

Comments to the Associate Editor

Dear Associate Editor, please find attached the revised version of the manuscript and the change-tracking version (kept in the procedure to easily identify all modifications made from the previous version, either recommended by the reviewers or being in straight line).

This is the second round of revision as we have previously answered the questions of Reviewers #1 and #2 who validated several aspects of this work (e.g. the organization of Section 2 into specific subsections for "water flow" and "morphodynamics"). Consequently, the two most recent reviews were attributed to Reviewers #3 and #4 in this reply.

As you have noticed, there was quite a big amount of work still to be done on the paper, and we are grateful to both reviewers for their careful reading of our previous attempts and also for the numerous suggestions they provided - which we generally followed, only to a few exceptions.

So you will find in this document our answers to the reviewers' comments and the associated changes made in the manuscript, indicating the corresponding lines in the revised version of the manuscript. We have also performed a few additional minor changes, correcting typos here and there, as well as a few lines attached to the description of Figure 10 (L961-963) whose header was also slightly changed.

We would like to draw your attention to the fact that some of the reviewers' comments were contradictory with recommendations made in the previous reviews, and this will probably happen in the future rounds of review, if any. Our point here is that we hope this work has become generic enough (although not exhaustive) and of sufficient scientific quality for the exposed analysis to be of interest for the readership of HESS.

Finally, this work aims at establishing the proposed classification procedure on firm grounds rather than providing an extensive database of test cases. However 68 new references have been added during this revision stage. These are listed as "additional references" at the end of this document.

Best regards from Montpellier,
Bruno Cheviron & Roger Moussa.

N. B.

- For simplicity, the previous versions of the summary, equations and figures have been deleted from the change-tracking version of the manuscript, provided as an additional file. Both the revised version and the change-tracking version therefore include the new Figures only.

- Some of the added references have been included in the Figures, therefore (slightly) changing the associated numerical values reported in the text (as well as in Fig.8).

- The list of plotted references has also been updated in the Appendix A (20 more entries)

- There was an error in the RANS equation of the HESS-D document : the correct equation is Eq.3 in the revised manuscript.

REVIEWER #3

General comment

The paper presents a review of modelling approaches to free-surface flow and morphology. It aims to characterise the modelling choices which are gathered from published studies, as function of the dimensional and dimensionless characteristics of the physical system to be modelled. The final goal is to establish a list of guidelines and best practices helping modellers in choosing the right model.

Four different flow models (namely the Navier Stokes Equations, the Reynolds-Averaged Navier Stokes Equations, the Saint Venant equations, and Approximations to the Saint Venant equations) and related morphological models are considered. After introducing the equations, the paper proceeds by reviewing the use of these models in published studies, to build up a classification of model usage as function of various dimensional and dimensionless criteria. In detail, initially the correlation between model choice and spatial and temporal scales of the study domain are analysed. Then, a classification of flow typologies is introduced, based on characteristic depths and slopes, and modelling choices are matched to flow typologies using the above schematisation. Finally, the usage of models as function of characteristic values of dimensionless numbers is analysed.

The paper raises a hot topic, since the increasing availability of computational tools based on different and competing mathematical models requires modellers to be increasingly aware of the range of validity and best usage of models. In fact, quoting Escauriaza et al. [2015] in his very recent review of morphodynamic models for gravel-bed rivers,

At present there is no systematic, reliable method to define the model category appropriate for a specific phenomenon in nature.

Thus, this paper aims to provide a number of criteria to address this issue.

Having said this, I have a number of major concerns. I think that the proposed classification criteria do not fully address the matter, and therefore the manuscript does not fully reach the objective (as stated in the abstract) *to help each modeller positioning his (her) choices with respect to the most frequent practices, within a generic, normative procedure*. In this respect, I think some different view could be incorporated, or at least discussed. Furthermore, I find the presentation of the morphodynamic part of models a bit generic, unnecessarily stretched over different subsections, and incomplete of some recent development. These concerns are discussed in detail below.

In addition to these, I also have a number of minor concerns and suggestions. Among these, I would suggest simplifying the language (especially in the introduction) for improving understanding. Minor issues are listed in detail below.

This is why I recommend a severe major revision.

Major issues

1/ I think that the present discussion on determinants of modelling choices misses a fundamental point: how modelling choices are determined by the objective of the study, i.e. by the natural phenomenon to be modelled. By this, I mean that the same flow and erosion event/process can be described by different models, depending on the required detail in the modelling description. One of the possible examples of this issue is the dynamics of a dune-dominated fluvial bed, which, depending on the modeller's focus, can be readily studied by the SV approximation with appropriate continuity models for sediment [e.g., Ribberink, 1987, Blom, 2008] or instead by much more refined flow and sediment transport models, which fully take into account the vertical coordinate and the distribution of forces on each sediment particle [e.g., Nabi et al., 2012, 2013b,a]. In my view, to select the most appropriate model for their applications, modellers shall compare the smallest spatial/temporal scale at which the selected model operates, and the scale of the phenomenon to be analysed. Therefore it is not an intrinsic scale of the physical system/domain, but instead the scale of the object of interest within that system, which dictates the modelling choice. The multiplicity of modelling approaches allowed for the same case may help to explain the overlap of approaches across scales which is observed in this manuscript.

We completely agree with this remark that "the end justifies the means" in the choice of modelling strategies. Following your suggestion (and these made by Reviewer #4) we acknowledge the idea that the objectives of a study show through the choice of appropriate subscales (say δL and δT) either directly cited in numerical procedures or as the spatial and temporal resolutions of data collection. Accordingly, we are now trying to correlate the choice of a model to the size of the domain (L , T , H) and to the size of its subscales (δL , δT) before examining the influence of flow typologies and dimensionless numbers.

The text of the manuscript has been changed in the abstract (L24, 26 & 32), the introduction (L84, 88, 90 & 151), the conclusion (L1013-1014 & 1033) and in several occurrences of the term "subscale" or in direct association (e.g. L540 & 544-545).

Figure 2 now has two subplots, Fig. 2a (unchanged) and Fig. 2b that positions modelling and/or data collection scales in the $(\delta L, \delta T)$ plane for complementary indications, to gain a more comprehensive view. The added Fig.2b and its legend appear hereunder.

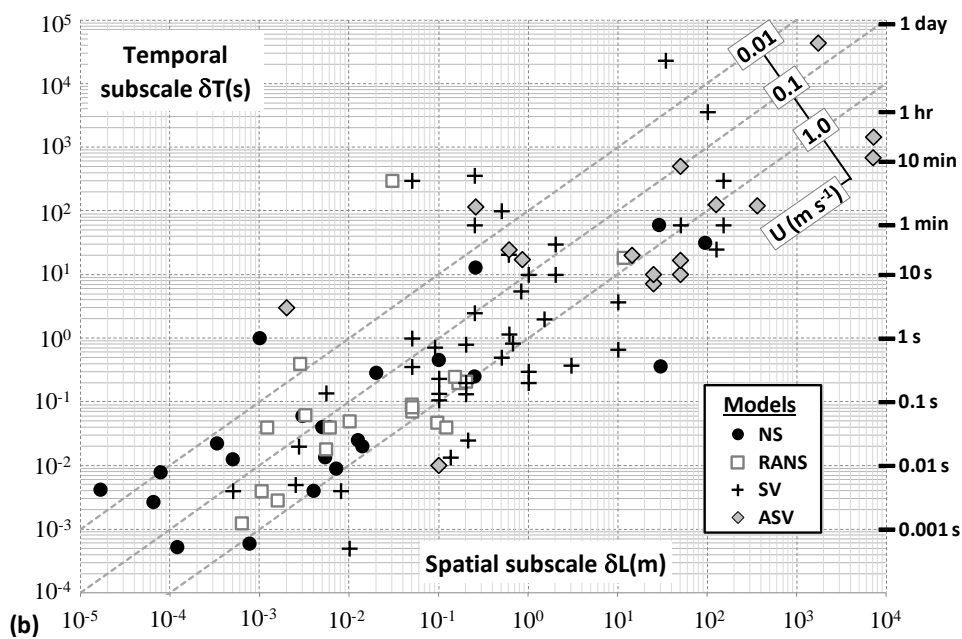


Figure 2 – How increasing (L, T) spatiotemporal scales (a) and $(\delta L, \delta T)$ subscales (b) of the flow domain tend to be associated with decreasing complexity in the choice of flow models, sorted here into four levels of refinement: Navier-Stokes (NS), Reynolds-Averaged Navier-Stokes (RANS), Saint-Venant (SV) or Approximations to Saint-Venant (ASV). A transverse analysis involves forming L/T ratios, searching for clues to model selection according to these "system evolution velocities" or governed by flow typologies that would exhibit specific L/T ratios (a). Unit values of the Courant number ($Cr = U\delta T/\delta L$) have been used to trace characteristic flow velocities of $U = 0.01, 0.1$ and 1 m s^{-1} and the indicative numerical stability criterion is $Cr \leq 1$: for given δL and U values, δT should lie behind the dotted line (b). Both plots were assembled from information available in the studies cited in Appendix A, selecting six textbook cases (sketches A to F, Table 1) for illustration (a).

