



1 **Travel time based thermal tracer tomography**

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12 **Abstract:**

13 Active thermal tracer testing is a technique to get information about the flow and transport
14 properties of an aquifer. In this paper we propose an innovative methodology, using active
15 thermal tracers in a tomographical setup to reconstruct cross-well hydraulic conductivity
16 profiles. This is facilitated by assuming that the propagation of the injected thermal tracer is
17 mainly controlled by advection. To reduce the effects of density and viscosity changes and
18 thermal diffusion, early time diagnostics are used and specific travel times of the tracer
19 breakthrough curves are extracted. These travel times are inverted with an eikonal solver
20 using the staggered grid method to reduce constraints from the pre-defined grid geometry and
21 to improve the resolution. Finally, non-reliable pixels are removed from the derived hydraulic
22 conductivity tomograms. The method is applied to successfully reconstruct cross-well
23 profiles as well a 3-D block of a high resolution fluvio-aeolian aquifer analog dataset.
24 Sensitivity analysis reveals a negligible role of the injection temperature, but more attention
25 has to be drawn to other technical parameters such as the injection rate. This is investigated in
26 more detail through model-based testing using diverse hydraulic and thermal conditions in
27 order to delineate the feasible range of applications for the new tomographic approach.

28



29 **1 Introduction**

30 Tracers are commonly used to get an insight into the hydraulic properties of the subsurface
31 on the aquifer scale and to identify dominant transport routes. Among the many tracers used
32 for aquifer characterization, heat is frequently injected as a thermal tracer in boreholes or
33 wells (Anderson, 2005; Hermans et al., 2015; Rau et al., 2014; Saar, 2011). From measured
34 breakthrough curves (BTCs), aquifer heterogeneity and preferential flow paths are inferred
35 (Bakker et al., 2015; Colombani et al., 2015; Klepikova et al., 2014; Leaf et al., 2012;
36 Macfarlane et al., 2002; Vandenbohede et al., 2008; Wagner et al., 2014; Wildemeersch et al.,
37 2014).

38 Main attributes of ideal tracers are their good detectability, their lack of influence on the flow
39 regime, and nontoxicity to the environment. Heat is an ideal choice because it is easily
40 detectable by means of traditional temperature sensors or distributed temperature sensors
41 (DTS), and it can be monitored continuously in-situ. Typically, background variations are
42 insignificant, and natural heating-cooling cycles have smaller frequencies than the
43 investigated thermal signals. It is also ideal because moderate changes in temperature do not
44 harm the environment and thus commonly no regulative constraints are imposed. However,
45 due to possible viscosity and buoyancy effects, and its relationship with hydraulic
46 conductivity (K), variation in temperature may modify the flow regime. (Ma and Zheng,
47 2010) concluded from numerical simulations that no substantial density effects occur when
48 heating groundwater up by 15°C. This same critical value is given by (Russo and Taddia,
49 2010), based on the recommendations by (Schincariol and Schwartz, 1990) that buoyancy
50 effects only appear at density differences higher than 0.8 kg m⁻³. However, this calculation is
51 only valid if the groundwater temperature is close to 0°C. By setting a starting temperature of
52 10°C (which is more realistic for a shallow aquifer in temperate climate), this critical density
53 difference is already reached at a heating threshold of 8°C. This value coincides with that by



54 (Ma et al., 2012), who refined their previous findings using field experiments and numerical
55 sensitivity analysis. Essentially, despite several appealing properties, such a tight range for
56 the temperature limits the viability of heat as a tracer. Viscosity and buoyancy effects may
57 render a reliable interpretation of thermal tracer tests impossible. Alternatively, techniques
58 have been developed that can handle broader ranges and are not prone to hydraulic effects of
59 temperature variation. This is the focus of our study.

60 Our starting point is the fact that for detecting preferential flow paths full analysis of thermal
61 transport behavior may not be necessary. If we focus on characteristic parameters such as
62 travel times or moments of the BTCs, the signal-to-noise ratio may be acceptable for much
63 broader temperature ranges. Travel times of traditional solute tracers are related to the
64 hydraulic properties of aquifers, assuming that the main transport process is advection. This
65 is the case given a sufficient ambient hydraulic or forced gradient during the experiment
66 (Doro et al., 2014; Saar, 2011). One important difference of heat tracer transport over
67 traditional tracers is that diffusion not only takes place in the pore fluid, but in the rock matrix
68 as well. So while the tracer front of a solute tracer tends to be sharp, the thermal tracer front
69 appears smoothed. This may make interpretation of BTCs more difficult.

70 Because thermal diffusion takes place, heat transport is affected not only by the hydraulic
71 properties but by the thermal properties of the aquifer material as well. However, contrasts in
72 thermal parameters are relatively small compared to contrasts in K , which typically spans
73 orders of magnitude (Stauffer et al., 2013). Porosity can also be influential in heat transport,
74 due to the high heat capacity contrast of water and rock components. Yet, natural variability
75 in porosity is commonly much smaller than that in K . Therefore variability observed in the
76 transport of a thermal tracer is caused mainly by heterogeneity of K .



77 In previous studies on thermal tracer testing, diverse set-ups have been chosen that differ with
78 respect to heating method, injection volumes, rates and temperatures, test duration, and well
79 configurations (Wagner et al., 2014). Mostly hot water is infiltrated in an injection well and
80 BTCs are recorded in one or more downstream observation well (Ma et al., 2012; Macfarlane
81 et al., 2002; Palmer et al., 1992; Read et al., 2013; Wagner et al., 2014; Wildemeersch et al.,
82 2014). Insight in aquifer heterogeneity is not well constrained by analysis of thermal signals
83 introduced and measured over long screens. To obtain a better definition of the heterogeneity,
84 observations in several wells or at different depth levels need to be compared. Ideally a
85 tomographic setup is chosen, where multiple point injection (sources) and observation points
86 (receivers) are used. By combined inversion of all signals, the spatial variations in K are
87 reconstructed. So far, however, this concept is more established in geophysics, and for
88 aquifer characterization in hydraulic tomography, which utilizes pressure signals from depth-
89 dependent pumping tests or multi-level slug tests (Cardiff et al., 2009, 2012; Illman et al.,
90 2010; Yeh and Zhu, 2007).

91 (Klepikova et al., 2014) presented a passive thermal tracer tomography application for
92 characterizing preferential flow paths in fractured media. Their method focused on modelling
93 the fracture network with a sequential method which involves first identifying the location of
94 fault zones on the temperature-depth profiles under ambient flow and pumping conditions.
95 Next, an inversion of the temperature profiles is conducted to obtain borehole flow profiles,
96 and the last step is to estimate the hydraulic properties from these flow profiles. This method
97 provides cross-well connectivities. The work by (Doro et al., 2014) is dedicated to the
98 experimental design of cross-well forced gradient thermal tracer tomography. In their
99 approach, a special multi-level injection system is necessary to induce the tracer into a
100 horizontal layer. They also recommend limiting the temperature range to avoid buoyancy
101 effects. Their proposed methodology to interpret the results is to use an inversion scheme



102 developed by (Schwede et al., 2014) for this specific experimental setup. This inversion
103 method utilizes the temporal moment of measured BTCs and hydraulic head data together in
104 a joint geostatistical inversion procedure (Illman et al., 2010; Yeh and Zhu, 2007; Zhu et al.,
105 2009). This procedure is computationally demanding, and it assumes a multi-Gaussian
106 distribution of hydraulic properties, which represents a strong restriction in comparison to the
107 true conditions in the field.

108 In our work, we suggest a travel-time based inversion procedure, which does not require a
109 priori structural or geostatistical assumptions and is computationally efficient. It is motivated
110 by (Vasco and Datta-Gupta, 1999), who presented a numerical approach to reconstruct the
111 hydraulic parameters of an aquifer using solute tracer injections in a tomographical setup. As
112 a core element, the transport equation is transformed into an eikonal problem using an
113 asymptotic approach for the tracer transport solution. Their approximation uses the similarity
114 of tracer front propagation to seismic and electromagnetic waves, but with the restriction that
115 the tracer front is abrupt. This approximation can be used for hydraulic signals as well (Vasco
116 et al., 2000), and the travel time of the hydraulic signal can be related to the hydraulic
117 diffusivity of the system. (Brauchler, 2003) further developed a travel time based inversion
118 for Dirac and Heaviside hydraulic sources, using the early time diagnostics of the signals. To
119 improve spatial resolution, they applied staggered grids (Vesnaver and Böhm, 2000) during
120 inversion. This inversion methodology was applied to several hydraulic laboratory and field
121 experiments (Brauchler et al., 2007, 2011, 2013b; Hu et al., 2011; Jiménez et al., 2013).
122 (Brauchler et al., 2013a) also utilized travel-times in a tracer experiment on rock samples in
123 the laboratory scale. Their work revealed that for those samples transport was dominated by
124 the rock matrix, but hydraulic parameters were not estimated.

