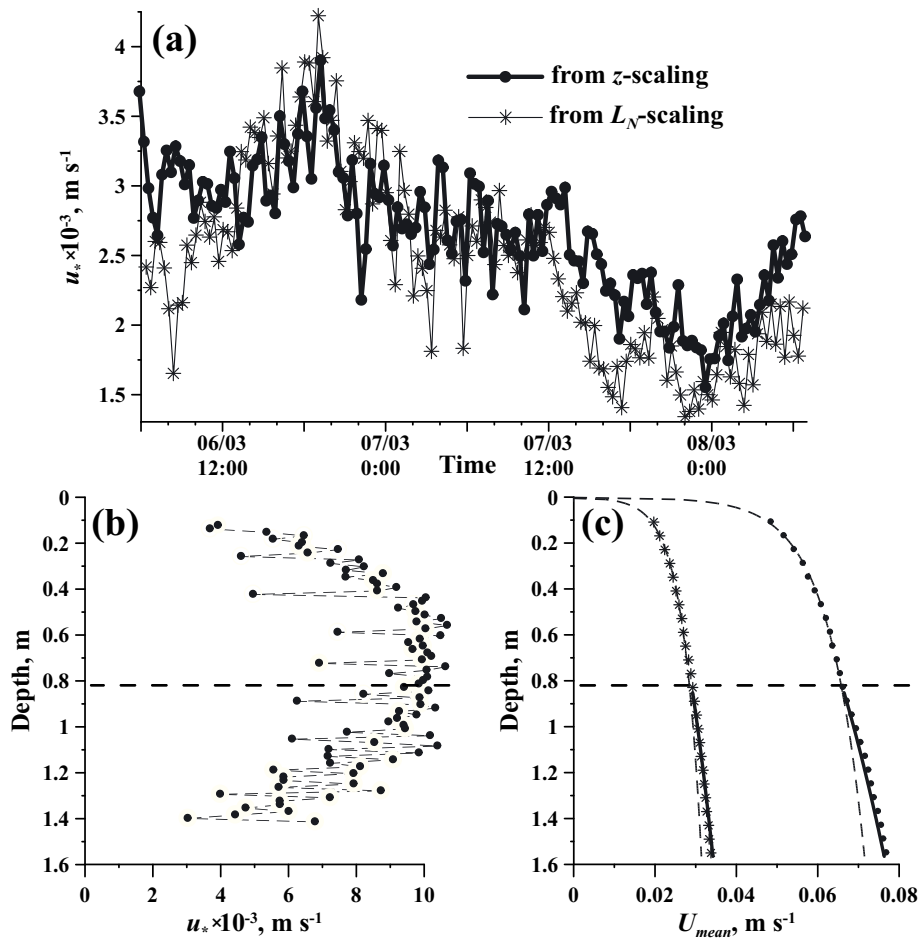


Figure 8. Ice boundary layer structure: (a) friction velocity u_* at station S1 as determined from the quasi-neutral production-dissipation balance (Eq. 7, thick solid line with symbols) and from the Dougherty-Ozmidov length scale (Eq. 13, thin lines) at the distance from the ice base nearly equal to the mean Ozmidov length $L_N \approx 0.8$ m; (b) mean vertical profiles of u_* , and (c) vertical velocity profiles for strong (circles) and weak (asterisks) currents. The two profiles in Panel (c) correspond to those in Fig. 6. Horizontal dashed line in Panels (b,c) marks the depth equal to the mean Ozmidov length $L_N \approx 0.8$ m

information; a care should be taken however about the thickness of the nearly logarithmic layer and its dependence on the under-ice stratification—the effect described above and discussed in more details below.

4.5 Heat budget at the ice base and relation of ice-water heat flux to under-ice mixing

The heat balance at the IWI was calculated by Eqs. (1) and (3) using data on temperature within the ice cover and ice thickness variability measured by the echosounder. The ice-water heat flux was generally stronger at Station S1, correlated with stronger currents and mixing intensities. Already at the beginning of the observations period in February, the upward conductive heat

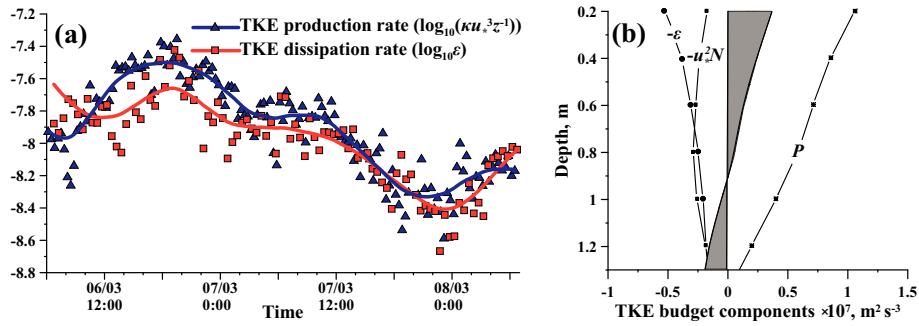


Figure 9. TKE production-dissipation balance in the ice boundary layer. (a) TKE production and dissipation rates at the depth of 0.85 m beneath the ice base. Bold lines are the values filtered by a moving average with a 6-hour window; (b) vertical profiles of the TKE budget components in Eq. (14) averaged over the 2 day period of measurements. Data are presented for the station S1 only; the S2 results are qualitatively the same.