2/ Another issue originates from the same observation. When comparing modelling choices for the same natural phenomenon, a more detailed description of the system, by the use of a more complex model, does not necessarily represent a refinement, but instead a fully different view on the same physical system. In other words, using a more complicated model does not necessarily allow to replicate what a simpler model would do, plus adding more information, but instead can produce a completely different outcome. An example is given by Sloff and Mosselman [2012],

who compared the results of two different continuity models (with and without mixed sediment) for the same river bifurcation. I would therefore like this issue to be discussed with reference to the modelling approaches analysed in this manuscript.

The recommended elements of discussion may now be found in the introduction (L113-117). *"Alternatives consist in examining the "scale matching" between available data and modelling aims (Lilburne 2002, Kim & Ivanov 2015) and the possibility to use a more complicated model not only because it replicates what a simpler model would do, plus additional information, but also because it offers different, specific outcomes (e.g. Sloff & Mosselman 2012)."*

3/ Furthermore, the manuscript misses the analysis of a critical point in modelling, which is cost-effectiveness and feasibility. By cost-effectiveness, I mean the possibility of reproducing the observed behaviours by including the minimum amount of processes [e.g., Escauriaza et al., 2015] and trying to minimize computational cost, which, although less of a limiting factor than in the past [e.g., Mosselman, 2012], still is a critical point in making modelling practically feasible. An effort towards simplicity could indeed help reducing parametrisation, data requirement, and thus minimizing well-known model shortcomings such as equifinality and other modelling mistakes [see, e.g., Mosselman and Le, 2016]. The principle of minimizing modelling effort could then help modellers in better placing their modelling efforts within the classification diagrams proposed in this manuscript, when working in regions characterised by significant overlap of different modelling choices. This concept is somehow buried inside the discussion (lines 960-964) but would in my view require some expansion and references to available studies on the matter.

We agree with this remark of analytical nature. We have thus added a few words on the subject in the conclusion (L1126). However, these aspects were already present in the same paragraph, especially when explicitly referring to the principle of parsimony (L1087-1089) and also L1123-1130 from a wider point of view on the philosophy of modelling (Section 4.2).

Finally, this paper is turned towards the examination of literature elements, its aim is not explicitly to address the question of good practices (see Figure 11). This is why we did not go very far in this direction yet.

4/ I have some concerns about the structure and content of Section 2, regarding erosion models. I am fine with the presentation of flow models from the most complete (NS) to the most simplified (ASV). However, the presence of individual subsections devoted to erosion model associated to each flow model looks unnecessary, because the difference between the erosion models in each subsection is unclear. Furthermore, the relevant equations are not presented, which does not help in understanding. Finally, some recent developments such as the direct numerical simulation of turbulent flows based on the NS equations and particle-based morphodynamic models [Kidanemariam and Uhlmann, 2014, Colombini, 2014] could be incorporated.

-Section 2 is the Section that we changed the most, following the recommendations. We agree that Section 2 was a bit clunky. We have now introduced the equations for the morphodynamic side, so that each level of refinement (NS, RANS, SV, ASV) is treated from the points of view of "1. Water flow, 2. Morphodynamics" modelling (thus also changing the titles of the dedicated subsections from "Erosion" to "Morphodynamics").

-A consequence is that the subsections (2.2.2, 2.3.2, 2.4.2 and 2.5.2) devoted to morphodynamics have been enriched and (partly) rewritten, with more references and hopefully stronger physical/analytical arguments for more balance between the "water flow" and "morphodynamics parts".

-A corollary is that the writing of most flow equations has been modified so as to ensure coherent formulations and easier comparisons between the flow equations and the newly introduced morphodynamic equations.

-On the other hand, we have kept this "1. Water Flow /2. Morphodynamics" structure in all subsections because the morphodynamic approach always adds complexity to the water flow issue, thus deserving specific, additional indications (this point was validated by Reviewers #1, 2 and 4).

- The mentioned two references and many others have been added to substantially enrich the NS subsection (2.2.2).

Minor issues

- In the title and throughout the paper: I consider "erosion issues" as an excessively limiting and possibly misleading term, as the modelling here considered both includes erosional and depositional processes. I would therefore change it into "morphodynamics".

We followed this suggestion and have changed most occurrences of the "erosion" term into morphodynamics, wherever possible and especially where "erosion" was previously used in the wide sense of erosion, deposition and transport.

We have also followed one of the recommendations of Reviewer #4 and explicitly included "hydraulics" in the title, now: **"Determinants of modelling choices for 1-D free-surface flow and morphodynamics in hydrology and hydraulics: a review"**

- Line 34: "help each modeller positioning his (her) choices" ! "help modellers in positioning their choices".

Done

- Line 80: "... through successive flow aggregations over various bed topographies" not clear.

Replaced by "involving several levels of flow aggregations in the streamwise direction" to introduce the idea of streamwise scenarios mentioned next.

- Line 83: "main rivers".

We think that "main rivers" is the appropriate term because we were mentioning progressively wider scales, starting from runoff initiation to the largest documented river flow.

- Lines 89-90: "which is the open normative procedure designed to allow comparisons between studies and to be fed to the community" not clear.

The last part of the phrase has been replaced by a new phrase: "This signature is thought normative enough to facilitate comparisons between studies."

- Line 92: "genericity" ! "generality".

Done

- Line 104: "Earth's surface" ! "earth surface".

Done

- Lines 109-111: "under the angle of connecting ... or debating the merits ...": not clear.

The phrase is now L111-113 and reads "...under the angle of hydrological and sedimentological pathways (see the review by Bracken et al. 2013) or questioning the merits of similitude laws and

these of upscaling methods in the description of hydrological processes (Strahler 1956, Blöschl and Sivapalan 1995, Slaymaker 2006)."

- Line 113: It is not clear what "This" refers to. In the same line, I have some trouble with the definition of "contextual" and "strategic". I would also ask for rewriting of the rest of the sentence, which is not clear to me.

The expression "scale matching" has been placed between quotation marks in the preceding phrase. The confusing sentence has been rewritten (L117-120).

" With similar goals but a different framework, this study proposes an overview on the most popular modelling practices, confronting the theoretical refinement of flow models to the spatiotemporal scales and characteristics of the described free-surface flows."

- Line 137: It is misleading here to mention "reduced complexity models" as a synonym of models in river morphodynamics based on the approximation of the sediment discharge being equal to the transport capacity calculated locally or including a transport distance. In fact, these assumptions are very widely applied in 1D and 2D models usually regarded as "physics based" (e.g., BASEMENT [Vetsch et al., 2014], and Delft3D [Sloff et al., 2001]). Conversely, the models usually regarded as "reduced complexity" [e.g., Murray and Paola, 1994] may apply even cruder estimates of sediment discharges.

The phrasing has been modified (L142-145).

"...flow energetics (Vanoni 1946, Hino 1963, Lyn et al. 1992, Mendoza & Zhou 1997) while most 1D or 2D physics-based models (e.g. Sloff et al. 2001, Vetsch et al. 2014) either assume the "transport capacity" (Foster & Meyer 1972, Bennett 1974) or "transport distance" schools of thoughts (see details in Wainwright et al. 2008)."

- Line 141: "erosion science".

Replaced by "morphodynamics"

- Line 145: "conceptual element", "contextual element". The meaning of these items in my view is neither obvious nor standard in the literature. A definition would be needed.

The formulation has changed to be more explicit (L150-154)

" The methodology consists in defining the "signature" of each case study as the chosen model refinement and modelling subscales vs. the given spatiotemporal scales, flow typology, and dimensionless numbers, hypothesizing the conceptual element (model refinement and spatiotemporal subscales) is the consequence of the contextual elements (flow scales, typology and dimensionless numbers)"

This is simply to outline/introduce the existence of a causal link between what is chosen (modelling strategy, the concepts) and what is given (the contextual elements).