125 In this study, we present a new formulation for inversion of spatially distributed hydraulic
126 conductivity using early tracer travel times. It follows the same principles as presented by



127 (Brauchler, 2003) for hydraulic tomography. Our objective is to obtain a versatile and
128 efficient technique for thermal tracer tomography, which, by focusing on early times,
129 minimizes the role of buoyancy and viscosity effects. In the following section, the new
130 inversion procedure is introduced. It is then applied to a three-dimensional (3-D) high
131 resolution aquifer analog of the Guarani aquifer in Brazil. We inspect the capability of the
132 new approach to reconstruct 2-D and 3-D sections with heterogeneous K -distribution, and
133 provide a sensitivity analysis of variable injection rates and temperature ranges. Finally, the
134 findings of exhaustive testing with variable field conditions and technical design parameters
135 are compiled to determine the application window of this new thermal tomography variant.

136 **2 Tomographical inversion procedure**

137 **2.1 Travel time inversion**

138 Under high Péclet number conditions, when it can be assumed that the thermal transport is
139 dominated by advection, the propagation of an injected thermal plume can be used to gain
140 information about the hydraulic properties of the investigated aquifer. Our goal is to calculate
141 the hydraulic conductivity, K , of the aquifer by inverting the advective thermal tracer
142 breakthrough times. (Vasco and Datta-Gupta, 1999) showed that the transport equation of a
143 solute tracer can be formulated as an eikonal equation, which is utilized to calculate K .
144 According to this work, a line integral can be written for tracer breakthrough times:

$$t_{st}(x_r) = \int_{x_s}^{x_r} \frac{ds}{v_{st}(s)} = \int_{x_s}^{x_r} \frac{\phi(s)}{K(s)i(s)} ds \quad (1)$$

145 Here, $t_{st}(x_r)$ is the breakthrough time of the solute tracer at the receiver (x_r), x_s is the source
146 location, v_{st} is the mean tracer velocity, ϕ is the aquifer porosity and i is the local hydraulic



147 gradient. This equation can be used for a thermal tracer (tt) by including the thermal
148 retardation factor, R :

$$t_{tt}(x_r) = \int_{x_s}^{x_r} \frac{ds}{v_{tt}(s)} = \int_{x_s}^{x_r} \frac{\phi(s)}{R K(s) i(s)} ds \quad (2)$$

149 Thermal retardation depends on the porosity of the aquifer, ϕ , the heat capacities of aquifer
150 matrix, C_m , and of water C_w :

$$R = \frac{C_m}{\phi C_w} \quad (3)$$

151 Changes in these parameters are commonly small compared to changes in K , thus the thermal
152 retardation can be approximated as a constant. For the same reason, ϕ and the hydraulic
153 gradient, i , are also considered fixed. Values of ϕ and C can be approximated from prior data,
154 while the hydraulic gradient between observation and injection is measured during the
155 experiment. With these assumptions and the use of standard tomography algorithms, a
156 solution can be found on a pre-defined grid.

157 The presented method uses a step-function injection temperature signal for the active thermal
158 tracer test. In this case the traveling time of the thermal tracer is associated with the
159 propagating thermal front. The tomographical concept requires multiple independent thermal
160 tracer injections at different depths. Temperature BTCs are recorded at multiple observation
161 points, for example at different levels in a downgradient observation well. As common
162 practice for such setups, the number of sources and receivers is one of the important factors
163 that defines the significance and resolution of the results.



164 2.2 Early time diagnostics

165 Compared to a conservative solute tracer, heat does not behave ideally. Diffusion is
 166 significant in aquifer matrix and pore fluid, while the viscosity and density of the
 167 groundwater are variable. Due to the highly diffusive behavior, the emerging thermal front
 168 cannot be considered as a sharp transition boundary. In order to obtain accurate results with
 169 the inversion, the complications from thermal diffusion need to be mitigated. Both diffusion
 170 and mechanical dispersion effects increase with travel time. Mitigation thus can be done by
 171 using an earlier characteristic time of the thermal front instead of the (peak) breakthrough
 172 time. The earlier characteristic time can then be corrected to the real breakthrough time using
 173 a conversion factor, as shown for hydraulic tomography by (Brauchler, 2003) with a
 174 correction for the specific storage coefficient.

175 The propagation of a thermal front far from the source is described as a one-dimensional (1-
 176 D) advection-diffusion problem considering thermal retardation.

$$R \frac{\partial T}{\partial t} = D \frac{\partial^2 T}{\partial x^2} - u \frac{\partial T}{\partial x} \quad (4)$$

177 The analytical solution to this problem is (Ogata and Banks, 1961):

$$T(x, t) = T_0 \left(\frac{1}{2} \operatorname{erfc} \left(\frac{Rx - ut}{2\sqrt{DRt}} \right) + \frac{1}{2} \exp \left(\frac{ux}{D} \right) \operatorname{erfc} \left(\frac{Rx + ut}{2\sqrt{DRt}} \right) \right) \quad (5)$$

178 The breakthrough time of the thermal tracer is associated with the peak of the first derivative
 179 of the temperature, and can be calculated analytically. During the breakthrough detection,
 180 instead of the temperature, the first derivative, T' , of the temperature is used as the observed
 181 signal.

$$t_{peak} = \frac{R\sqrt{9D^2 + u^2x^2} - 3DR}{u^2} \quad (6)$$

182 Early time characteristic values can be described proportionally to the peak value:



$$T'(x, t) = \alpha T'(x, t_{peak}) \quad (7)$$

183 with

$$\alpha = \frac{T'(x, t)}{T'(x, t_{peak})} = \frac{T'(x, \tau_{\alpha} t_{peak})}{T'(x, t_{peak})} \quad (8)$$

184 Substituting the peak time solution into this expression yields:

$$\alpha = \frac{\exp\left(-\frac{(\tau_{\alpha} - 1)(u^2 x^2 - 18D^2 \tau_{\alpha} - \tau_{\alpha} u^2 x^2 + 6D \tau_{\alpha} \sqrt{9D^2 + u^2 x^2})}{12D^2 \tau_{\alpha} - 4D \tau_{\alpha} \sqrt{9D^2 + u^2 x^2}}\right)}{\tau_{\alpha}^{\frac{3}{2}}} \quad (9)$$

185 where

$$\tau_{\alpha} = \frac{t}{t_{peak}} = \frac{1}{f_{\alpha}} \quad (10)$$

186 and τ_{α} is the relative time to the peak time, and f_{α} is the transformation factor.

187 Although Eq. (9) has three additional parameters, velocity (u), distance (x) and dispersion

188 coefficient (D), the function is not sensitive to these values, and under realistic conditions its

189 shape remains the same. Neglecting the second order terms of velocity, the expression can be

190 simplified to

$$\alpha = \frac{\exp\left(\frac{6(\tau_{\alpha} - 1)}{4\tau_{\alpha}}\right)}{\tau_{\alpha}^{\frac{3}{2}}} \quad (11)$$

191 This equation can be solved analytically for τ_{α} , although infinite numbers of transcendent

192 solutions exist (the presented solution is only valid if α and f_{α} are positive). To have an

193 analytical solution for τ_{α} -values between 0 and 1 (times before the peak time), the -1st branch

194 of the Lambert Omega function is applied. The final expression for the transformation factor

195 reads



$$f_{\alpha} = -\text{LambertW}\left(-1, -\frac{\alpha^2}{e}\right) \quad (12)$$

196 This finding corresponds with the transformation factor used in hydraulic tomography
197 presented by (Brauchler, 2003; Hu et al., 2011). In order to apply the conversion, the
198 temporal scale of the record must be adjusted to the time of the thermal front arrival. In
199 practice, this time is when the first increase on the temperature derivative record can be
200 observed.

201 The application of early time diagnostics is illustrated on Figure 1. We are mainly interested
202 in advective transport. However, thermal diffusion may also be significant, smoothening and
203 expanding recorded temperature BTCs, and thus affecting also its derivative. The
204 identification of the peak time through the derivative T' is challenging due to the flatness of
205 the curve at the maximum value of the peak. However, using the early time diagnostics (step
206 1), only the value of the peak must be known for Eq. (7). In step 2, the desired fraction of the
207 peak value (α) and the associated time ($\tau_{\alpha} t_{peak}$) must be found on the measured T' curve.
208 Finally, in step 3, the time is corrected to a calculated peak time using the transformation
209 factor according to Eq. (12). In this step, the temperature curve is extrapolated from the
210 fraction time, and by this the effect of diffusion is taken into account. Note that the time zero
211 of the correction is when the thermal front reaches the receiver. This time can practically be
212 chosen when the earliest identifiable temperature change appears at a receiver.