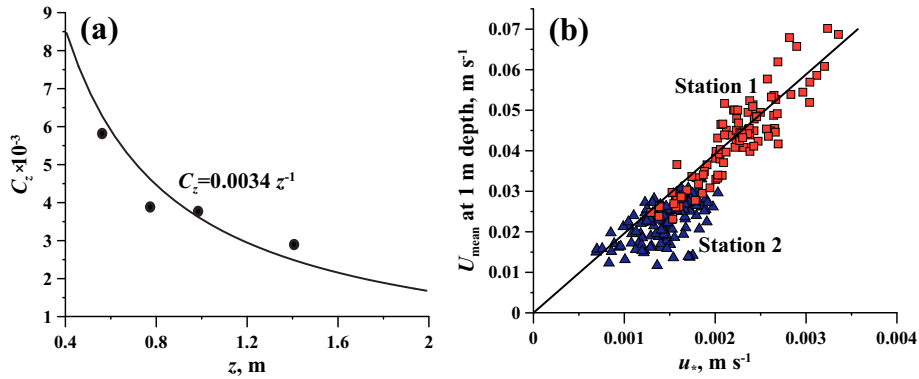


Figure 10. Bulk characteristics of the ice boundary layer: (a) bulk transfer coefficient for momentum $C_T = C_z$ as a function of distance from the IWI. Dots are values calculated using the horizontal velocity data from point measurements by 2-d INFINITY loggers, line is an analytical approximation; (b) relationship between the mean flow velocity at $z = 1$ m and the turbulent stress.

loss Q_{ci} of up to 80 W m^{-2} was compensated to 30-50% by the heat supply from the water column $Q_{cw} = Q_{iw}$. The latter significantly reduced the ice growth rate and latent heat release (cf. red and blue areas in Fig. 11). During periods of currents intensification at S1 (24 Feb - 07 Mar), the snow cover reduced the heat release to the atmosphere that also lowered the conductive flux at the ice base Q_{ci} (Fig. 11a). As a result, the heat flux from water to the ice at S1 (Fig. 11a) exceeded that from the ice to the atmosphere, producing melting at the ice base (negative Q_L) despite continuing surface cooling (positive Q_{ci} remained positive with values of up to 40 W m^{-2}). After 25 Mar, ice cover started to melt at both stations, coinciding with an increase of air temperatures above the freezing point (Fig. 2). Quantitatively, the ice-water heat fluxes at S2 were in the range of $5\text{-}10 \text{ W m}^{-2}$, which agrees with estimations from earlier lake studies. However, $Q_{cw} = Q_{iw}$ at S1 had appreciably higher values, reaching up to 40 W m^{-2} at their peaks.

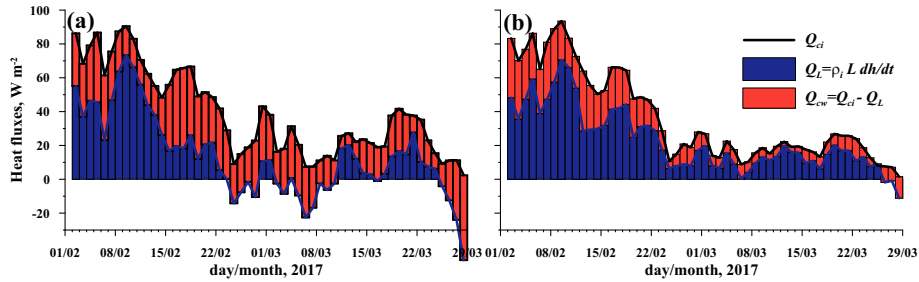


Figure 11. Daily averaged heat balance at (a) station S1 and (b) station S2.

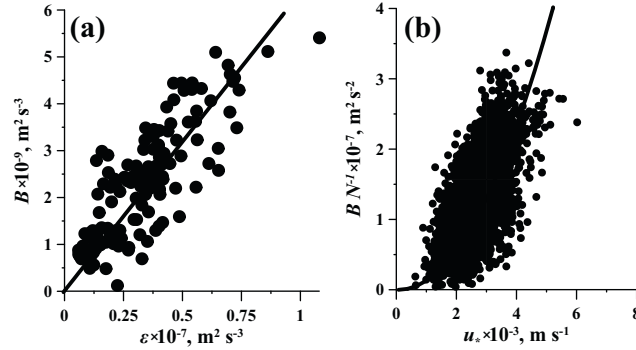


Figure 12. Buoyancy flux at the ice-water interface B as a function of (a) TKE dissipation rate ε and (b) friction velocity u_* .

— An attempt to link the ice-water heat flux with the mixing characteristics in the stratified boundary layer in form of a bulk relationship (Eq. 19) provided a remarkable result: the heat flux at the IWI and the dissipation rate of the TKE are linked linearly (Fig. 12a) or, in terms of buoyancy flux B :

$$B = g\alpha Q_{cw} = 0.065 \varepsilon. \quad (24)$$

- 5 Here, ε is taken at the distance $z_{crit} = 0.85$ m from the ice base that corresponds to the boundary between the ice-adjacent sublayer and the linearly stratified boundary layer (see Section 4.2). The linear correlation between the ice-water heat flux and ε supports the scaling (20) in the stratified boundary layer under ice: ~~From from~~ the two bulk relationships Eq. (18) and Eq. (20), the former suggests $Q_{cw} \propto \varepsilon^{1/3}$ $Q_{iw} \propto \varepsilon^{1/3}$ and the latter agrees with the observed linear dependence $Q_{cw} \propto \varepsilon$ $Q_{iw} \propto \varepsilon$. Here-
- 10 with, the result discards the widely used bulk relationship $Q_{cw} \propto u_*$ $Q_{iw} \propto u_*$ and suggests that Q_{cw} Q_{iw} scales as a friction velocity squared (Eq. 20) or, in terms of Eq. (19), $C_Q \propto u_*$. The dependence $Q_{cw} \propto u_*^2$ $Q_{iw} \propto u_*^2$ is well supported by our data, though the scaling is less apparent at very low u_* due to the lack of values at low turbulence levels. Still, the data on $Q_{cw}(u_*)$ $Q_{iw}(u_*)$ (Fig. 12b) clearly demonstrate the inappropriateness of the “quasi-neutral” scaling (Eq. 18). Instead, the flux can be parameterized as

$$B_{iw} = 0.015 u_*^2 N_S, \quad (25)$$

where N_S is the quasi-constant buoyancy frequency in the boundary layer $0.4 < z < 10$ m.