- Sections 2.1, 2.1.1, 2.1.2: The content of these very short subsection is essentially an introduction to the more detailed description of models in the following subsections. I therefore advise to simply remove these headers and make it just an introduction to the following content.

Done. These headers have been merged into an introduction to Section 2, with modifications to the paragraph that previously described erosion (see the full paragraph L163-190).

- Section 2.2.1: Authors may insert a sentence to clarify that the list of models here considered does not enclose all the possible modelling choices in the literature.

Done. We have added that the choice made here was "among many other possibilities"

- Line 157: "from the richness of their physical basis" ! something like "depending on the degree of refinement in their physical description".

Done. We replaced the wrong formulation by ", i.e. depending on the number and nature of the indications included in their physical description".

- Line 160: please introduce the acronym "ASV".

Done.

- Line 161: "Diffusion Wave" ! "Diffusion Wave Equation", "Kinematic Wave" ! "Kinematic Wave Equation".

Done.

- Line 169: "... the examination of erosion issues from the angle of decreasing refinement ... as a whole" not clear.

Suppressed because the paragraph has been rewritten.

- Line 171: "disconnection" ! "discrepancy" or "inconsistency".

Done.

- Equation (1) is not the full set of the 2D x-z NS equations. The vertical momentum equation and the continuity equation are missing. Authors are in general advised to present mathematical models more thoroughly.

This paper focuses on 1D approaches thus on the projection of the momentum equation on the x axis. The continuity equation is not shown because (i) a full presentation of the mathematical models is not the purpose here and (ii) the momentum equation (on x) seemed enough to illustrate the main differences between the NS, RANS, SV and ASV refinements.

- Lines 194-200: I do not understand why the discussion over turbulence models, until the development of the RANS equations, is placed here, since the RANS are then introduced in another section.

We agree. These six lines have been placed in the RANS subsection (2.3.1, L269-275). The following three lines have been placed at the end of the NS paragraph (section 2.2.1, L211-214).

- Section 2.2.2: Here "erosion issues" are presented without mentioning any continuity model for sediment. A presentation of the the continuity framework could be useful, as the resulting equations are themselves part of the hydro-morphodynamic mathematical model.

As previously mentioned, we added equations for the morphodynamics with associated comments and literature elements throughout section 2.

- Line 219: "debating the case of turbulence damping ..." please rephrase for clarity.

Replaced by "examining turbulence damping"

- Line 221: "the matter is not free from doubt today" please rephrase for clarity.

In coherence with the introduction of a NS-level equation for morphodynamics and the associated comments (L216-237) we have rephrased the sentence and added a few lines for a wider overview on

the topic (following one of your more general suggestions in the major issues you have raised) (L255-265).

"The matter is not completely free from doubt today (Kneller & Buckee 2001) though the diagram proposed by Elghobashi (1991, 1994, p310) to describe the regimes of interactions between particles and turbulence seems rather widely accepted. For the most dilute suspensions ($c_d < 10^{-6}$) the sediment load is not supposed to have any influence on turbulence characteristics. For the intermediate case ($10^{-6} < c_d < 10^{-3}$) the sediment load is supposed to enhance turbulence only if the particle response time is at least two orders of magnitude greater than the Kolmogorov time scale, i.e. the characteristic time for the turbulent eddies to vanish: for the same sediment load and water viscosity, larger particles tend to enhance turbulence while smaller particles tend to damp it. For dense suspensions ($c_d > 10^{-6}$) frictional drag, abrasion due to impacts of the travelling particles and increased flow viscosity have been described prone to enhance the detachment capacities of loaded flows (e.g. Alavian et al. 1992, Garcia & Parker 1993)."

- equation (2): same comment as for equation (1).

See the answer to this comment.

- Line 248: "judiciously".

We deleted this term.

- Line 255: "private hunting grounds".

We used "an interesting feature" instead.

- Lines 272-274: The hypothesis of shallow water ($H \ll L$), which limits the admissible free-surface slope and implies quasi-hydrostatic pressure distribution over the vertical, shall be mentioned when introducing the Saint Venant equations.

We agree. These elements now appear in the text (L337-340).

"The SV equations also termed "shallow water equations" assume the $H \ll L$ hypothesis of shallow water which limits the admissible free-surface slope and implies a quasi-hydrostatic pressure distribution over the vertical. The integration process from NS to SV (Chow 1959, Abbott 1979)..."

- equation (3): same comment as for (1) and (2).

See the answer to these comments.

- equations (5) and (6) represent a generalisation of (3). My advice is just to present (5) and (6) instead of (3) for brevity.

We tend to disagree with this comment. In our opinion, the SV level of refinement is a tilting point in complexity between (roughly) the world of environmental fluid mechanics (NS, RANS) and that of hydrology (ASV). So we think we should keep (3), (4) and (5)+(6) to illustrate the gradual shift from one world to the other. This deliberately places emphasis on the central position of the SV-type approaches and the variety of applications they allow.

- Line 309: "erosion-hydrology": definitely "morphodynamics" or "hydro-morphodynamics".

Yes: "hydro-morphodynamics" seems more appropriate.

- Line 310: the Exner equation is presented here, but never properly introduced in the text.

An extended formulations of the Exner equations is now present in the manuscript and discussed L379-387.

- Line 312: there has been quite more recent work on the diffusive character of the Exner equation. See for instance Furbish et al. [2012] and related papers.

Same answer as above.

- Line 312: I do not agree with the Authors on the point that "most studies ... take particle velocity equal to water velocity", as I could not think of a single study in which this assumption is done.

We agree that "most studies" was probably a too strong formulation, but on the other hand the type of assumptions made on "particle velocity relative to fluid velocity" seem to depend on the purposes of the studies (thus, plausibly, on the scientific background of their authors). The formulation has changed (L389-390).

"Conversely, in the field of hydrology, numerous citing papers discard one or several terms from the Bennett (1974) equations, typically taking particle velocity equal to water velocity."

For example, the catchment-scale erosion models not interested in detailed morphodynamics often assume that particles are transported into suspension, at the average flow velocity. Most of the times, such simplified erosion models will just predict erosion, transport and deposition of sediment masses (with no geometrical argument). Precisions have been added in the text accordingly, in subsection 2.5.2, L487-496, when introducing the ASV-level equation for morphodynamics.

Generally speaking, the Exner equation requires the evaluation of solid discharges to be fed to the continuity framework, not of particle velocities. Although some empirical formulae for particle velocity based on the hydrodynamic variables exist, these are not of immediate use for evaluating bedload fluxes. In fact, even if the particle velocity was known, bedload discharges would come from the product of particle velocities and the thickness of the transport layer, once again not precisely known. Therefore assumptions on particle velocities are not generally made. Theoretical studies though exist [Furbish et al., 2012, Ancey and Heyman, 2014, Ballio et al., 2014] in which the complex linkage between particle velocity distribution and sediment discharge in the continuity framework is addressed.

Changes have been made in accordance in the text (L381-383).

- Line 361: "shear stresses are generally calculated from near bed laminar or near laminar profiles". Could Authors provide reference to this statement? I disagree with it, because the sediment transport formulae mentioned in line 365 have all been derived for turbulent natural flows.

We agree for the general case and have changed the phrase (L440-443).

"Although derived for turbulent natural flows, shear stresses may also be calculated from near-bed laminar or near-laminar velocity profiles, sometimes with the regularising hypothesis that detachment and transport occur just above the criterion for incipient motion (see the review by Lajeunesse et al 2010)."

- Line 479: "This dispersion contains a lot of information". I would delete this sentence as unnecessary.

Done.

- Section 3.1.1: In Figure 2 I cannot easily gather a proper trend from the figure in terms of model choice as function of the L/T system speed. Instead, a trend is detectable if we consider the L scale

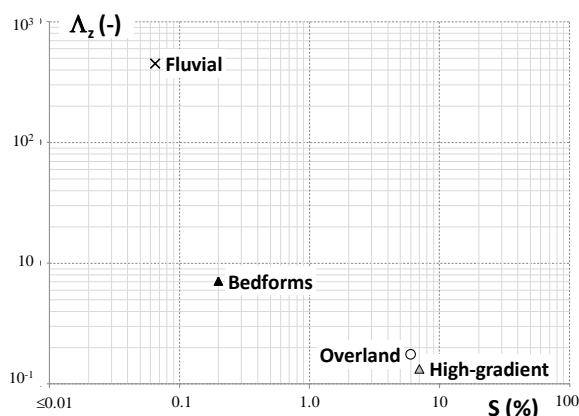
only (e.g., the Navier-Stokes equations seem to have been applied to smaller domains than the ASV equations, which makes sense). Could the Authors discuss on this point within the section?