213 **2.3 Staggered grids and null-space energy**

214 To invert the tracer travel times the SIRT algorithm (simultaneous iterative reconstruction
215 technique) is used to solve the eikonal problem, implemented in GeoTOM3D (Jackson and
216 Tweeton, 1996). The algorithm calculates the transport trajectories between the sources and
217 receivers and solves the line integral of Eq. (2) along the trajectories – in a curve based 1-D



218 coordinate system. For solving the line integral, the solution domain is discretized to a grid.
219 Initially a homogeneous velocity field is defined, and then the velocity values of the cells are
220 updated iteratively to minimize the difference between the inverted and recorded travel times.
221 In order to provide the uniqueness of the solution, an even-determined problem is needed and
222 thus the number of grid cells should be kept close to the number of measurements (source-
223 receiver combinations). The spatial distribution of the trajectories is never uniform over the
224 domain, the result quality can differ in space, and the result can be non-unique.

225 For discretization, instead of constructing a static regular grid, the staggered grid method
226 (Vesnaver and Böhm, 2000) was used. Solving the problem on a regular grid would highly
227 constrain the freedom of the solution to the geometry of the used grid and the source-receiver
228 locations. By applying the staggered grid method this constrain can be overcome, with the
229 benefit that the nominal spatial resolution is increased. Otherwise, for a good spatial
230 resolution using one fine grid, a large number of sources and receivers would be required or
231 regularization terms would have to be applied. Staggered grids were successfully employed
232 for hydraulic tomography by (Brauchler, 2003) and for solute tracer tomography by
233 (Brauchler et al., 2013a). In this staggered variant, the problem is solved on different
234 vertically and horizontally shifted versions of a low-resolution regular grid. The inverted
235 results are different for the shifted grids, which is exploited by arithmetically averaging these
236 results to arrive at a final tomogram. The inversion will be stable because of the coarse grids,
237 while the resolution of the averaged tomogram will be as small as the displacements.
238 Although this means that the travel time inversion step will be performed multiple times for
239 one tomogram, it is still computationally affordable due to the marginal computation demand
240 of a single coarse grid resolution.

241 To characterize the reliability of the results, the null-space energy map is computed. This
242 method has been applied for hydraulic tomography in several studies (Brauchler et al., 2013a,



243 2013b; Jiménez et al., 2013) and uses the distribution of the inverted transport paths over the
244 inversion grid. The null space energy map is calculated from the singular value
245 decomposition (SVD) of the tomographic matrix, which contains the length of each inverted
246 transport path in each grid cell. Values of the null space energy map are between 0 and 1,
247 thus higher values mean higher uncertainties. Based on the null-space energy map, non-
248 reliable pixels can be deleted from the tomogram. The resulting full inversion procedure,
249 starting with the tracer data and ending with the reliable part of the final K -tomogram is
250 depicted in Figure 2.

251 **3 Application case**

252 **3.1 Aquifer analog model**

253 The presented methodology is developed and tested on the Descalvado aquifer analog
254 (Höyng et al., 2014) that is implemented in a finite element heat transport model (Figure 3).
255 This analog represents a 3-D high-resolution dataset obtained from mapping an outcrop of
256 unconsolidated fluvio-aeolian sediments in Brazil. These sediments host parts of the Guarani
257 aquifer system, one of the world's largest groundwater reservoirs. The analog is based on five
258 vertical outcrop sections that are recorded during ongoing excavation and interpolated by
259 multi-point geostatistics following the procedure by (Comunian et al., 2011). The spatial
260 extent of the analog is $28 \text{ m} \times 7 \text{ m} \times 5.8 \text{ m}$ (x, y, z). Hydraulic conductivity, K , and porosity,
261 ϕ , data was documented on sub-decimeter scale, in three parallel and two perpendicular
262 profiles during excavation. (Höyng et al., 2014) distinguish nine different hydrofacies (H1-9)
263 of similar hydraulic properties, which form the primary building blocks and which determine
264 the structural heterogeneity of the characterized volume. In order to ease the interpretation of
265 results, the focus is on major architectural elements which are the four zones that form the



266 characteristic layers of the formation (Table 1). These can be easily distinguished visually by
267 the dominant color in the selected color scale in Figure 3: the blue top low conductive zone,
268 the red central conductive zone, the orange lower-central zone and the yellow bottom zone. In
269 order to use this analog for thermal transport simulations, the original dataset (Höyng et al.,
270 2014) is expanded with estimated thermal properties (heat capacity, C , thermal conductivity,
271 λ) assigned to the different hydrofacies units (“thermofacies”). These properties were
272 calculated based on porosity and available lithological information (Bayer et al., 2015).

273 The Descalvado aquifer is built up mainly by highly conductive sand and gravel with a
274 layered structure. The average hydraulic conductivity value is approximately $K = 10^{-4} \text{ ms}^{-1}$
275 and the largest difference between two adjacent hydrofacies is three orders of magnitude.
276 Locally, low- K clay intraclasts exist that even induce higher variations. But, due to sizes of
277 only a few centimeters and a marginal volumetric share, they are negligible for flow and
278 thermal transport simulation. Thermal heterogeneity among the different facies units is
279 controlled by differences in porosity, because the mineral composition does not substantially
280 vary. When clay intraclasts are ignored, thermal conductivity spans from $\lambda = 2.6$ to 3.2 Wm^{-1}
281 s^{-1} , and the volumetric heat capacity ranges between $C = 2.4$ and $2.6 \text{ MJ m}^{-3}\text{K}^{-1}$. The global
282 thermal isotropic micro-dispersivity in the model is set to about the average grain size, $\beta =$
283 0.1 mm .

284 Flow and transport is simulated as coupled processes using the software FEFLOW (Diersch,
285 2014), and the SAMG algebraic multigrid solver (Thum and Stüben, 2012). The analog is
286 embedded into a larger domain with extrapolated homogeneous layers, to minimize lateral
287 boundary effects. The model mesh is generated with the Triangle algorithm (Shewchuk,
288 1996) and progressively refined towards the analog. Close to wells, the elements are refined
289 to mm scale. The total extent of the model is $118 \text{ m} \times 117 \text{ m} \times 15.7 \text{ m}$ consisting in a total of



290 1'664'626 triangle prism elements. In the center of the model, the resolution of the finite
291 element mesh is similar to or finer than the resolution of the original aquifer analog dataset.

292 The aquifer is assumed to be confined. In order to simulate initial steady-state conditions with
293 regional groundwater flow in the direction of the long axis, x , constant head boundary
294 conditions are imposed at the perpendicular sides of the model and no-flow conditions at the
295 other model faces. The constant head values are specified to impose an average hydraulic
296 gradient according to Table 2. The initial temperature of the model is set to 10 °C. This value
297 is also used as a boundary condition at the sides of the model, which yields isothermal initial
298 conditions.

299 **3.2 Experimental setup**

300 We present reconstructions of K -fields of 2-D and 3-D analog sections. These sections are
301 called tomograms. 2-D profiles represent vertical cross-sections between an injection (source)
302 and an observation (receiver) well, while data of three observation wells is utilized for 3-D
303 reconstruction. We specify a base case, which serves as our principal study case, and
304 additionally inspect the performance of the methodology by varying the experimental design
305 and profile. Note that independent of the dimensions of reconstructed sections, always the
306 full 3-D analog model was used to simulate the thermal tracer propagation and resulting
307 travel times.

308 Focus is set first on 2-D reconstruction. Three profiles in the central plane of the aquifer are
309 selected (Figure 3). This central plane constitutes a mapped outcrop section with relatively
310 high facies variability. It contains heterogeneous structures of different sizes and contrasts,
311 and it is chosen for being sufficiently far away from the analog boundaries. The location of
312 profile 1 is depicted in Figure 3. Figure 4a shows the relative locations of an upstream



313 injection well and downstream observation well used for all three 2-D profiles. The distance
314 between the wells is 5 m for an investigated area of 5 m x 6 m.

315 To examine further the role of aquifer heterogeneity, two additional profiles from the central
316 plane of the analog are investigated. In both cases, the source-receiver geometries are kept the
317 same (Figure 4a). Profile 2 shows a similar layered structure as profile 1, but with less small-
318 scale heterogeneities. The central conductive zone is thicker, providing better connection
319 between the two wells. In profile 3, the central conductive zone is discontinuous, creating a
320 different hydrogeological situation, with weaker connection between the two wells.

321 In the simulated setup, 6 sources and 6 receivers are employed (Figure 4a), resulting in a set
322 of 36 source-receiver combinations. The sources are defined as point injections with constant
323 injection rates during the entire simulation time. The used injection temperature signal
324 delineates a Heaviside step function, where the instantaneous change in temperature is
325 arbitrarily set at 0.1 days after the start of simulation, which marks the beginning of the
326 experiment. In order to record BTCs in all observation points even at very small injection
327 rates and temperatures, extremely long simulation times are used (50 days). However, most
328 of the breakthroughs occur during the first five days of the simulation.

329 The crucial technical design parameters for the experiments are the injection rate, Q , and the
330 injection temperature (or temperature difference, ΔT , in comparison to ambient aquifer
331 conditions). The base values of these two parameters are selected after preliminary field
332 testing (Schweingruber et al., 2015) as $Q = 1 \text{ ls}^{-1}$, and $\Delta T = 20 \text{ }^\circ\text{C}$. These parameter values
333 and hydraulic model settings are varied in the ranges listed in Table 2 in the sensitivity
334 analysis presented in Section 4.3.