5 Discussion

Our study presents the first detailed assessment of mixing conditions under the ice cover of Lake Baikal and their effect on the growth and melt of the ice cover. The seasonal ice cover is the Baikal's inherent feature, whose role in functioning of this unique ecosystem remains not fully understood. In this regard, the outcomes of the study underscore the importance of the lake-wide circulation for the ice cover duration and ice thickness. The applicability of the results extends however beyond the specific Baikal conditions. Lake Baikal shares the major features of the seasonal ice cover on other lakes, as well as on inland and marginal seas allowing extension of the results on other ice-covered waters. Besides, the ice-water boundary layer in Lake Baikal owns a remarkable feature relevant to fundamental problems of environmental fluid mechanics: A strong boundary-layer flow on the background of permanent stable density stratification. In our study we successfully tested an alternative approach to the traditional Monin-Obukhov similarity theory, based on the Dougherty-Ozmidov scaling. We also revealed several important facets of the turbulent energy budget in stratified boundary layers, as well as established a relationship between the shear turbulence under ice and the heat flux at the ice-water boundary.

The high values of water-ice heat fluxes in Lake Baikal and their apparent relationship to the intensity of under-ice [circulation currents](#) were previously noted by Aslamov et al. (2014a, 2017). In the present study, the measured fluxes reached up to 40 W m^{-2} at their peaks, which is an order of magnitude higher than values reported from small [lakes Kirillin et al. \(2018b\)](#)-[arctic lakes \(Kirillin et al., 2018b\)](#) and comparable to the highest reported values [of oceanic fluxes to the ice at the melting stage \(Peterson et al., 2017\)](#) [in alpine thermokarst ponds \(Huang et al., 2019\)](#) and [in the ocean \(Gallaher et al., 2016; Peterson et al., 2017\)](#). Concurrent registration of fluxes, current velocities and dissipation rates of the TKE reveals the direct link between turbulence production by the velocity shear and ice growth (ablation). The finding contradicts to the conventional assumption on the major role of convection produced by the solar radiation penetrating the ice in the under-ice mixing of freshwater lakes. While radiatively-driven convection is a prominent feature of freshwater lakes (Farmer, 1975; Yang et al., 2017; Volkov et al., 2018), and effectively mixes the upper water column of Lake Baikal in winter (Granin et al., 2000; Jewson et al., 2009), its effect on the boundary mixing and heat transport to the ice base appears to be restricted by the stable stratification in the relatively thin interfacial layer with water temperature increase from the melting point at the IWI to that of the convectively mixed layer (Kirillin et al., 2018b). As a result, the energy of convection partially dissipates within the convective layer and is partially spent for entrainment of deeper waters at the base of the mixed layer (Mironov et al., 2002). The rate of the energy dissipation produced by convection in small lakes (Volkov et al., 2018) is about 10^{-9} - $10^{-8} \text{ W kg}^{-1}$, which is roughly an order of magnitude lower than that measured in this study. Consequently, the turbulence budget in the boundary layer differs significantly from that reported in previous lake studies. The TKE production averaged over the entire observations period prevails over dissipation at $z < L_N$, while in the deeper part of the boundary layer ε slightly exceeds the production. Herewith, the boundary mixing continuously pumps turbulent energy downwards, contributing to destroying of stratification in the main stratified layer SL $\delta < z < h_S$. The opposing trend is created by solar radiation, which increases the temperature of the convectively mixed

layer T_m and creates by this the upward exponential decrease of temperature from T_m to the freezing point T_f . The observed quasi-linear stratification in the SL is apparently a result of the two opposing forces leading to a nearly steady-state conditions characterized by a constant N within the layer $\delta < z < h_S$.

5 A particular advance in estimation of the turbulent energy budget under ice was achieved by direct estimation of the TKE dissipation rate using the velocity structure method. By this, we (i) avoided applying of the Taylor’s “frozen turbulence” hypothesis, which remains questionable at relatively low current velocities under ice, and (ii) were able to trace the vertical distribution of ε across the boundary layer (at least, a part of it, covering ~ 2 m). On one hand, this extended information allows better quantification of the turbulent structure, on the other hand, it poses some fundamental questions on the major forces behind the under-ice mixing and heat transport to be discussed below.