Yes, we added a few elements on this point (L587-592).

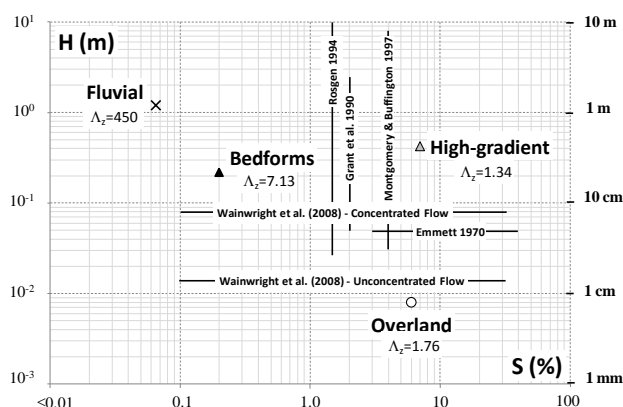
"Most applications find themselves in the $10^{-2} < L/T < 10^2 \text{ m s}^{-1}$ range, exhibiting no clear difference between the NS, RANS, SV or ASV refinements. Conversely, this indicates that each level of refinement has been used to model high or low system evolution velocities, sometimes by relying on specific (adapted or upgraded) formulations of the systems of equation (see for example the hybrid NS-SV level of refinement needed for detailed morphodynamics, especially to reproduce the long-term evolution of bed topography)."

• Figure 5: I have some concerns over the use of a dimensional vertical coordinate such as H for discriminating between flow typologies. In my view, depth does not provide a unique criterion unless it is compared with the size of roughness. This is somehow obviated within the figure by plotting a non-dimensional threshold ($\lambda_z = 100$), but it is then inconsistent with the axis.

When submitting the paper we had these two options. The plot of Λ_z vs. S is shown below.



The S- Λ_z plane shows less differences between median cases than the S-H plane, so the latter was considered to better outline the differences between flow typologies. However, we agree with the reviewer that the indication of Λ_z values was not properly given in the Figure used in the manuscript. We have modified this figure to explicitly mention the Λ_z value that goes with each (S, H) pair of values.



The legend has been modified accordingly, as well as the comments in the text.

The legend to Fig.5 now reads: "Median position of the studies belonging to the "Overland", "High-gradient", "Bedforms" and "Fluvial" flow typologies, plotted on the (S: slope, H: water

depth) plane, with indication of the associated inundation ratio ($\Lambda_z=H/\varepsilon$) This figure was assembled from information available in the studies cited in Appendix A."

The text now reads (L790-791)

"The simplest way to proceed is to work in the (S, H) plane, then to indicate the values of Λ_z for each pair of (S, H) values".

- Section 3.3.1: Please remark that dimensionless numbers arise in the non dimensionalisation of systems of governing equations, i.e., they are an inherent feature of the model in use to describe the physical system.

This is now mentioned in the text. It was a good occasion to add a few missing elements on how to use/see the dimensionless numbers in this study (L893-900).

"On the one hand the dimensionless numbers arise in the non-dimensionalisation of the systems of governing equations, being an inherent feature of the model. On the other hand only the selected dimensionless numbers appear in the non-dimensional formulation of the equations, from appropriate arrangements of their terms, and this choice indicates which are the physical processes of interest for the modeller. Finally, not all dimensionless numbers can be made explicit in the simplest mathematical models (especially the ASV models) but their values can always be calculated, thus correlated (or not) with the use of one or the other of the flow models."

- Lines 806-810: "accelerated by pressure effects" does not necessarily correlate with supercritical flows. Furthermore, the discussion over the direction of propagation of waves is sound only if the bed is fixed. With movable bed, more characteristics come into play, and the identification of a sub- and super-critical regime is not entirely possible [Lyn, 1987, Lyn and Altinakar, 2002].

We agree and the correction is: *"for example flows accelerated by pressure effects"* (L924).

The difficulties that arise for movable beds have also been mentioned (L927-929): *"However, the presence of a movable bed makes the identification of sub- and supercritical regimes less obvious, as additional phenomena come into play (Lyn 1987, Lyn & Altinakar 2002)."*

- Lines 811-813: I have a concern over the use of the bed slope S. Actually, the free-surface slope appears to be much a stronger control over the flow characteristics.

This is also our impression but the bed slope is easily observable (and the free surface slope may be calculated, and so is the friction slope, from S)

- Line 843: "angle" ! "point of view".

Done.

- Line 930: "each modeller ... his or her" ! "modellers ... their".

Done.

- Line 933: "comprehensive view" ! "comprehensive set" or "database".

Done ("database").

REVIEWER #4

MAJOR COMMENTS

For the manuscript, “Determinants of modelling choices for 1-D free-surface and erosion issues in hydrology: a review”, this study attempts to present a normative classification through a comprehensive literature review. I applaud their efforts in incorporating a wide range of studies with various scales, typology, and dimensionless numbers. I feel in overall the manuscript has been well written, so I contend this deserves to be published in HESS if a couple of reservations will be addressed.

- **Literature:** I appreciate with such a comprehensive review, but my first reservation is whether this could sufficiently reflect the current developments of knowledge in our community. The reason I thought is the total number of cited references except for those used in the appendix is about 370, but only 16 references (less than 5%) are recently published within 5 years (since 2011). Readers can wonder all recently reported state-of-the-art studies are well addressed and reflected in this review. I would appreciate if authors can update some parts with new published contents if any.

We agree with this reservation and we have tried to address it in the new version of the manuscript that has now 21 new references in the Appendix (179 cases) and 68 new references in the manuscript, for a total of 409 references, among which 60 (14%) have been published since 2010. All figures have been modified accordingly, except Fig.1 and 4 that are of a conceptual nature. A new Fig.2b was also plotted from the references in the Appendix (to show the δT vs. δL modelling subscales, as described hereunder).

Some examples, although not limited to, are

(1) friction coefficients in overland flow (not river flow) can be explained by using many dimensionless variables but show mixed trends in controlled conditions (Kim et al., 2012, WRR, “*Hydraulic resistance to overland flow on surfaces with partially submerged vegetation*”). Since overland flow move on surfaces with partially submerged roughness elements and very shallow depths of flow, we could not directly employ empirical relationships developed for river flow. As authors also mentioned, the inundation ratio (Lawrence, 1997) is often a key indicator to differentiate overland and river flows, but most of studies (in Fig. 10) did not focus on very small order of magnitude ($\ll 1$) on the ratio (e.g., overland flow on vegetated area). In the above literature, a couple of experimental data was represented for cases with the small ratio numbers. It would be great if authors can discuss and incorporate (in L132, L612, L645, L816, or somewhere) how the friction is addressed in overland flow with larger elements.

There was indeed the need to mention studies more turned towards flows through emergent vegetation. We added a few elements in the section that describes the dimensionless numbers, in the paragraph on the inundation ratio (L936-940).

" *The encountered values of Λ_z are very high for rivers flowing on smooth, cohesive, fine-grained beds ($\Lambda_z > 100$) and very low for all types of flows between emergent obstacles ($\Lambda_z < 1$, Ferro 2003, Hogarth et al. 2005, Canovaro & Solari 2007, Ferguson 2007, Lamb et al. 2008) including flow through vegetation (see Järvelä 2004, Holden et al. 2008, Gumiere et al. 2011a, Kim et al. 2012, Nepf 2012).*"

(2) Please add/introduce erosion controls (e.g., in L335-337) of, for example, scale effect (Kim et al., 2016. *Environmental stochasticity controls soil erosion variability* Sci. Rep. 6, 22065), shielding effect, nonuniqueness (Kim and Ivanov, 2014. *On the nonuniqueness of sediment yield at the catchment scale: The effects of soil antecedent conditions and surface shield*. WRR, 50, 1025-1045; Nearing et al., *Sediment yields from unit-source semiarid watersheds at*

Walnut Gulch. WRR, 43, W06426, 2007), and micro-scale variability (Risse et al., *Assessment of error in the universal soil loss equation*. Soil Sci. Soc. Am. J. 57, 825–833, 1993; Kinnell, *Why the universal soil loss equation and the revised version of it do not predict event erosion well*. Hydrol. Processes, 19, 851–854, 2005), and in L822-834, present an additional importance of shields number that explains the time scale of non-uniqueness (WRR, 50, 1025-1045).