335 In practice, the source of the injected water can be the investigated aquifer, but note that in
336 this case heating has to be well controlled to keep the injection temperature constant. During



337 a field experiment, the recorded data is always distorted by noise. With the commonly used
338 temperature sensors, this noise is considered very small (Wagner et al., 2014), but still the
339 sensitivity of the temperature sensors is limited. To take this into account when simulating
340 the receiver points, those where the temperature changes are smaller than 0.1°C are ignored
341 for the inversion. In addition, source-receiver combinations with geometric angles larger than
342 40° were not used, following the suggestion of (Hu et al., 2011) for hydraulic tomography in
343 layered aquifers.

344 For the 3-D reconstruction, an exemplary case is defined with 1 injection and 3 observation
345 wells forming a triangular prism (Figure 4b) located close to profile 1. The base face is an
346 isosceles triangle, and the observation wells are located along the baseline. The axis of this
347 triangle is at the line where the 2-D profiles are located. The distance between the injection
348 well and the central observation well is 6.5 m and the length of the triangle base is 3 m. The
349 configuration of the individual wells is the same, resulting 18 observation points, and 108
350 source-receiver combinations in total. The experiment was simulated using the base values
351 from Table 2, employing the same Heaviside injection signals as at the 2-D cases.

352 **4 Results and discussion**

353 The following, results are structured into four major parts. The first part is the inspection of
354 the inverted tomograms for the three 2-D and one 3-D analog profiles. The second part is the
355 validation of the method using the result of the 3-D reconstruction. The third part is a
356 sensitivity analysis of the inversion procedure with respect to experimental settings such as
357 injection rate and temperature. The fourth part reveals the application window of travel-time
358 based thermal tomography through rigorous testing with different sections, changing
359 hydraulic conductivity contrasts and varying experimental parameters.



360 **4.1 Reconstruction of hydraulic conductivity profiles**

361 The left column of Figure 5 depicts the analog profiles, and these are contrasted with the
362 inverted ones on the right. For better comparability, the original analogs are up-scaled to the
363 same grid as used for the results with $0.125 \text{ m} \times 0.125 \text{ m}$ cell size. Figure 5a represents the K
364 distribution of the aquifer analog at profile 1. It is characterized by an overall layered
365 structure, and it shows highest variability with small scale facies patches in the central part
366 between $z = 2 \text{ m}$ and $z = 4.5 \text{ m}$. Of major interest is the red central conductive zone
367 (hydrofacies H4) at around $z = 4 \text{ m}$ with non-uniform thickness. In the field, it can cause flow
368 focusing and promote preferential flow. This zone is even more pronounced in profile 2
369 (Figure 5c), but not continuous in profile 3, where only laterally high-conductive wedges can
370 be found. In all profiles, the underlying lower central zone is dominated by the orange facies
371 H5. With the embedded small-scale layered and cross-bedded elements, this zone will give
372 insight in the competence of the inversion procedure to resolve local, decimeter-scale
373 structures.

374 BTCs from 36 source-receiver combinations are recorded for one tomographic experiment,
375 but only 34 of them are used for the inversion, since the applied 40° angle criteria between
376 sources and receivers excluded two combinations. During staggering, the tomographic
377 inversion is performed on 16 different spatially shifted coarse grids. The uniform cell size of
378 these low resolution grids is $0.5 \text{ m} \times 0.5 \text{ m}$. In total, 30 iterations are done per inversion, and
379 the inverted velocities are restricted within a range of physically possible tracer velocities.
380 Note that the inversion algorithm allows to provide constrains in velocity and if they are not
381 set appropriately, it can produce outlier pixels close to the sources and receivers, where the
382 flow is focused. Velocity limits (i.e. expected high and low values for K and i) can be
383 calculated using prior information, and the method is not sensitive to small changes in their
384 values. The 16 coarse tomograms are merged together into a fine staggered grid, with a



385 resolution of 0.125 m x 0.125 m. The total computational time for reconstructing one profile
386 was around 10 minutes on an office PC (Intel® Core™ i7-4770 CPU 3.40 GHz).

387 After calculation of null-space energy maps, a threshold of 85 % is found suitable to
388 constrain the *K*-tomograms. In other words, only with null space energy of less than 85 % (or
389 vice versa, with a reliability of at least 15 %) pixels are shown in the final reconstructed
390 profile. As illustrated in Figure 5b, d, and f, this yields fringed edges in the *K*-tomograms,
391 and some grayed gaps in the interior. Since the null space denotes local coverage of transport
392 trajectories, there are some regions which are unsatisfactorily accessed. As expected, these
393 are mainly close to the boundaries of the inspected profile and not in the reach of the source-
394 receiver couples. By changing the arbitrary null space energy threshold, masking of areas of
395 low reliability may be accentuated or mitigated. The most suitable value of the threshold,
396 however, is based on expert knowledge and is set depending on the requirements of the
397 specific case. Experience shows that modifying this value (by 5-10 %) has a minor influence
398 on the visualized structures of major interest, because the null-space energy of the highly
399 conductive zones tends to be very small.

400 The reconstructed profiles in the right column of Figure 5 shed a first light on the capabilities
401 of thermal tomography. First, we observe that for all profiles, the upper zone (in blue) cannot
402 be reconstructed by the inversion. Typically a considerable fraction of it is masked in gray
403 due to the limited contribution to heat transport, which is not surprising due to the low
404 hydraulic conductivity of this zone. In contrast, the tomographic approach identifies the
405 location of the highly-conductive upper central zone (in red) rather well. This zone delineates
406 the fastest travel route between the wells for the heat tracer. Between the upper (blue) and
407 central (red) zones is the strongest contrast in the profiles. This strong contrast shadows the
408 top of the tomograms, because the transport is short-circuited through the high-*K* zone,
409 resulting that it appears upshifted on the tomogram. When the contrast is smaller, such as in



410 profile 3, this shadow effect is weaker, and it is possible to gain better insight into the low
411 conductivity zone (Figure 5e-f).

412 A striking feature is that the tomographic approach resolves the continuity of the highly-
413 conductive upper central zone in profiles 1 and 2, and it detects the discontinuity in profile 3.
414 Furthermore, the inverted value of hydraulic conductivity of this zone ($K = 8 \times 10^{-4} \text{ ms}^{-1}$) is
415 comparable to the original model ($K = 1.38 \times 10^{-3} \text{ ms}^{-1}$). For the lower central zone, we obtain
416 a similarly good match with an inverted value of $1.6 \times 10^{-4} \text{ ms}^{-1}$ in comparison to the original
417 value of $2.96 \times 10^{-4} \text{ ms}^{-1}$ for the dominant hydrofacies H5 (Table 1). This is remarkable,
418 having in mind that related travel time-based techniques of hydraulic tomography have
419 shown to be suited for structural reconstruction, but to a lesser extent for hydraulic parameter
420 estimation. In those studies, parameter values were obtained by ex-poste calibration with the
421 full forward model [Hu *et al.*, 2011; Jiménez *et al.*, 2013; Hu *et al.*, 2015].

422 The promising findings as depicted in Figure 5 support the applicability of travel-time based
423 tracer inversion for thermal tomography, even though thermal diffusion tends to blur
424 advective travel times, which hinders a reliable inversion. However, by taking early arrival
425 times of the recorded BTCs, this effect is minimized. Likewise, when preferential pathways
426 exist, these will be detected by the first thermal breakthrough, which is least influenced by
427 diffusion. As a result, travel-time based thermal tomography appears especially suited for
428 locating and characterizing high-conductivity zones.

429 With the 36 source-receiver combinations, exact profile reconstruction is not possible, since
430 the tomograms appear to be smoothed. Fine-scale differences in the form of the high-
431 conductivity zone are not reproduced in the tomograms. This is the same for the small facies
432 mosaics that originally occur in the mainly orange lower central zone. This zone seems mixed
433 with the lower yellow zone, and the hydraulic conductivities of both zones are slightly



434 underestimated. Despite the minor hydraulic contrast between both layers, however, the
435 tomograms indicate locally a facies transition (especially in Figure 5f). This is not identified
436 in the tomogram of profile 1 (Figure 5b). Here most small-scale structures exist in the lower
437 central part above. These cannot be resolved, but they detract from the transport routes of the
438 thermal tracer and thus induce noise in the reconstructions of the lower central and bottom
439 layer.

440 Figure 6 shows the reconstruction of the selected 3-D section. The result is presented the
441 same way as the 2-D profiles, using an upscaled version of the original analog for
442 comparison. 3-D staggering is employed resulting in 64 coarse grids in total. This requires 64
443 individual inversions and thus a computational time that is drastically longer than in the 2-D
444 cases. With 20 iterations per inversion, the total computational time on the same PC (Intel®
445 Core™ i7-4770 CPU 3.40 GHz) was around 1 hour for 3-D inversion. The spatial resolution
446 of the coarse grid is $0.5 \text{ m} \times 0.5 \text{ m} \times 0.5 \text{ m}$ and of the staggered grid thus is $0.125 \text{ m} \times 0.125$
447 $\text{m} \times 0.125 \text{ m}$.