10 Close to the ice base, vertical profiles of the TKE dissipation rates decayed as $\varepsilon \propto z^{-1}$, supporting the scaling of the turbulent mixing length with the distance from the solid boundary, similar to the neutral or nearly-neutral conditions. This fact gives a solid background to estimation of the shear velocity u_* from the mean velocity profiles: The latter method is often uncertain, given both depth of the logarithmic layer and the ice roughness are not known a priori. However, the thickness of the layer with $\varepsilon \propto z^{-1}$, varies depending on the ~~current~~currents speed and mixing intensity. At strong currents, dissipation followed
15 the scaling across the entire depth of the high-resolution velocity measurements of 1.4 m; at $\varepsilon \ll 10^{-7}$, it reduced to several tens of centimeters. The maximum mixing length scales remarkably well against the Dougherty-Ozmidov length scale: using N_S as the characteristic value for stratification, $L_N = 1.2$ m at $\varepsilon = 10^{-7}$ W kg $^{-1}$ and $L_N = 0.4$ m at $\varepsilon = 10^{-8}$ W kg $^{-1}$. Qualitatively the result can be treated as follows: the shear at the ice base dominates in the turbulence production at distances from the ice base less than L_N , while farther from the source of the shear the stratification limits the size of the turbulent
20 eddies. This structure is also supported by comparison of the friction velocities computed from the two mixing lengths, Eq. (7) and Eq. (13): u_* tends to be overestimated by the Ozmidov scaling at high mixing rates close to the ice base, and becomes lower than that produced by z-scaling at larger distances from the ice. At $z \approx L_N$ both estimations provide equal results (see Fig. 8). In the TKE budget, the depth $z > L_N$ is also a turning point, where Eq. (14) shows a close balance between turbulence production and damping.

25 Generally, the simplified model of the TKE balance Eq. (14) was well supported by the data. The discrepancies included the prevalence of the TKE production over the damping terms closer to the ice cover and a slightly lower production than the sum of the lost terms at $z > L_N$. The imbalance can be tentatively attributed to downward advection of the TKE. Besides, the balance was estimated in assumption of the constant buoyancy frequency N , neglecting a possibly stronger loss of the TKE closer to the ice, where N increases. We also neglected possible transport of turbulent energy produced by convection due to
30 solar heating. The latter can however be assumed small: within the stratified layer, the radiation levels under ice produced the destabilizing buoyancy flux of only $\mathcal{O}(10^{-10})$ W kg $^{-1}$, and the convectively mixed layer was located several meters beneath the deepest point covered by measurements.

The proposed scaling of the TKE budget differs from the conventional MOST approach by using the Dougherty-Ozmidov length scale instead of the Obukhov length scale, i.e. by replacing the surface buoyancy flux with the mean stratification as
35 a major scaling variable. This alternative approach is convenient for analysis of observational data: in contrast to the surface

fluxes, N is easily measurable in oceanic and lake studies. Noteworthy, both scaling approaches have been shown to be interchangeable (Grachev et al., 2015). In the particular case of ice-covered waters, the D-O approach is also more physically sound than MOST, since the surface buoyancy flux does not dominate the turbulent conditions under ice, it is rather a result of the upward heat transport from deeper waters. One of the derivatives of the D-O boundary layer scaling is the relationship $\varepsilon \propto u_* N \varepsilon \propto u_*^2 N$, which also implies that the length scale $u_* N^{-1}$ can be used instead of L_N without any basic changes in the model assumptions. This scaling was previously considered e.g. by (Zilitinkevich and Mironov, 1996) and can be preferable for testing and refining the model parameters if observations of u_* are available at high resolution.

The presence of an ice-adjacent interfacial sublayer $0 < z < \delta$ is another remarkable feature of the shear-dominated IBL. The thickness of the layer ($\delta \approx 0.2$ m) was close to the smallest estimate based on the solar radiation of the layer created between the fixed temperature at the ice base and a convectively mixed homogeneous layer beneath. Our temperature measurements were too scarce to trace the evolution of its thickness at variable current velocities. Some insight into the genesis of the layer can be however obtained by assuming the largest value of the Dougherty-Ozmidov length scale $L_N \approx 0.85$ m as the maximal thickness of δ . In this case, the mean buoyancy frequency in the interfacial sublayer $N_\delta \approx 0.02$ s⁻¹. That leads to the buoyancy Reynolds number $Re_b \approx 16$ and the gradient Richardson number $Ri_g \approx 0.4$. Both values are close to the critical values of $\mathcal{O}(10^1)$ and 0.25, respectively. The layer $0 < z < \delta$ may be therefore assumed to stay in near-critical turbulence-free state. In the rest of the boundary layer, both Ri_g and Re_b are far beyond the critical values, indicating developed turbulence.

The following structure can be tentatively drawn from the above analysis (Fig. 13): the background vertical temperature (density) distribution is formed by absorption of solar radiation and the upward heat flux to the ice, producing the profile typically observed in ice-covered lakes without strong currents: a nearly homogeneous convectively mixed layer with a relatively thin stratified layer on top. Mixing by the velocity shear reduces the density gradient. In the top layer, shear mixing is balanced by solar heating from above, so that the density gradient tends to the critical value between turbulent and non-turbulent states. In the upper part of the homogeneous layer, a weaker nearly linear density gradient forms. In the larger part of this layer, at $z > L_N$, the turbulence production is balanced by the stratification, in accordance with “z-less” linear scaling, Eq. (13). Beneath the depth L_N , the scaling suggests $\partial U \partial z^{-1} = const$ that is also supported by the measured mean velocities (Fig. 8). The stratification is nearly linear suggesting also that ε is nearly constant across it. The thickness of the layer depends apparently on the scales of the under-ice flow. The shape observed temperature profiles slightly deviate from linearity, with a tendency to re-stratification corresponding to weaker currents at S2 (Fig. 3b). Apparently, variations in the current speeds may affect the general tendency to linear stratification, producing changes in the heat content of the upper water column caused by horizontal advection of heat, or by re-stratification during periods of low shear. We do not possess enough well-resolved data on temperature to investigate these effects.