Thanks for the suggestions; we added the references that seemed in direct connection with the purpose of the 2.4.2 subsection, a part of which has been rewritten (L410-416).

"On the other hand, many erosion controls have received attention within the SV or ASV formalisms, i.e. without explicit descriptions of particle-scale flow features: micro-scale variability (Risse et al. 1993, Kinnell et al. 2005), local sheltering effects (Nearing et al. 2007, Kim & Ivanov 2014), slope effects (Polyakov & Nearing 2003), particle-size effects (Van Rijn 1984a, Hairsine & Rose 1992a, Sander et al. 2007, Wainwright et al. 2008), flow stratification effects (van Maren 2007), the effects of hyperconcentrated flows (Hessel 2006). Bedload transport (e.g. Van Rijn 1984b, Julien & Simmons 1985, Hairsine & Rose 1992b, Wainwright et al. 2008) has also motivated the search for dedicated formalisms."

Erosion equations: The second suggestion is, as mentioned by other reviewers, about erosion equations. Unlike flow equation, erosion part is still difficult to grasp because the description of manuscript is more or less written in a narrative way or a list type with many citations only (do not have details).

Both reviews suggest introducing the equations that describe erosion or actually morphodynamics (i.e. erosion, transport and deposition) as required in the other review. We have thus selected the most characteristic formulations of the momentum or mass conservation equations for the solid phase at each description level (NS, RANS, SV and ASV). In the comments next to the equations we have also tried to present and discuss some alternative formulations (Section 2).

We copy here the answer to Reviewer #3 on these aspects.

-Section 2 is the Section that we changed the most, following the recommendations. We agree that Section 2 was a bit clunky. We have now introduced the equations for the morphodynamic side, so that each level of refinement (NS, RANS, SV, ASV) is treated from the points of view of "1. Water flow, 2. Morphodynamics" modelling (thus also changing the titles of the dedicated subsections from "Erosion" to "Morphodynamics").

-A consequence is that the subsections (2.2.2, 2.3.2, 2.4.2 and 2.5.2) devoted to morphodynamics have been enriched and (partly) rewritten, with more references and hopefully stronger physical/analytical arguments for more balance between the "water flow" and "morphodynamics parts".

-A corollary is that the writing of most flow equations has been modified so as to ensure coherent formulations and easier comparisons between the flow equations and the newly introduced morphodynamic equations.

-On the other hand, we have kept this "1. Water Flow /2. Morphodynamics" structure in all subsections because the morphodynamic approach always adds complexity to the water flow issue, thus deserving specific, additional indications.

- The mentioned references and many others have been added to substantially enrich the NS subsection (2.2.2).

I would suggest to explain flow/erosion phenomena with **physical and empirical components**. For example, in rivers, advection and diffusion are of primary physics while suspended and bed

load are main sources which are not directly resolved by physics but addressed by empirical, experimental relationships. In hillslopes, advection is the only driver while splash, rainfall/flow detachments and rainfall/flow entrainments are sources which are only computed by empirical equations. Similarly, what are the main physics and sources in NS-coupled-erosion models? What can or cannot be resolved by physical conservation laws or what can be lumped in each typology case (step-pools, pool-riffles, etc.)?

We agree with this remark and have added some general context elements at the beginning of Section 2 (L177-180).

"Depending on the refinement of the coupled flow and morphodynamics models as well as on flow typology, a clear trend is that some elements are explicitly addressed whenever possible, e.g. particle advection and diffusion, while others are most often parameterised, e.g. particle detachment from excess bed shear stress and friction laws in general."

- **Hydrologic viewpoints:** Although authors frequently employed the term, hydrology, the hydrologic viewpoints are not well revealed. For example, According to the schematic view in Fig. 4, flow depth is always larger than the height of obstacles (rocks, boulders, vegetation, etc.). In hydrology viewpoint, the flow depth is mostly much smaller than that of larger elements.

This is true. As Fig. 4 was intended to be as generic as possible, we changed its legend to explain how it also covers the cases of small inundation ratios (though side views do not allow a proper representation of tortuous flows).

A phrase was added in the legend of Fig.4.

"The very small inundation ratios ($\Lambda_z < 1$) typical of overland flows in hydrology (flows through emergent obstacles, including vegetation) correspond to ε values larger than H values (tortuous flows are best seen in the top views of Fig.8)."

Overland flow often has a slope larger than 10 % up to 100 % (45 degree) in Fig. 5. The median slope 6 % in overland flow seems to be very mild.

We tend to agree but this is a consequence of the chosen classification of cases, esp. to separate overland flows from high-gradient flows on criteria that rely on water depth. On the one hand this may be seen as a bias in the representation, on the other hand larger water depths result in many differences in flow typology even for equivalent slopes and inundation ratios, which is another way to look at Fig.5.

In shields diagram, most data comes from the experiments in hydraulic conditions while data obtained from hillslope erosion studies seems not to be incorporated. Although most of theories in hillslope erosion was indeed borrowed from the results of river erosion studies, it would be better if authors mention about how those are different since many watershed-based, hydrology-viewpoint erosion studies have been done.

The point raised here is now mentioned in the introduction (L90-92) to state that we aim at gathering the hydrological and hydraulic points of view within the same normative framework.

"This signature is thought normative enough to facilitate comparisons between studies, encompassing both the hydrological (i.e. more "natural") and hydraulic (i.e. more "controlled") contexts."

However, a few words have been added in the comments to the Shields diagram (L953-955)

"To search for additional indications, the points in Fig.9 have been sorted by flow depths with the arbitrary $H=5$ cm threshold. Other case classifications may be relevant, for example identify the hydrological and hydraulic contexts."

- **Scale:** Another suggestion is about scale. I am wondering about “Problem/Domain” scale and “Resolving” scale for computations. Please refer to the below classification table, fill in blank if possible with authors’ language, and modify/reflect something useful for this review if any. Some part of description was used in Kim and Ivanov, *A holistic, multi-scale dynamic downscaling framework for climate impact assessments and challenges of addressing finer-scale watershed dynamics*. J. Hydrol. 522, 645–660, 2015. The most often used modelling in hydrology is based on watershed scale in which, as an example, length scale is say 10 km and time scale is event, seasonal, or annual. Authors used the term, hydrology in title and others, but there is no even typical context addressed in Table 1.
[Table 1 not shown]

- The issue of "domain scale" vs. "resolution scale" as also been raised by the Reviewer #3 (Major issue 1/) so the provided answer is copied hereunder.

We completely agree with this remark that "the end justifies the means" in the choice of modelling strategies. Following your suggestion (and these made by Reviewer #4) we acknowledge the idea that the objectives of a study show through the choice of appropriate subscales (say δL and δT) either directly cited in numerical procedures or as the spatial and temporal resolutions of data collection. Accordingly, we are now trying to correlate the choice of a model to the size of the domain (L, T, H) and to the size of its subscales (δL , δT) before examining the influence of flow typologies and dimensionless numbers.

The text of the manuscript has been changed in the abstract (L24, 26 & 32), the introduction (L84, 88, 90 & 151), the conclusion (L1013-1014 & 1033) and in several occurrences of the term "subscale" or in direct association (e.g. L540 & 544-545).

Figure 2 now has two subplots, Fig. 2a (unchanged) and Fig. 2b that positions modelling and/or data collection scales in the (δL , δT) plane for complementary indications, to gain a more comprehensive view. The added Fig.2b and its legend appear hereunder.