448 To assess the reliability of the inverted result, the null-space energy map is calculated. For the
449 3-D application a limit of 95% of reliability is used to accept reconstructed voxels. Lower
450 values would substantially reduce the reconstructed volume, since non-reliable voxels are not
451 presented. Generally, the reliability and thus overall result quality of the 3-D analysis is
452 worse than for the 2-D cases. This is due to the fact that the inverted transport paths cover
453 less of the domain of interest.

454 Figure 6a depicts the upscaled analog model, sliced to half at the central plane where the
455 injection well is located. The same way of presentation is used for the reconstruction in
456 Figure 6b. To highlight the differences to the 2-D results, the inverted high- K zone is
457 presented for the whole domain without slicing it to half. The central slice of the 3-D



458 reconstruction is similar to profile 1, because the injection well located at the same location
459 (Figure 5a) and the observation wells located only 1.5 m further away. However, by blanking
460 unreliable voxels in the 3-D visualization it is difficult to compare the 2-D and 3-D
461 reconstruction in Figures 5b and 6b. At first sight, the reconstructed features of the 3-D and
462 the 2-D inversion are similar. A pixel-to pixel comparison using the central plane of the 3-D
463 reconstruction shows that the difference to the reconstructed values of profile 1 is less than 30
464 %. This demonstrates that especially for systems with mainly horizontal structures such as the
465 sedimentary aquifer here, results in 2-D are only slightly improved in a 3-D inversion.
466 Comparing the full profile, the inverted K values are lower than at the 2-D cases, but still in
467 the same magnitudes as the original values of the aquifer analog (central conductive zone:
468 $3 \times 10^{-3} \text{ ms}^{-1}$ inverted to $1.4 \times 10^{-3} \text{ ms}^{-1}$ original, middle zone: $1 \times 10^{-4} \text{ ms}^{-1}$ inverted to 3×10^{-4}
469 ms^{-1} original).

470 In Figure 6b, the central conductive zone of the aquifer is localized mainly at the lateral
471 boundaries close to the wells. Centrally, K values are underestimated and smooth channels
472 appear between injection and observation wells, delineating the suspected main transport
473 paths of the tracer. Similar as in the 2-D reconstructions, the central part of these channels is
474 vertically upshifted. The top low- K zone is not reconstructed, but fragments of it appear in
475 the results, marking the location of the contrast boundary on the bottom of this zone. The
476 contrast between the two lower zones can be identified laterally but not centrally – same as in
477 the 2-D profile.

478 4.2 Validation

479 For validation, the reconstructed 3-D K -field is implemented in a numerical model with the
480 same settings as that used for the forward simulations with the original analog data. Here,
481 homogeneous thermal properties are assumed. In total 9 observation wells with 6 observation



482 points in each are used to validate the inverted result (Figure 4b). A full tomographic
483 experiment is simulated with 6 independent warm water injections using the same
484 configuration as the original simulated experiment. The recorded BTCs are compared with
485 simulations with the aquifer analog dataset. The differences in the breakthrough times are
486 used for the validation.

487 Considering the good reconstruction of the high- K zone, which is most relevant for the
488 thermal transport, we can expect that at most of the observation points the difference would
489 be small. This is exactly what Figure 7 shows, where the distribution of the differences is
490 presented as histogram. Most of the values are close to zero showing a good validation of the
491 result. There are two groups of outliers marked with yellow color. The negative outliers are
492 associated with the observations in the top low- K zone where the inversion was not sufficient.
493 Here the predicted heat transport is faster than in the aquifer analog. The second outlier group
494 is related to the underestimated K of the lower central zone (Figure 3). The difference in the
495 breakthrough times becomes most significant at observation points that are furthest from the
496 injection well.

497

498 **4.3 Sensitivity analysis**

499 The experimental setup may be crucial for the quality of the inversion results. For example, it
500 is well known from related tomographic inversion studies that the feasible resolution depends
501 on arrangement and the numbers of sources and receivers (Cardiff et al., 2013; Paradis et al.,
502 2015). Here we focus on two technical design parameters, which are particularly crucial for
503 thermal tomography when using heated water: the injection temperature and the injection
504 rate. In the following sensitivity analysis, we question whether these need to be carefully



505 tuned or not. Profile 1 is chosen for investigation depicted again in Figure 8a and 9a. Note
506 that for forward simulation of travel times, always the full 3-D analog model is used.

507 We first inspect the role of the temperature of the injected water. In all of our models, the
508 ambient groundwater temperature is considered uniform and 10 °C. Viscosity and density
509 effects increase with the temperature difference, ΔT , in comparison to the ambient
510 groundwater. These effects may distort the results of inversion, and thus a maximal
511 difference of $\Delta T = 8\text{-}15$ °C has been suggested for thermal tracer testing (Doro et al., 2014;
512 Ma and Zheng, 2010; Russo and Taddia, 2010). This severely constrains the applicability of
513 heat as an active tracer, because it complicates interpretation of BTCs influenced by
514 buoyancy forces. For our tomography, we examine a ΔT from 5°C to 80°C to cover the full
515 range of technical possibilities. The injection rate is kept at $Q = 1 \text{ ls}^{-1}$.

516 Figure 6 depicts the inverted K -tomograms for $\Delta T = 5, 10, 20, 40$ and 80 °C. The results
517 show that the inversion method is not very sensitive to ΔT . The tomograms slightly vary, but
518 they all maintain the major features, and especially the central high conductive zone is
519 identified similarly in all variations. Even with an extreme value of $\Delta T = 80$ °C, no distortion
520 appears. This is surprising because buoyancy effects are significant under such conditions.
521 This is attributed to the use of early time diagnostics, which are mainly controlled by
522 advective transport even if substantial thermal and density gradients prevail in the aquifer.
523 Being able to inject water with high temperature is considered advantageous, because this
524 means that a strong signal is introduced, a high signal-to-noise ratio can be achieved, and a
525 greater aquifer volume can be accessed. In practice, of course, maintaining a constant
526 injection temperature at high temperatures can be a technical challenge, and requires more
527 sizeable heating devices.



528 The sensitivity of the injection rate, Q , is investigated on a range of four orders of magnitude,
529 $Q = 10^{-3}$, 10^{-2} , 10^{-1} , 1 and 10 ls^{-1} (Figure 7). The injection temperature is fixed at $\Delta T = 20 \text{ }^\circ\text{C}$.
530 At small injection rates, the heat introduced to the aquifer is small; hence there is no
531 detectable breakthrough at most of the observation points. As shown in Figure 9b, little
532 insight is obtained with $Q = 10^{-3} \text{ ls}^{-1}$, and the quality of the results is poor. Increasing the
533 injection temperature can improve the quality of the result in this case.

534 By raising the injection rate, the reconstructed continuity of the central conductive zone
535 improves (Figure 9c-e). For our particular case, this is attributed to the setup. Since the top
536 two observation points are located in the upper low conductivity zone, this influences the
537 reconstruction of the central high conductive zone.

538 In contrast, at the highest simulated injection rate of $Q = 10 \text{ ls}^{-1}$, the derived tomogram is
539 unsatisfactory (Figure 9f). This is caused by the highly distorted flow field. Our inversion
540 procedure is based on the assumption that the hydraulic gradient between the two wells is
541 (approximately) constant. This is not valid anymore, and the relation between inverted mean
542 tracer velocity and hydraulic conductivity is not linear. This effect appears only at very high
543 injection rates, in this case at $Q = 10 \text{ ls}^{-1}$, which exceeds technical possibilities (with an
544 injection temperature of $\Delta T = 20 \text{ }^\circ\text{C}$ this would mean 840 kW of thermal power for the
545 experiment). The intensity of the effect of Q settings varies between the different zones. For
546 instance, the lower part of the tomograms in Figure 9 is not affected.

547 **4.4 Application window**

548 The insight gained from variable injection rates and temperatures revealed that the presented
549 tomographic inversion method is robust within a broad range, but has limitations. But what
550 exactly are the limits? We tested a broad range of different scenarios to delineate a general
551 application window, where the inversion method can be used to reconstruct the distribution of



552 K in an aquifer. The parameters listed in Table 2, injection temperature, injection rate and
553 ambient hydraulic gradient, were systematically varied within the given ranges. These ranges
554 were rigorously set, and for reaching possible theoretical limits, some scenarios even
555 exceeded the technically feasible range. Additionally, in the three profiles (Figure 3), the
556 contrasts in the values of K were artificially modified. This was done by expanding or
557 squeezing the original value range for a profile by a factor (range multiplier) between 0.1 and
558 100. As a result, the original structures of the analog were kept while the variance was
559 changed.