We did not consider the rotational effects on the boundary layer characteristics. A slight Ekman-like rotation of the mean current was observed under ice (Fig. 3), and the Ekman scale $u_* f^{-1}$ based on the mean shear velocities is $\mathcal{O}(10^1)$ m. Herewith, the Coriolis force may have an effect on the thickness of the boundary layer. At weak stratification, effects of both N and f on the boundary layer dynamics may appear comparable, so that the model will perform better if the length scale L_N (or its equivalent $u_* N^{-1}$) is replaced by the combined length scale $u_* (fN)^{-1/2}$. The latter scaling was proposed by (Pollard et al.,

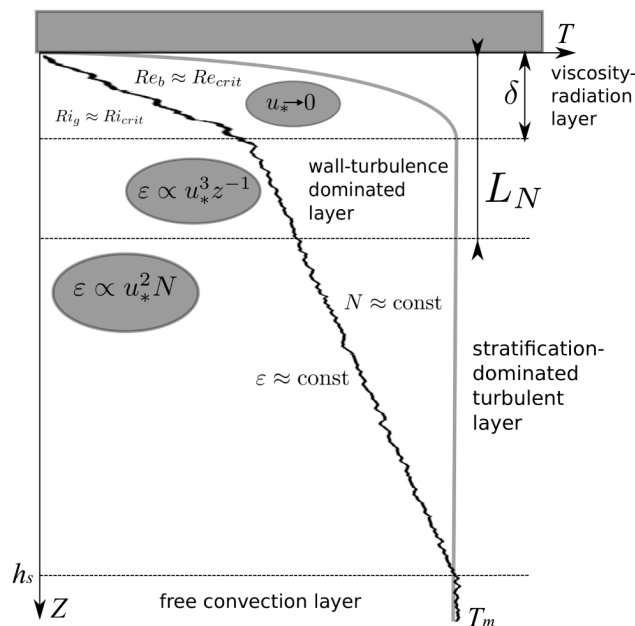


Figure 13. Formation of the vertical temperature profile in a shear-driven ice boundary layer. Gray line is a no-shear original profile, black line is the result of shear-driven turbulence on the background of stable stratification. See text for notations.

1972). Zilitinkevich and Mironov (1996) suggested that the fN -scaling has rather limited area of application. However low flow velocities and weak stratification are typical for conditions under ice, and the Pollard et al. (1972)-scaling may find its application in ice-covered waters. More data are required here, in particular, on the fine density and flow structure over the entire Ekman layer.

- 5 The main motivation of the study was seeking for the relationship between the shear mixing and stratification on one hand and the heat release from water to the ice cover on the other hand. In this regard, the scaling of the heat (buoyancy) flux with the shear velocity squared and the buoyancy frequency under ice (Eq. 4.520) is our major finding with implications to a wide range of the ice-related problems. The simple “bulk” approximation Eq. (19) is widely used in models of ice-water interaction, but its validity was never thoroughly tested before. The simple relationship $Q_{cw} \propto u_* \Delta T$ $Q_{iw} \propto u_* \Delta T$, equivalent to Eq. (19), failed
- 10 to describe the heat flux dynamics in the ice boundary layer, replaced by the scaling $Q_{cw} \propto (g\alpha_T)^{-1} u_*^2 N$ $Q_{iw} \propto (g\alpha_T)^{-1} u_*^2 N$. The result suggests that the effect of the currents on the decrease of the arctic sea ice may appear much stronger than assumed by the present model projections, based on the bulk estimation $Q_{iw} \propto u_*$. Further decline of the arctic ice cover may result in both increase of the under-ice current speeds due to changed global circulation and increase of the stratification due to warming of the under-ice waters. Both factors, according to our scaling will accelerate the vertical heat transport from water to the ice,
- 15 resulting in a positive feedback on the ice melt. Hence, incorporation of the under-ice density stratification in the heat flux parameterizations may improve significantly the outcomes of the oceanic ice models, especially for the ice melt periods, when freshening of the under-ice boundary layer produces strong salinity stratification.

In contrast to the surface heat flux, the bulk formulation Eq. (6) worked well for the momentum flux thanks to the nearly constant shear conditions close to the ice base. The result is of practical use in simple models of ice covered seas and lakes, where u_* can be directly approximated from the mean current speeds at a certain distance from the IWI (Fig. 10). The speeds should be however known at distances from the ice $z < L_N$, otherwise, the stratification effects on turbulence production make
5 the bulk formula not representative.

6 Conclusions

We investigated the fine vertical structure of turbulence characteristics in the boundary layer of Lake Baikal and proposed a model of stratified turbulent ice boundary layer based on the Dougherty-Ozmidov length scale of turbulence. The-In contrast to small lakes, where solar radiation dominates ice-water heat exchange, the water-ice heat flux in Lake Baikal was strongly
10 affected by the mean flow shear, similar to the ice boundary layer in the ocean. The shear-produced mixing was counteracted by the stable density stratification beneath the ice cover. Absolute values of the water-ice fluxes were an order of magnitude higher than those in no-shear conditions, ensuring basal ice melt even during cooling at the ice surface. The ultimate result consists in scaling of the water-ice heat flux against the shear velocity squared. The result suggests large errors in the heat flux estimations, when the traditional “bulk” approach is applied to stratified conditions with strong shear. It also implies that
15 under-ice currents may have a much stronger effect on the ice melt than estimated by traditional models.

Data availability. Data are available from the authors by request

Author contributions. GK, IA and NG conceived the study; IA, RZ, and GK designed and performed the field experiments; GK developed the model; GK, IA, VK, and NG performed field data analysis; GK wrote the manuscript with contributions from all co-authors

Competing interests. The authors declare no competing interests

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Interactive comment on “Turbulence in the stratified boundary layer under ice: observations from Lake Baikal and a new similarity model” by Georgiy Kirillin et al.

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We thank the Reviewer for the careful reading of our study and the thoughtful comments. We particularly welcome the comments from the oceanographic point of view, which help to clarify the similarities and differences between freshwater and oceanic ice-covered systems. The discussion will surely help our study to reach the oceanographic audience, and enable implementation of our results on wider scales. Below is the point-to-point reply to the comments.