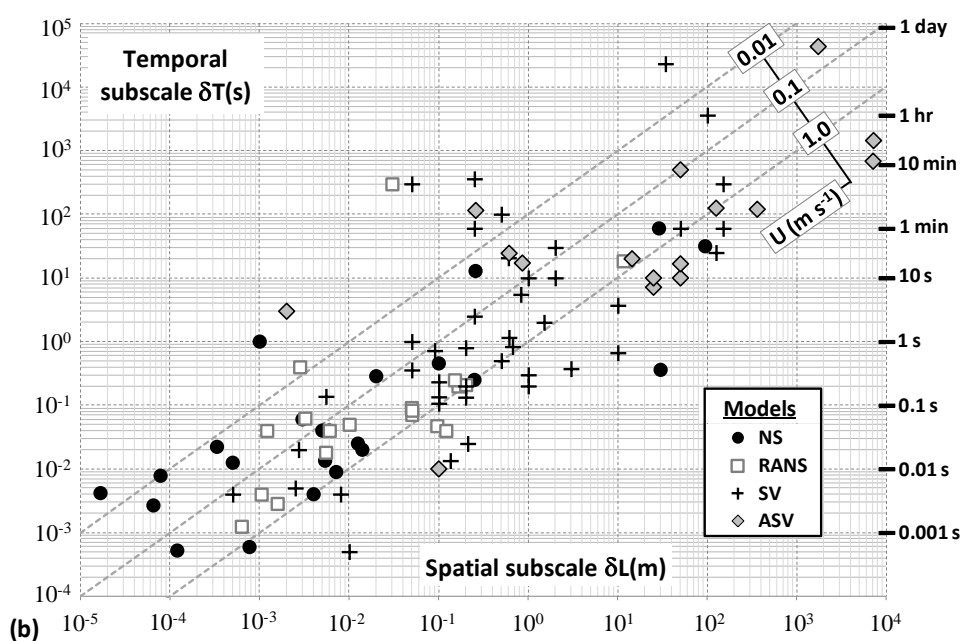


Figure 2 – How increasing (L, T) spatiotemporal scales (a) and (δL , δT) subscales (b) of the flow domain tend to be associated with decreasing complexity in the choice of flow models, sorted here into four levels of refinement: Navier-Stokes (NS), Reynolds-Averaged Navier-Stokes (RANS), Saint-Venant (SV) or Approximations to Saint-Venant (ASV). A transverse analysis involves forming L/T ratios, searching for clues to model selection according to these "system evolution velocities" or governed by flow typologies that would exhibit specific L/T ratios (a). Unit values of the Courant number ($Cr=U\delta T/\delta L$) have been used to trace characteristic flow velocities of $U=0.01$, 0.1 and 1 m s^{-1} and the indicative numerical stability criterion is $Cr \leq 1$: for given δL and U values, δT should lie behind the dotted line (b). Both plots were assembled from information available in the studies cited in Appendix A, selecting six textbook cases (sketches A to F, Table 1) for illustration (a).

- The Kim & Ivanov (2015) paper is now cited in the introduction (L115) but we could not include an additional table as we wanted to keep the structure of the paper unchanged (as suggested by Reviewers #1, 2 and 3).

- Table 1 was only intended to show the envelope of cases in Fig.2a with the associated numerical values: we do not think it is useful there to explicitly mention the differences between applications in hydrology and hydraulics. Conversely, we have tried to better address the hydrological contexts, elsewhere in the manuscript, following your suggestions. There are too many occurrences of the term "hydrology" to mention, and these are not always accompanied by the term "hydraulics".

- **Title:** I suggest title to use better words for "free-surface flow" and "hydrology". The impression of the term, free surface flow seems to be related to a topic for tracking and locating the free surface between water and air, but there is no discussions on this (e.g., volume of fluid technique). Also, the reviews of this study are more or less focused on the viewpoints of hydraulics and fluid mechanics, not hydrology. At least to me, I expected hydrology could be used when it is based on watershed scale. I agree watershed-based erosion models are also commented in the manuscript, but their relative importance might be less than 20-30 %. My tentative suggestion on the title is "*Determinants of modelling choices for 1-D surface flow and erosion problems in hydrology and hydraulics: a review*" If authors want to use "hydrology", I hope they can describe more on hydrology viewpoints.

- We partly agree with the suggestions here. We wish to keep "free-surface flow" because this terminology comes in opposition with "full" or "partially filled flows" (most often in ducts or circuits). From this point of view, dealing with "free-surface flows" means dealing with "environmental surface fluxes" but the former formulation seems more appropriate than the latter.

- We have added hydraulics in the title to assume the difference between the applications and scales more or less implicitly referred to by "hydrology" or by "hydraulics".

- Conversely, we deleted "in hydrology" from the title of Section 4.2 in order to regain more generality ("Research challenges and philosophy of modelling" applies to both the hydrological and the hydraulic points of view)

- Following the suggestion of Reviewer #3, we have also replaced "erosion" by "morphodynamics" (see the associated comments in this reply). The title now reads

"Determinants of modelling choices for 1-D free-surface flow and morphodynamics in hydrology and hydraulics: a review"

- Last but not the least, we can usually select a proper model in many hydrologic and hydraulic applications. For example, for watershed modelling, ASV with transport capacity concept; for flooding and river sedimentation, SV coupled with Exner equation; at much finer scale for structures NS-based erosion models has been widely used. I think there exists a certain rule of

thumb that everyone can agree to choose a model. Can authors present a couple of examples when people can misuse or select incorrectly numerical models?

We could certainly think of several examples in which numerical models have been misused or selected incorrectly, but this would first require the definition of good practices. This could be the object of a future, complementary paper (see Fig.11 in the conclusion).

In the present study, we rather focus on reporting the choices made and seeking their determinants. However, even if we do not examine the relevance of these choices, it is still interesting to identify the atypical choices. This is pretty much the object of Table 1 and that of the A, B, C, D, E and F points in Fig.2a, 3, 6a and 7a.

All typology authors mentioned exists at the same time within a larger watershed. How do authors present any suggestion or implication on watershed-based hydrology modelers? In this point of view, I contend a certain level of coupling will be necessary to address the details; it is associated that people are currently trying to combine many numerical models in many discipline (see below literature), which will be more facilitated in near future as computing power is increasing. L451-455

In our opinion the necessity for "a certain level of coupling" comes with the choice of a flow model in the list (NS, RANS, SV, ASV). The latter is influenced by the scales (L, T, H) while the choice of a coupling could rather be influenced by that of a resolution scale, i.e. the modelling subscales (δL , δT).

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Kim et al., 2013. Modeling erosion and sedimentation coupled with hydrological and overland flow processes at the watershed scale. Water Resour. Res. 49, 5134–5154.

MINOR COMMENTS

- L169: Please clarify what “this” refers to.

This section has been rewritten.

- In 3.1.1: For domain length (L), it makes sense to use length for river, but how did you compute the length scale for watershed (i.e., square shape)?

As the paper focuses on 1D modelling, the square shapes of watersheds are not explicitly treated. This is equivalent to say that the modeler has to explicitly indicate which of the 1D element(s) are addressed in the watershed when indicating the L, T and H scales.

- L700-702: Is it correct that Fig. 4e is not mentioned in the description of the line?

Yes this is correct.

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**Determinants of modelling choices for 1-D free-surface flow and
morphodynamicserosion issues in hydrology and hydraulics: a
review**

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Abstract

This review paper investigates the determinants of modelling choices, for numerous applications of 1-D free-surface flow and ~~morphodynamic~~~~erosion~~ equations in hydrology and hydraulics, across multiple spatiotemporal scales. We aim to characterize each case study by its signature composed of model refinement (Navier-Stokes: NS, Reynolds-Averaged Navier-Stokes: RANS, Saint-Venant: SV or Approximations of Saint-Venant: ASV), spatiotemporal scales and subscales (domain length: L from 1 cm to 1000 km; temporal scale: T from 1 second to 1 year; Flow depth: H from 1 mm to 10 m, spatial step for modelling: δL , temporal step: δT), flow typology (Overland: O, High gradient: Hg, Bedforms: B, Fluvial: F) and dimensionless numbers (Dimensionless time period T^* , Reynolds number Re , Froude number Fr , Slope S , Inundation ratio Λ_z , Shields number θ). The determinants of modelling choices are therefore sought in the interplay between flow characteristics, cross-scale and scale-independent views. The influence of spatiotemporal scales on modelling choices is first quantified through the expected correlation between increasing scales and decreasing model refinements (though modelling objectives also show through the chosen spatial and temporal subscales). ~~identifying then~~ Then flow typology appears a secondary but mattering determinant in the choice of model refinement. This finding is confirmed by the discriminating values of several dimensionless ~~numbers, that~~ numbers, which prove preferential associations between model refinements and flow typologies. This review is intended to help ~~each~~ modellers in positioning their ~~his (her)~~ choices with respect to the most frequent practices, within a generic, normative procedure possibly enriched by the community for a larger, comprehensive and updated image of modelling strategies.

Keywords

Free-surface flow, modelling strategy, cross-scale analysis, flow typology, dimensionless numbers.