560 Each inverted K -distribution was compared with the (scaled) analog profile, qualitatively and
561 quantitatively. A first visual test showed whether major structures were reconstructed and the
562 geometries are similar, especially focusing on the conductive zones (Figure 3). Only
563 acceptable tomograms were kept for the subsequent quantitative analysis.

564 The quantification is based on an estimated connectivity time between the sources and the
565 receivers. The connectivity time is calculated by converting the K -tomogram into a velocity
566 field, using the Darcy equation. With this velocity field, the shortest travel route and time is
567 calculated between all possible source-receiver combinations using the A* pathfinding
568 algorithm (Hart et al., 1968). The difference between the connectivity times on the original
569 model and the inverted result is used to quantify result quality, and by this, define an optimal
570 application window for the method.

571 For condensing the results into a normalized parameter space and plotting them in a 2-D
572 coordinate system, two dimensionless parameters are selected: The thermal Péclet number
573 (Pe_t) to characterize the hydraulic conditions of the subsurface and the effective injection
574 power to describe the used technical parameters of the experiments. Pe_t is calculated
575 separately for the four identified zones of the aquifer:



$$Pe_t = \frac{C_w q d}{\lambda} \quad (13)$$

576 where q is the Darcy velocity, λ is the thermal conductivity and d is the length scale, which is
577 here set to unity thickness of the aquifer ($d = 1$ m). The used technical parameter effective
578 injection power, P' , is defined as:

$$P' = Q' \Delta T = \frac{Q}{qd} \Delta T \quad (14)$$

579

580 where the effective injection rate Q' represents a normalized rate related to prevailing
581 groundwater flow velocity and calculated for the given length scale, d . Note that Pe_t and P'
582 are not completely independent; using a higher injection rate can increase the Pe_t of a zone.
583 Thus, the defined coordinate system is not orthogonal.

584 After evaluating approximately one hundred different experimental scenarios, resulting in
585 over 350 data points, the application window of the method is identified. In figure 10
586 continuous lines mark strict boundaries between feasible and unfeasible regions, and dashed
587 lines denote an approximate boundary where the result quality of tomograms start to decrease
588 in the lateral direction.

589 If Pe_t is below a critical value, the inversion method is not able to provide any hydraulic
590 information for the investigated zone because the assumption that the heat transport is
591 advective is not valid anymore. In this region, the heat transport is governed by thermal
592 diffusion, and no information on K can be extracted from the heat tracer data. A good
593 example of this is the top low conductivity zone on Figure 8b-f, which is not reconstructed
594 properly in any of the presented tomograms. Zones characterized by such low Pe_t are
595 typically short-circuited via adjacent conductive zones. The critical Pe_t number rises non-
596 linearly with the increase of P' . By raising Pe_t with higher injection rate, advection can be



597 promoted in these zones. This provides some information for the tomogram, but the flow
598 field is not short-circuited via an adjacent zone (Figure 7f), yielding a shadow-zone (top low-
599 K zone).

600 At low P' , the amplitude of the tracer breakthrough tends to be too small to be measured in
601 enough observation points to successfully perform the K reconstruction. This strict limit for
602 the application window is due to the assumed 0.1 °C limit for temperature measurement
603 accuracy. It can be overcome by increasing the injection rate or temperature.

604 The result quality gradually declines towards high Pe_t and high P' . This is caused by the
605 distortion of the flow field from high injection rates (see Figure 9f). Reconstructions,
606 therefore, may still be acceptable beyond the given dashed boundary. Note that in practice,
607 this region is infeasible, hence barely relevant. This is because it corresponds with an
608 injection power of 500 kW - 1 MW, and thus this region is also technically infeasible or at
609 least not favorable.

610 **5 Conclusions**

611 Early arrival times of tracer breakthrough curves (BTCs) are specifically suited for
612 identifying highly conductive zones in heterogeneous aquifers. In our study we formulated a
613 procedure for combined inversion of multiple early arrival times measured during cross-well
614 tracer testing. A tomographic set up with multi-level tracer injection and observation was
615 implemented in a model with a 3-D high-resolution aquifer analog, and we examined the
616 capability of the inversion procedure to reconstruct the heterogeneous distribution of K . Heat
617 was selected as a tracer, which offers several advantages in comparison to many solute
618 tracers, but its applicability is traditionally considered limited due to the higher diffusion and
619 coupled thermal-hydraulic processes.



620 It is demonstrated that the tomographic interpretation of heat tracer signals is well suited for
621 characterization of aquifer heterogeneity. By picking early arrival times, the impact of
622 thermal diffusion, buoyancy, and viscosity variation is minimized and in this way, inversion
623 becomes quasi insensitive to the temperature range. The presented application window of
624 thermal tracer tomography is wide, and it covers three orders of magnitude for thermal Péclet
625 numbers and five orders of magnitude for injection power. A key principle is that the
626 transport in the aquifer is dominated by advection, and injection of hot water causes minor
627 distortion. This can be controlled, for instance, by establishing a forced gradient between
628 injection and observation point by operating an adjacent pumping well.

629 The travel-time based inversion is a fast and computationally efficient procedure, which
630 delivers a tomogram in a few minutes with six sources and receivers. It is revealed that not
631 only structures of mainly highly conductive zones could be reconstructed, but also the values
632 of K were closely matched. This is appealing keeping in mind that the presented eikonal
633 inversion is based on a rough approximation of groundwater flow and transport by a wave
634 equation. Yet when close to strong contrast boundaries, the procedure is not able to
635 reconstruct low conductivity zones due to short circuit-shadow effects. To reconstruct these
636 hidden features, a full heat transport model calibration would be required.

637

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643

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785 **Tables**

786 **Table 1:** Hydraulic conductivity, K , porosity, ϕ , thermal conductivity, λ , and bulk heat
 787 capacity, C , for the nine facies that build up the Descalvado analog. The four zones are
 788 introduced for discussion of results and listed here with the major facies components.

zones	facies number (original code)	K [m s^{-1}] ¹	ϕ ¹	λ [$\text{Wm}^{-1}\text{s}^{-1}$] ²	C [$\text{MJ m}^{-3}\text{K}^{-1}$] ²
top low conductive	H1 (St,f)	$6.23 \cdot 10^{-6}$	0.24	3.19	2.49
	H2 (St,m2)	$2.49 \cdot 10^{-5}$	0.29	2.85	2.60
	H3 (St,m1)	$5.97 \cdot 10^{-5}$	0.29	2.85	2.60
central conductive	H4 (Sh/Sp,m1)	$1.38 \cdot 10^{-3}$	0.33	2.61	2.69
lower-central	H5 (SGt,c)	$2.96 \cdot 10^{-4}$	0.32	2.66	2.67
	H6 (SGt,m)	$9.44 \cdot 10^{-5}$	0.32	2.66	2.67
	H7 (Sh/Sp,m2)	$7.77 \cdot 10^{-5}$	0.29	2.61	2.69
bottom	H8 (Sp,f)	$1.63 \cdot 10^{-4}$	0.25	3.12	2.51
	(clay)	$7.84 \cdot 10^{-8}$	0.29	1.90	3.00

789 [¹Höyng *et al.* 2014, ²Bayer *et al.* 2015].

790 **Table 2:** Parameterization of experimental setups, with base values and minimum-maximum
 791 ranges.

parameter	base case	minimum	maximum
injection rate, Q [ls^{-1}] ¹	1	10^{-3}	10
groundwater	10	-	-



temperature, T [°C]			
injection temperature difference, ΔT [°C]	20	5	80
regional hydraulic gradient, i	0.01	10^{-3}	0.1
K range multiplier	1	0.01	100

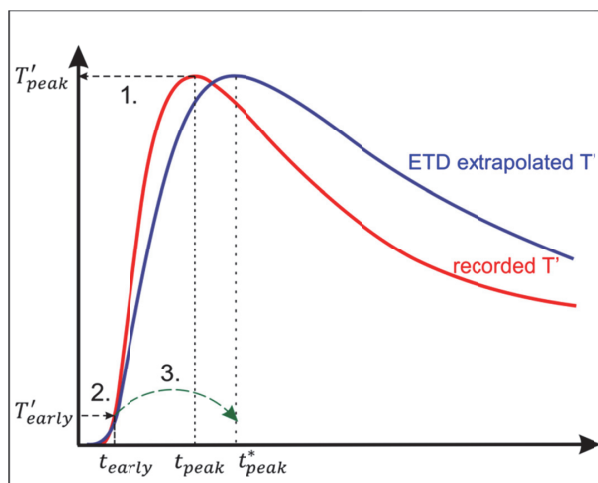
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Figures

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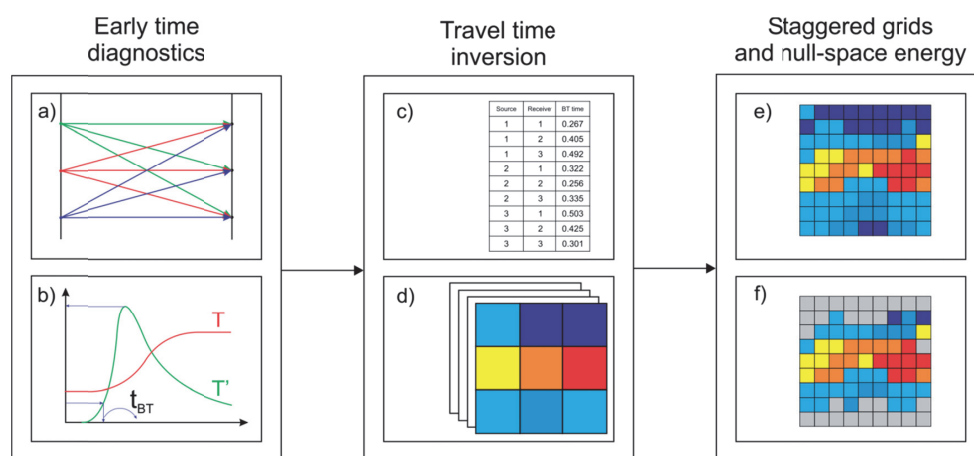
796 **Figure 1.** Three steps of applying early time diagnostics on a BTC. 1. Identify the peak T'
797 value on the recorded BTC. 2. Find the early time value to the corresponding fraction of the
798 signal. 3. Extrapolate the early time to the ideal peak time using the transformation factor, f_α .