Pg. 2 Line 8: Reference in manuscript about sea ice loss attributed “primarily” to basal ice melt (ocean-to-ice) as opposed to surface ice melt (air-to-ice) is still under debate.

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Ice mass balance (IMB) observations shows that the amount of surface (atmospheric) and basal (oceanic) melt varies with each year (some years the top melts about the same as the bottom). I recommend rewording this sentence to state that a significant component of sea ice volume loss occurs from the sea ice bottom... (or something like “due to ocean-to-ice heat fluxes”)

When pointing on the primary role of the basal melt, we referred to:

Carmack, et al.: Toward quantifying the increasing role of oceanic heat in sea ice loss in the new Arctic, Bulletin of the American Meteorological Society, 96, 2079-2105, 2015.

We appreciate that uncertainty may still exist in the estimates of the relative role of surface and basal melting, and rephrase the sentence in a less categorical way, as suggested.

Pg. 2 Lines 16-28:

I’m somewhat confused on the persistence of the interfacial layer (IL) during SML free convection. How is the IL maintained if the strongest heating (solar) is in the layer closest to the ice base. I assume this is due to either high sensible heat losses caused by the negative ice temperature gradients (thermal conductivity), or latent heat losses to the lake ice base (or combination of both in March); either way, the negative heat budget despite solar heating near the ice-water interface should briefly be addressed here...perhaps the best solution in the intro is to capture the dominating heat loss term during the period of your study (latent heat or sensible heat).

The raised issue is indeed quite important for understanding the under-ice boundary layer dynamics and is characteristic of fresh (and brackish) waters. In short: The stable density stratification prevents convective mixing despite the negative buoyancy production by the decrease of the solar radiation with depth.

Thanks to the freshwater density anomaly, density increases with temperature

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in fresh and brackish waters at temperatures below the maximum density point (≈ 3.98 °C for freshwater). At low salinities, water temperatures remain always higher than that of ice; hence the stably stratified interfacial layer (IL) with downward temperature increase is inevitable under the ice base. As a result, the heat budget in the IL is governed by the balance between the radiation absorption and—as Reviewer rightly suggested—upward heat conduction to the ice base. The balance was considered by Barnes and Hobbie (1960, referred in the paper), who proposed an elegant analytical solution of the heat transport equation for the thickness of the IL, temperature profile within it, and the water-ice heat flux in purely conductive conditions (see also Mironov et al. 2002, referred in the paper, for an extended discussion). Noteworthy, the IL is not purely conductive even in small lakes, characterized by intermittent wave-generated turbulence (Kirillin et al. 2018, referred in the paper), making the estimation of the ice-water heat exchange in the absence of mean flow particularly challenging. The deformation of IL in the presence of a strong mean flow is considered in Discussion section of the paper. We add a reference to the conduction-radiation model of Barnes and Hobbie to the introduction.

Pg. 2 Lines 21-23: In Arctic Ocean air-ice-water interactions, entrainment of subsurface heat (usually the near-surface temperature maximum (NSTM)) is hard to achieve with static instabilities (e.g. brine rejection), this is usually reserved for stronger dynamic (sheared) instabilities. If this statement (lake heat entrainment with static instabilities) has been demonstrated by previous work, please reference.

Convection due to solar heating in freshwaters is more energetic than mixing by brine rejection and represents a classical natural example of penetrative (entraining) convection in the absence of the mean flow shear with analogies in atmospheric, oceanic and astrophysical flows. The first detailed analysis was performed by Farmer (1975, referred in the paper), a review is given by Kirillin et al. (2012, referred in the paper).

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Pg. 3 Section 2: Regarding the geostrophic currents in Lake Baikal, request there be some background provided in this section as to the source of this current (pressure gradient force created by???? and spatial scale drives a low Rossby Number environment, etc.)

Being meticulous, we did not mention a geostrophic character of the observed currents in Section 2. Still, since Lake Baikal in winter is completely isolated from the direct contact with the atmosphere by the ice cover, the pressure (density) gradient is the major driver of under-ice flows. The characteristic velocity and spatial scales of the observed under-ice currents are (10^{-2} - 10^{-1} m s⁻¹) and (10^4 - 10^5 m), respectively (Figs. 1 and 5 of the paper). Hence, away from boundaries it is balanced mainly by the Coriolos force. The under-ice current of the same scales is persistent in this part of Lake Baikal (Aslamov et al. 2014, 2017, cited in the paper). Therefore, the geostrophy was mentioned in the abstract of the paper. The origin of the pressure gradient forcing is not exactly known: large-scale density fields under the Baikal ice were not measured. The usual suspects are the the wind-topography interactions, which create a regular large-scale snow cover pattern on the ice surface. As a result, the horizontal temperature gradients in under-ice waters are created by the inhomogeneous heating by solar radiation, as documented by Aslamov et al. (2017, cited in the paper). We add this information to the revised version and remove the only use of the word “geostrophic” from the abstract.

Pg. 8 Lines 17-18: Why were there more current meters deployed at S1 and not at S2?

The available amount of loggers did not allow to obtain a detailed vertical resolution at both stations. Therefore, bulk of the loggers were deployed at the primary site, with only one logger deployed at the reference site to background the AQUADOPP measurements.