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1 Introduction

Free-surface flow models cover a wide range of environmental and engineering applications, across multiple spatiotemporal scales, involving several levels of ~~through successive~~ flow aggregations ~~in the streamwise direction~~, over various bed topographies: these govern both the qualitative (flow typology) and quantitative (dimensionless numbers) flow characteristics. Each case study may thus be positioned along "streamwise scenarios" (from runoff initiation to the main rivers) from unequivocal indications of the spatiotemporal scales and subscales, flow typology and associated dimensionless numbers. This literature review investigates the determinants of choices made for 1-D free-surface flow and ~~erosion~~ morphodynamic modelling in hydrology and hydraulics, seeking links between contextual information (spatiotemporal scales, flow typologies, dimensionless numbers) and conceptual descriptions (data collection and/or calculation subscales, refinement of the flow equations or, equivalently, richness of the physical basis). The entire set of descriptors, *i.e.* model refinement, spatiotemporal scales and subscales, flow typology and dimensionless numbers, constitutes the signature of a study. ~~This signature is thought normative enough to facilitate which is the open normative comparisons procedure designed to allow comparisons~~ between studies, encompassing both the hydrological (*i.e.* more "natural") and hydraulic (*i.e.* more "controlled") contexts. ~~and to be fed by the community.~~

For the sake of ~~generality~~ genericity, this review addresses a wide range of spatiotemporal scales, starting at the smallest plot scales (spatial scale: domain length $L < 10$ m; time scale: duration of the process $T < 10$ s; flow depth: $H < 1$ cm, Fig.-1), those of runoff genesis, overland flow hydraulics and detailed particle-scale physics (Horton 1945, Emmett 1970, Feng & Michaelides 2002, Schmeeckle & Nelson 2003). The intermediate scales of catchment and hillslope processes are these expected to exhibit the widest variety of flow typologies thus modelling strategies (Croke & Mockler 2001, Parsons et al. 2003, Aksoy & Kavvas 2005, Mosselman 2012). The larger river basin scales ($L > 100$ km; $T > 10$ days; $H > 1$ m) are also handled here, relevant for river flow modelling, flood prediction and water resources management (Nash & Sutcliffe 1970, Rosgen 1994, Loucks & van Beek 2005) with regional surface-subsurface interactions (De Marsily 1986), non-point pollution,

fluvial sediment budgets and global biogeochemical cycles (Walling 1983, Milliman & Syvitski 1992, Syvitski & Milliman 2007).

On the ~~e~~Earth's surface, flow aggregation in the streamwise direction occurs across several geomorphic thresholds (Kirkby 1980, Milliman & Sivitsky 1992, Church 2002, Paola et al. 2009), through a succession of flow typologies (Emmett 1970, Grant et al. 1990, Rosgen 1994, Montgomery & Buffington 1997). Flow aggregation in space and time is described, through the width function and geomorphological unit hydrograph concepts (Kirkby 1976, Robinson et al. 1995, Agnese et al. 1998), under the angle of ~~connecting-scale~~ hydrological and sedimentological pathways (see the review by Bracken et al. 2013) or ~~questioningdebating~~ the merits of similitude laws ~~and these ofversus~~ upscaling ~~methodsissues~~ in the description of hydrological processes (Strahler 1956, Blöschl and Sivapalan 1995, Slaymaker 2006). ~~An-a~~Alternatives consists in examining the "scale matching" between available data and modelling aims (Lilburne 2002, Kim & Ivanov 2015) ~~and the possibility to use a more complicated model not only because it replicates what a simpler model would do, plus additional information, but also because it offers different, specific outcomes (e.g. Sloff & Mosselman 2012).~~ ~~This raises technical (contextual) as well as strategic (conceptual) issues, handled here from an-~~ With similar goals but a different framework, this study proposes an overview on the most popular modelling practices, confronting the theoretical refinement of flow models to the ~~specific-nominal~~ spatiotemporal scales and characteristics of the free-surface flows described ~~processes at play~~.

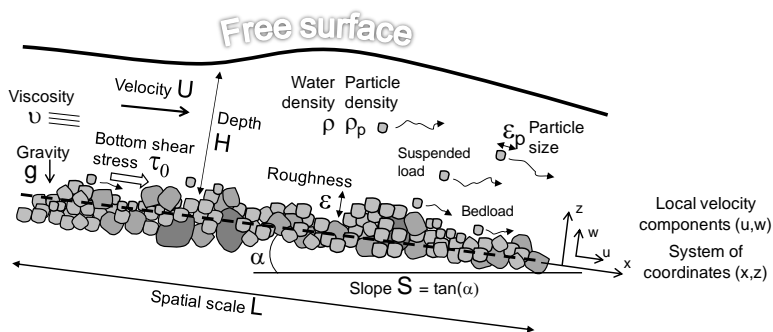


Figure 1 - Quantities most often used in the literature of free-surface flow and morphodynamic erosion modelling, with explicit reference to the (L, T, H) spatiotemporal scales of interest. This review is limited to 1D (x) spatial representations for simplicity, focusing on the streamwise (x) component of the mass and momentum conservation equations. The streamwise length (L) and velocity (U) suggest a natural time scale $T_0=L/U$ for the propagation of information, waves or perturbations, to be compared with the time scales (T) opted for in the literature.

Many papers or handbooks have summarised free-surface flow modelling and numerical techniques in hydraulics (King & Brater 1963, Abbott 1979, Cunge et al. 1980, Carlier 1980, French 1985) or hydrology (Chow 1959, Kirkby 1978, Beven 2000, Elga et al. 2015, Paniconi & Putti 2015) for various contexts, purposes and flow typologies. Less works have discussed the concern of *ad hoc* friction laws (Leopold et al. 1960, Gerbeau & Perthame 2001, Nikora et al. 2001, Roche 2006, Burguete et al. 2008), at the microscopic or macroscopic scales (Richardson 1973, Jansons 1988, Priezjev & Troian 2006, Smith et al. 2007, Powell 2014) although friction, flow retardation and energy dissipation processes are closely related to bedforms, thus plausibly govern flow typologies then, possibly, modelling choices. Often outside any focus on friction, numerous works have provided wide overviews on erosion modelling (Ritchie & McHenry 1990, Laflen et al. 1991, Merritt et al. 2003, Aksoy and Kavvas 2005, Boardman 2006). Morphodynamic Erosion models that lean on the most sophisticated flow models calculate explicit particle detachment, transport and deposition from velocity fields or flow energetics (Vanoni 1946, Hino 1963, Lyn et al. 1992, Mendoza & Zhou 1997) while most 1D or 2D physics-based models (e.g. Sloff et al. 2001, Vetsch et al. 2014) reduced complexity models either assume the "transport capacity" (Foster & Meyer 1972, Bennett 1974) or "transport distance" schools of thoughts (see details in Wainwright et al. 2008).

This multidisciplinary review (hydrology, hydraulics, fluid mechanics and morphodynamics erosion science) searches for the determinants of modelling choices. It focuses on hydrology but borrows from hydraulics and fluid mechanics, also when addressing morphodynamic erosion issues (erosion, transport and deposition of bed particles). The methodology consists in defining the "signature" of

each case study as the chosen model refinement and modelling subscales vs. the given spatiotemporal scales, flow typology, and dimensionless numbers, hypothesizing the conceptual element (model refinement and spatiotemporal subscales) is the consequence of the contextual elements (flow scales, typology and dimensionless numbers). The methodology consists in defining the “signature” of each case study as the chosen model refinement and the given flow typology, spatiotemporal scales and dimensionless numbers, hypothesizing the conceptual element (model refinement) is the consequence of the contextual elements. The paper is organized as follows: section 2 sorts the flow equations into four levels of refinement, section 3 plots these refinements versus the spatiotemporal scales of the studies, also depicting the influence of flow typologies and dimensionless numbers. Section 4 discusses the results and future research leads. Some of the best documented references among the cited literature have been gathered in Appendix A: most figures in this manuscript were plotted from this database.

2 Flow models

2.1 List of flow models

~~2.1.1 Water flow~~

Free-surface flow equations in the literature may roughly be sorted into four levels of decreasing refinement, i.e. depending on the number and nature of the indications included in their physical description from the richness of their physical basis. The choice made here (among many other possibilities) includes the Navier-Stokes equations (noted NS: Navier 1822, Stokes 1845), their average in time termed Reynolds-Averaged Navier-Stokes equations (RANS: Reynolds 1895, for turbulent flows), the depth-averaged Saint-Venant equations (SV: Saint-Venant 1871) and further approximations (referred to as ASV for Approximations to Saint-Venant), among which the Diffusive Wave Equation (DWE: Hayami 1951) and Kinematic Wave Equation (KWE: Iwagaki 1955, Lighthill & Whitham 1955).