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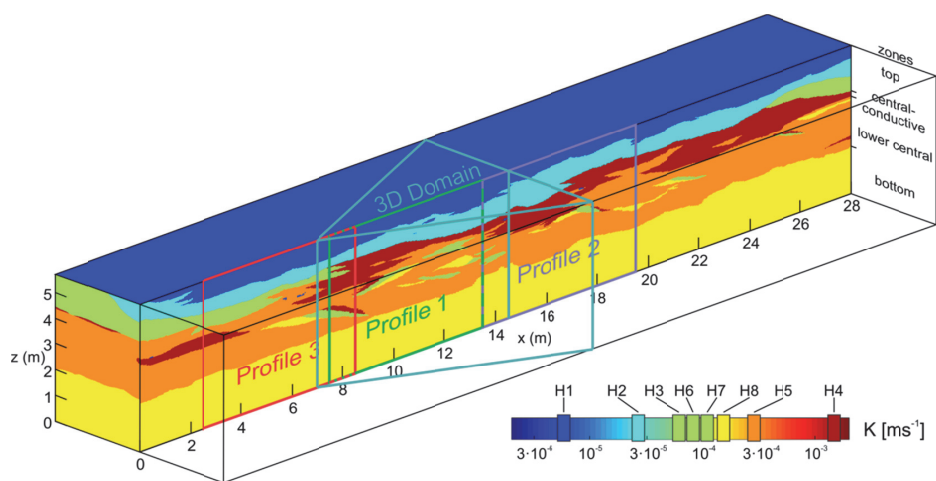
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803 **Figure 2.** Major steps of inversion methodology: a) Conceptual setup of thermal tracer
804 tomography, b) breakthrough time detection using the early arrival times, c) tomographical
805 breakthrough time dataset, d) inverted tomograms applying the eikonal solver on different
806 shifted grids, e) high-resolution tomogram after merging the staggered results together, and f)
807 non-reliable pixels are masked after null-space energy calculation.

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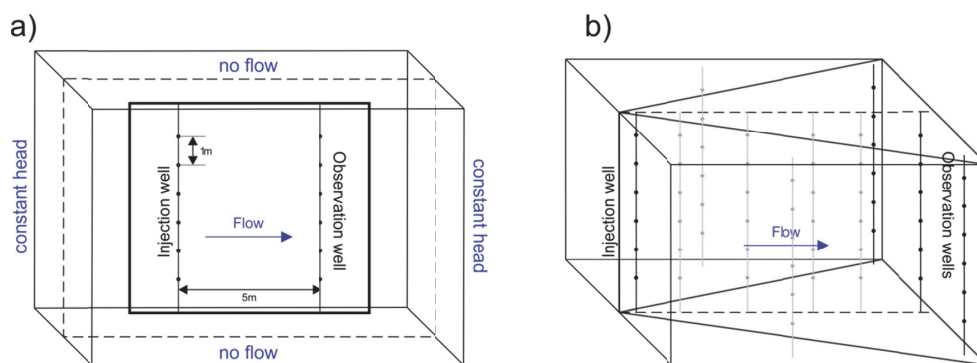


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810 **Figure 3.** Vertical cross section through the center of the 3-D Descalvado analog dataset. H1-
 811 8 represent the hydrofacies units (ignoring clay intraclasts). The location of the three 2-D and
 812 one 3-D profile is marked with different colors.

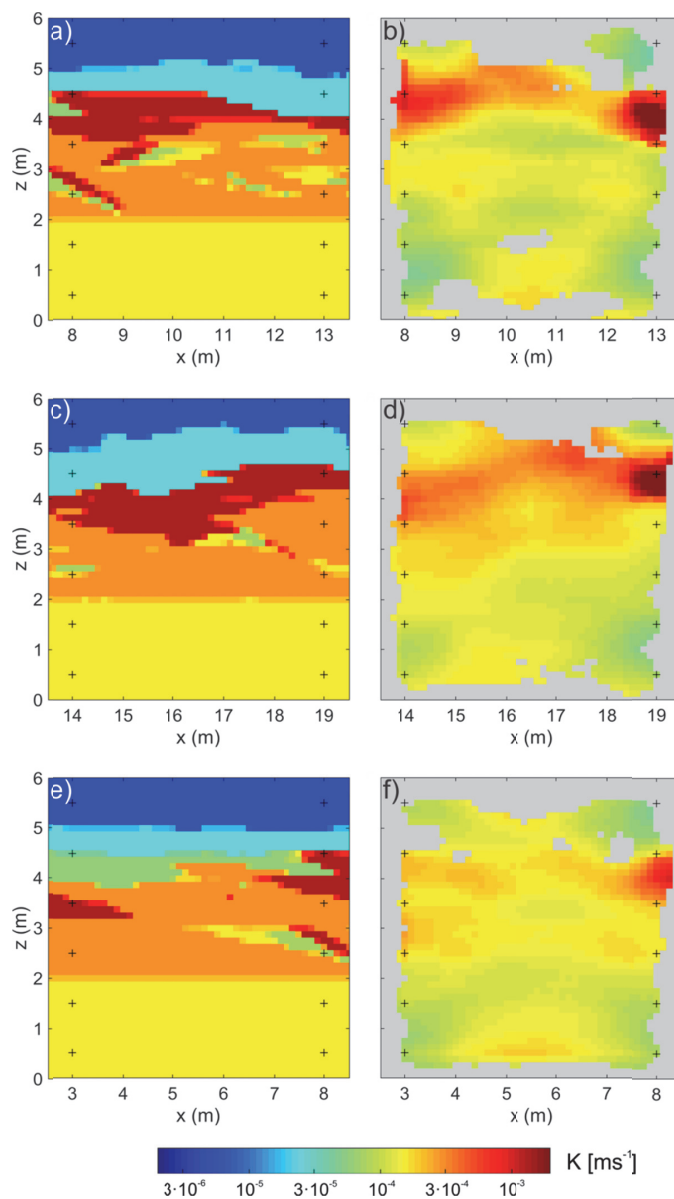
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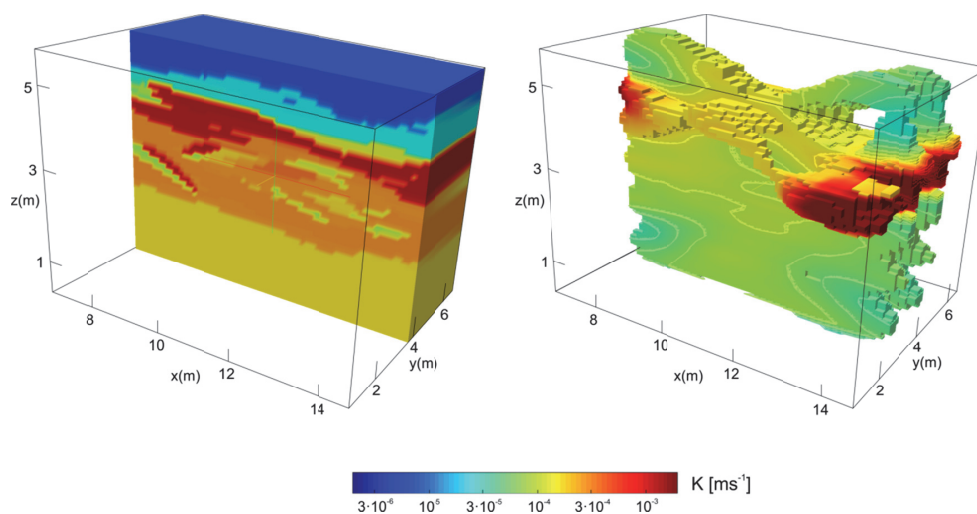
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816 **Figure 4.** a) Simulated experimental configuration and numerical model boundary
 817 conditions. The tomographical setup consists of six sources in the injection well and six
 818 receivers in the observation well. b) Setup of the 3-D experiment with one injection and 3
 819 observation wells. Additional wells used for validation are marked with grey color.