Pg. 10 Figure 2: Perhaps I missed this in the results discussion, but why did the pen-

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etrated solar radiation drop off substantially after Feb. 22 when topside solar radiation increased (Figs 2c and 2d). Was there a snow event(s)? The only reference I can find to Fig. 2d is on page 13 and only accounts for the mean daily under-ice short-wave radiation ($I_0 = 9.7 \text{ W m}^{-2}$) and a range of through-ice radiation (8-18%). These results do not match with the results in Figs. 2c and 2d where the transmissivity (solar(under-ice)/solar(top-ice)) between March 6th and March 16th appears to be well below 1% with under-ice radiation values of $<2 \text{ W m}^{-2}$. The extremely low under-ice radiation values heavily skew the 9.7 W m^{-2} average over the study period and likely affects the intensity of short-wave induced convective overturning in the SML. There appears to be two “modes” to this dataset: 1) light snow cover prior to 22 February with active SMLs and strong ILs; and 2) moderate-heavy snow cover after 22 Feb with inactive SMLs and weak ILs. Request clarification on how this transition in the steady state condition was handled and why it is appropriate to conduct DO scaling model validations across these varying conditions.

It is a valuable comment for specifying the background conditions behind the under-ice boundary layer formation. Indeed, the drop of the under-ice solar radiation was caused by a (relatively light, $\approx 0.5 \text{ cm}$) snowfall, which prevented the light penetration through the otherwise transparent congelation ice. As correctly mentioned by the reviewer, variations in the under-ice radiation could have affected the temperature distribution under ice by slowing down or even canceling the warming in the convectively mixed layer at depths below 10 m. However, the estimates of the convective velocities w_* (Section 4.3) demonstrated that the radiation was of minor importance for the mixing conditions in the boundary layer compared with the shear instabilities u_* . Also, the background stratification did not change much during the period of observations because of the large thickness of the nearly homogeneous “free convection layer” and the strong effect of the shear mixing on the stratification above it (see Fig. 13 and the discussion around it). Therefore, shear and stratification in the boundary layer remain to be the major factors determining water-ice heat transport, making the D-O scal-

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ing to the ideal choice for parameterization of boundary fluxes. When modeling fluxes at longer time scales, changes in the background stratification caused by radiation variability are directly accounted for in the scaling by variations in the buoyancy frequency N .

We mention this fact in the revised paper and adjust the lower value of the sub-surface radiation.

Pg. 17 Lines 2-4: Once again, it appears that the event (likely snow) heavily impacted these results during the 24 Feb – 07 Mar period. If the 22 Feb event is snow, I anticipate it would affect several areas of the heat budget and near interface buoyancy to include lowering iceto-air sensible heat fluxes and destabilizing the IL (less downwelled solar radiation) allowing turbulent (shear) eddies access to the ice base. If this were indeed the case, it should probably be integrated into the discussion, if not, recommend addressing the cause of the significant change in heat balance conditions in Fig. 11 after 22 Feb (similar to the previous comment for page 10).

Similar to the previous comment, the remark correctly refers to the synoptic variability, which was only briefly mentioned in the original manuscript: apart from reducing the under-ice solar radiation, the snow cover reduced the heat release at the ice surface (the ice surface heat budget was not a subject of our study, though). As a result, the conductive heat flux at the ice base Q_{ci} reduced, while, as we mentioned in the paper, remained positive, with values up to 40 W m^{-2} . The major point here is: the turbulence due to the mean flow produced water-ice heat fluxes was sufficient to initiate ice melt at its base (cf. the negative peak in the fluxes in Fig. 11a and the velocity peak in Fig. 5a on 06-07.03.2017, Station 1).

We amend the sentence to clarify this issue.

Pg. 19 Lines 14-16: Not entirely accurate, oceanic fluxes during the 2014 MIZ experiment in the Beaufort Sea were $>100 \text{ W m}^{-2}$ with Autonomous Ocean Flux Buoys (Gallaher et al., 2016) and nearly 200 W m^{-2} in the Greenland Sea during the 1983/84

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MIZ experiment (McPhee et al., 1987)

The comment echoes the remark of Wenfeng Huang (see our response in the HESSD discussion). Our sentence is however formulated in a quite accurate way: the fluxes we reported are indeed significantly higher than those in small ice-covered lakes and are comparable to those reported in the ice-covered seas. Shallow alpine thermokarst ponds and drifting ice during strong storms in the ocean are the extreme examples worth mentioning, so we add references to Huang et al. (2019) and Gallaher et al. (2016).

Pg. 21 Line 6: I did not see an isothermal/isopycnal (homogeneous) layer mixed layer in the data (Fig. 3); perhaps, near-homogeneous is more appropriate.

Agreed. “Nearly homogeneous” is added.

Pg. 22 Line 16: Interesting idea to scale this DO scaling approach to sea ice modeling; however, near-freezing freshwater and seawater ice-water boundary layers have notable differences. Things that come immediately to mind are: 1) the rotational Ekman layer plays an important role (which was not tested in your study) in the deeper dynamically developed ocean boundary layers (20-35 cm/sec free drift ice speeds); 2) temperature becomes the equivalent of a passive tracer (no buoyancy contribution) in seawater above 25 ppt; and, 3) bulk parameterizations using MO scaling have worked pretty well when validated against eddy correlation and thermal dissipation observations. I will admit, that during calm wind conditions in the presence of significant meltwater (melt pond drainage), this parameterization does not perform well and is similar to your study minus the temperature stratification from solar heating. For this paragraph, I would recommend rewriting to target the potential benefit of this approach during weak atmospheric forcing over sea ice during the melt season.

A very appreciated while predictable comment. We agree that the M-O scaling would work in ice-covered seas at appreciable shear mixing, especially during ice growth. The D-O scaling would generally work in this case too, since both

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scalings can be derived from each other (see Grachev et al. 2015, cited in the paper). During the melting phase, in turn, the traditional M-O scaling can produce significant errors in the rates of basal ice melt, especially at strong under-ice stratification. Besides, the D-O scaling has several theoretical and practical advantages: Stratification in the boundary layer, not the buoyancy flux at the ice base, is the mechanism directly damping the turbulence under ice that makes N to a “natural choice” for scaling. Also, D-O scaling is easily applicable in practice: stratification is directly obtained from observations/models, which is often difficult with the boundary buoyancy flux. We use this opportunity to argue briefly, why the D-O scaling can be advantageous for modeling of the ice-ocean boundary layer:

(1) We have outlined in our Discussion how the Ekman forcing can be incorporated in the D-O scaling (Page 21, Lines 22-30). It does not seem to be a crucial issue however: The conventional M-O scaling does not include the Coriolis force either. See Zilitinkevich and Mironov (1996, referred in the paper) for a discussion on the scales, where the Coriolis forcing is important for the boundary layer scaling.

(2) Even if temperature has no contribution to the buoyancy in saline waters, the sea-ice boundary layer during the melt phase tends to be strongly stratified due to freshening (e.g., by basal ice melting, freshwater inflows, or melt pond drainage, as mentioned by the Reviewer), so that N becomes to the major factor dumping the shear-produced turbulence at the ice-water interface. The M-O scaling is apparently less relevant in this case.

(3) The neutral (M-O) scaling worked fine for the *momentum* fluxes in Lake Baikal too (in our Discussion we propose an explanation why). The impetus for the development of the new scaling behind the boundary-layer modeling was provided by the apparent inconsistency of the M-O scaling for the *scalar* (heat) fluxes, crucial for correct estimation of the ice melting rates. We strongly believe the

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scaling is more physically sound in ice-covered seas and is, at least, worth trying in marine models.

Pgs. 22-23 Section 6: Conclusion seems a little abbreviated, recommend recapping a few more of your findings.

We extend the conclusions taking into account the HESSD discussion.

Minor comments

We incorporated all suggestions in the revised version

Interactive comment on Hydrol. Earth Syst. Sci. Discuss., <https://doi.org/10.5194/hess-2019-608>, 2019.

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Interactive comment on “Turbulence in the stratified boundary layer under ice: observations from Lake Baikal and a new similarity model” by Georgiy Kirillin et al.

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We sincerely thank Dr. Huang for the valuable contribution to the discussion on our study. Indeed, the recently appeared study of Huang et al. (2019) on heat budget of a Tibetan ice-covered pond represents a great evidence of the large importance of the water-ice heat flux for the ice cover growth and melt: The absolute values of the flux in a thermokarst alpine lake are comparable or even larger than those found in Lake Baikal that makes the heat balance of two systems similar. Both in our study and in the study of Huang et al. (2019), the water-ice heat flux was estimated from the heat balance at the ice-water interface, using similar measurement setup with frozen ther-

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mistor chains. In that way, the results are inter-comparable, and we add in Discussion a reference to the study of Huang et al. The possible way of parameterization of the ice-water heat flux in thermokarst Tibetan lakes remains to be an intriguing question. The D-O scaling proposed in our study can only be used if information is available on the TKE production or dissipation rates, with subsequent parameterizations of the turbulence intensity based on the large-scale forcing. The latter is apparently different in thermokarst lakes than in Lake Baikal or other large lakes and seas. In deep Lake Baikal, in turn, the water column remains relatively cold throughout the ice season, but the strong upward heat flux is conditioned by the strong turbulent mixing due to large-scale under-ice currents. The upward heat release in the rather shallow (< 2 m mean depth) thermokarst lake is produced mainly by the heating of the water column due to the solar radiation, which is very strong over the Tibetan Plateau. As a result, the water temperatures under ice achieve values of up to 9°C , creating a strong gradient at the ice-water interface. The apparent source of turbulence is in this case free convection due to the negative buoyancy flux. In the typical conditions during the ice cover melt in Tibetan lakes, the mean water column temperature is above the maximum density values of $\approx 3.98^{\circ}\text{C}$. In this case, the convection is very non-stationary (see e.g. Kirillin and Terzhevik 2011), so that its parameterization is non-trivial. As rightly mentioned by Huang et al. in their discussion, direct measurements of turbulence under ice could provide the necessary quantitative information on the turbulence mixing in these quite specific conditions.

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We thank sincerely the Reviewer for the careful review and valuable suggestion on our study, which helped to improve the manuscript. Below we reply to the Reviewer's comments, providing only titles from original review:

Scaling arguments

- Both Q_{cw} and Q_{iw} were used interchangeably in the original version, which can be of course misleading to the reader - thank you for mentioning it. We have consistently changed all Q_{cw} 's with Q_{iw} .

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- Mironov et al. (2002) is cited right after Eq. (8)
- We added “kinematic” to viscosity for clarity.

Results

- Interesting questions are raised here about the relationship between the flow strength and the temperature profile in the upper 10 m of the water column. Qualitatively, the profiles at both sites are closer to straight lines than to an exponential “zero-shear” profile from Fig. 13, with the profile at lower current speeds (S1, Fig.3a) is slightly “convex”, i.e. closer to the theoretical curve from Fig. 13 and stronger currents correspond to a slightly “concave” shape of the temperature profile (S1, Fig.3a). In our discussion (Page 21, Fig. 13), we propose a nearly steady-state balance between stratification and shear as an explanation of the quasi-linear temperature profile. The deviation from the straight line, and as a result, the variations in the heat content of the upper water column could be caused by horizontal advection of heat by currents, or by restratification during periods of low shear. We do not possess enough well-resolved data on temperature to investigate these effects. We clarified this point in the text and also added time/date to Fig. 3.

Other minor points

Thankfully accepted.

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