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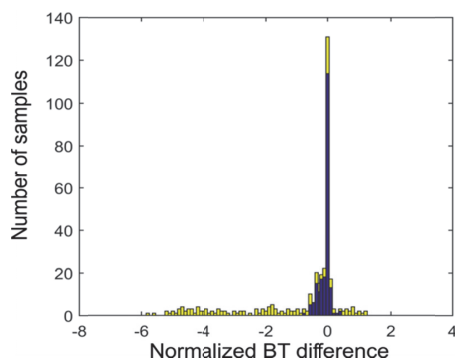
821 **Figure 5.** Reconstructed hydraulic conductivity profiles (see Figure 3): a) profile 1 – original,
822 b) profile 1 – reconstructed, c) profile 2 – original, d) profile 2 - reconstructed, e) profile 3 –
823 original, and f) profile 3 - reconstructed.



824

825 **Figure 6.** 3D reconstruction result: a) investigated subdomain of the upscaled aquifer analog

826 b) reconstructed domain with additional contour lines and unsliced high- K zone.

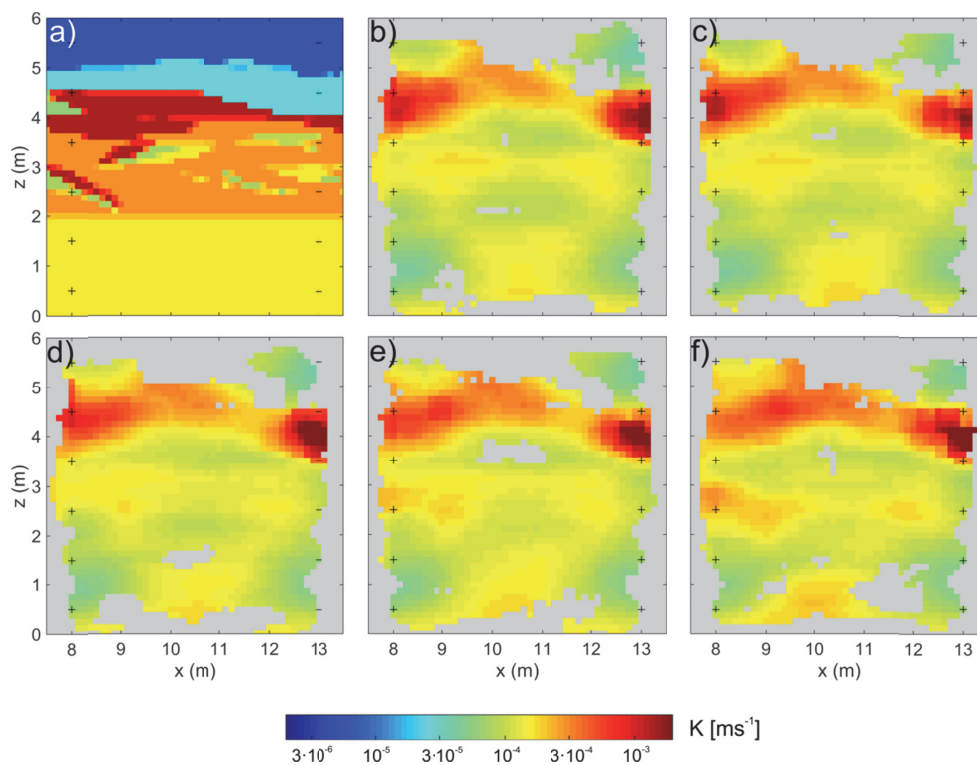


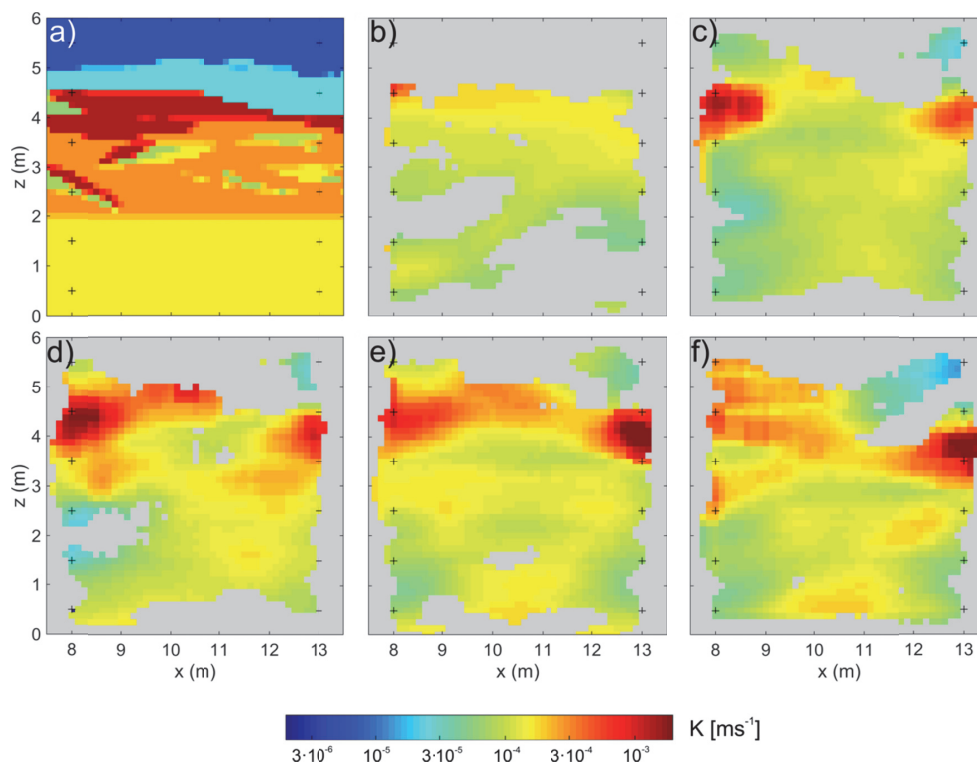
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828 **Figure 7.** Absolute differences of breakthrough times between the inverted and the original

829 model (normalized to the mean of the breakthrough times). Yellow color marks the known

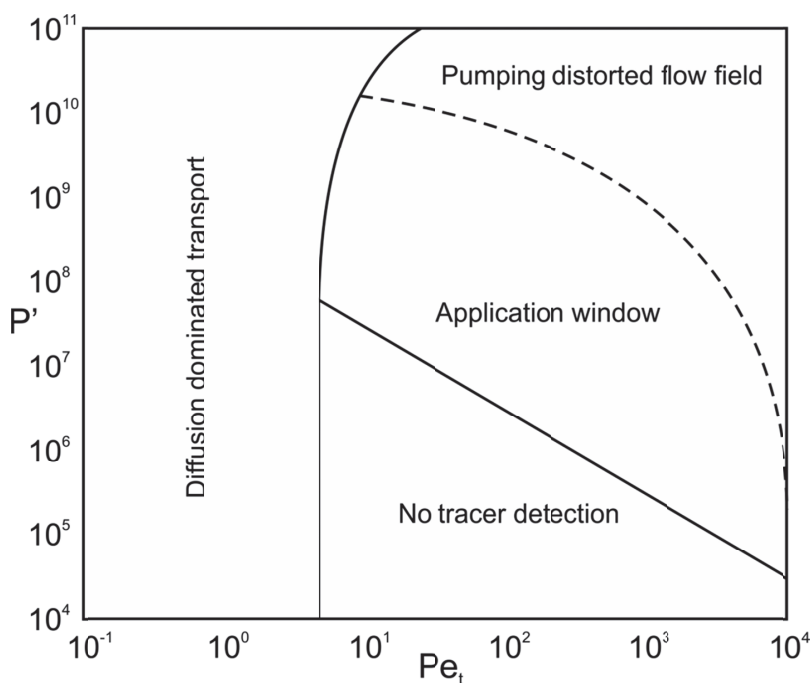
830 outliers, such as observation points in the top low- K zone and the far end of the domain.





834

835 **Figure 9.** Hydraulic conductivity K reconstructions with different injection rates. a) Original
836 K profile 1, b) $Q = 0.001 \text{ ls}^{-1}$, c) $Q = 0.01 \text{ ls}^{-1}$, d) $Q = 0.1 \text{ ls}^{-1}$, e) $Q = 1 \text{ ls}^{-1}$, f) $Q = 10 \text{ ls}^{-1}$.



837

838 **Figure 10.** The proposed application window of the thermal tracer tomography – related to
839 P' and Pe_t . At low injection power, the temperature change at the observation points is below
840 $0.1\text{ }^\circ\text{C}$ and no detection is possible. If Pe_t is below a critical value, the heat transport is
841 diffusion-dominated, and no hydraulic information can be inverted from the tracer travel
842 times. At very high P' the high injection rate distorts the flow field and the results.

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