2.1.2 ~~Erosion~~

In association with the flow equations, ~~The associated erosion~~ the equations ~~(not shown)~~ describing morphodynamic processes (particle erosion, transport and deposition) either issue from environmental fluid mechanics (e.g. Lyn 1987, Ribberink 1987, Elghobashi 1994) or from the ~~are based on a~~ representation of detachment and transport ~~more focused on hillslope processes~~ ~~on hillslopes~~ (Bennett 1974, Van Rijn 1984a, b, Wainwright et al. 2008), ~~arising from previous works in~~ streams (Einstein 1950) ~~and/or through the~~ channel networks (Du Boys 1879, Exner 1925, Hjulström 1935, Shields 1936, Bagnold 1956). Depending on the refinement of the coupled flow and morphodynamics models as well as on flow typology, a clear trend is that some elements are explicitly addressed whenever possible, e.g. particle advection and diffusion, while others are most often parameterised, e.g. particle detachment from excess bed shear stress and friction laws in general.

Friction is the link between water flow and erosion issues in terms of physical processes at play at the particle scale, or at the scale of the erodible bed asperities. On the one hand, this advocates the examination of erosion issues from the angle of decreasing refinements of the "flow and morphodynamics~~erosion~~" models seen as a whole (e.g. expecting the most complicated erosion processes to be out of reach of the simplest combined models). On the other hand, there might be a certain ~~inconsistency~~ ~~disconnection~~ between the refinement of the flow model and that of the chosen friction and erosion models, so the determinants of modelling choices should also be sought elsewhere: in flow typologies dictated by friction and flow retardation processes but also in "erosion ~~characteristics~~ ~~types~~", seen through a dimensionless descriptor (Section 3).

2.2 Navier-Stokes

2.2.1 Water flow

The Navier-Stokes (NS) equations have suitable simplifications for the shallow water cases ($L \gg H$) commonly used to describe free-surface flows. The three-dimensional fluid motion problem is reduced here to a two-dimensional description, whose projection along the streamwise axis writes:

$$\rho \left(\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + w \frac{\partial u}{\partial z} \right) + \frac{\partial p}{\partial x} = \rho g_x + \frac{\partial N}{\partial x} + \frac{\partial \tau}{\partial z} \quad (1)$$

where ρ is water density [ML^{-3}] assumed constant for incompressible flows, u is the local water velocity in x [LT^{-1}], t is time [T], x is the longitudinal distance [L], w is the local water velocity in z , z is the vertical coordinate [L], τ is the tangential stress in x [$\text{ML}^{-1}\text{T}^{-2}$], p is the local pressure [$\text{ML}^{-1}\text{T}^{-2}$], g_x is the projection of gravity g on x [LT^{-2}], N [$\text{ML}^{-1}\text{T}^{-2}$] is the normal stress in x (accounting for example for non-hydrostatic pressure effects) and τ [$\text{ML}^{-1}\text{T}^{-2}$] is the tangential stress in x , due to water [$\text{ML}^{-1}\text{T}^{-2}$] which is noted τ_0 on the bed in Fig.-1. The normal and tangential stresses also write $N = \mu \partial u / \partial x$ and $\tau = \mu \partial u / \partial z$, respectively, where μ [$\text{ML}^{-1}\text{T}^{-1}$] is the dynamic viscosity.

The Navier-Stokes equations stay valid throughout the full range of flow regimes, scales and contexts. They are preferentially used where much complexity is needed, often when relevant simplified flow descriptions could not be derived, for example for particle-scale applications (Chen & Wu 2000, Wu & Lee 2001, Feng & Michaelides 2002), overland flow (Dunkerley 2003, 2004) or flows over pronounced bedforms (Booker et al. 2001, Schmeeckle & Nelson 2003). A very wide review of numerical methods and applications for the NS equations is provided by Gresho & Sani (1998) and a benchmark of numerous solvers by Turek (1999). The general trend is that improvements in efficiency of the algorithms have approximately kept pace with exponential improvements in computer power over the past 50 years (Moore 1965, Mavriplis 1998, Koomey et al. 2010, Mosselman & Le 2016) which tends to push the limitations of numerical methods further away.

There are many turbulence models (e.g. DNS Direct Numerical Simulations, LES Large Eddy Simulations and RANS Reynolds Averaged Navier Stokes) suitable for free surface flow modelling (Katopodes & Bradford 1999). Direct Numerical Simulations explicitly resolve all turbulence scales at the cost of more than Re^3 calculations (Härtel 1996) while Large Eddy Simulations (Smagorinsky 1963, Leonard 1974) filter out the smallest scales and resolve only the larger ones. The RANS equations (Smith & McLean 1977, Rödi 1988) do not resolve any scale but the stress terms used for their closure have proven useful for the modelling of near bed turbulent patterns (see next subsection). The general trend is that improvements in efficiency of the algorithms have approximately kept pace with exponential improvements in computer power over the past 50 years (Moore 1965, Mavriplis 1998, Koomey et al. 2010) which tends to push the limitations of DNS and LES further away.

2.2.2 ~~Erosion~~ Morphodynamics

One of the earliest modern contributions on the rheology of two-phase flows is due to Einstein (1906) with the recognition that the viscosity of a mixture increases with the volumetric concentration of solid particles, at least for "slow flows". Brinkman (1947), Happel & Brenner (1965) then Leal (1980) studied the shearing strength of multiphase viscous flows while Batchelor (1974) and Russell (1981) addressed turbulent flows. Drew (1983) provided a general framework for the "*mathematical modelling of multiphase flow*", cited as a predecessor by Elghobashi (1994) who described particle-laden turbulent flows, discarding several assumptions (e.g. compressibility, phase change and thermodynamic effects) to yield a momentum conservation equation suitable for most natural flows and purposes:

$$\rho_k \left(\frac{\partial c_k u_k}{\partial t} + \frac{\partial c_k u_k^2}{\partial x} + \frac{\partial c_k u_k w_k}{\partial z} \right) + c_k \frac{\partial p_k}{\partial x} = \rho_k g_x + \frac{\partial c_k N_k}{\partial x} + \frac{\partial c_k \tau_k}{\partial z} + M_k \quad (2)$$

where the subscript k is an index for the phase (carrier: k=c, dispersed phase: k=d), c_k (-) is the local volumetric fraction ($c_c + c_d = 1$), u_k [LT^{-1}] and w_k [LT^{-1}] are the local velocities in x and z, respectively,

ρ_k [ML⁻³] is density, p_k [ML⁻¹T⁻²] is pressure, N_k [ML⁻¹T⁻²] and τ_k [ML⁻¹T⁻²] account for local non-hydrostatic pressure and shear stress effects, respectively, and M_k [ML⁻²T⁻²] is the momentum exchange term between phases. The exchange term vanishes for "one-way" couplings in which particles move in response to water motion (dispersed flows or dilute suspensions with $c_2 < 10^{-6}$) but should be kept for "two-way" couplings (dispersed flows with $10^{-6} < c_2 < 10^{-3}$ with non-negligible solid-fluid interactions, at the necessity of iterative resolution procedures) and also for "four-way" couplings (dense suspensions or collision-dominated flows with $c_2 > 10^{-3}$). In the latter case, additional models are needed to simulate particle-particle or particle-scale interactions (Nabi et al. 2012, 2013a,b) in the form of collisions, buoyancy and local pressure, drag or viscosity effects to be included in the above N_k and/or τ_k stresses (Drew 1983, Elghobashi 1994, Fernando 2012).

Several types of practical applications dictate the use of high-level formalisms in the description of particle detachment and transport, typically to handle explicit bed geometries and alterations (Colombini 2014, Kidanemariam & Uhlmann 2014), for example jet scours and regressive erosion (Stein et al. 1993, Bennett et al. 2000, Alonso et al. 2002), diverging sediment fluxes in canals (Belaud & Paquier 2001) or incipient motion conditions, calculated from grain size, shape and weight (Stevenson et al. 2002). The NS formalism is especially appropriate ~~also needed~~ to describe strong water-sediment couplings, *i.e.* couplings in which the solid phase exerts an influence on the liquid phase, acting upon velocity fields, flow rheology and erosive properties (Sundaresan et al. 2003). Such couplings may be sorted by increasing sediment loads, from dispersed multiphase flows (Parker & Coleman 1986, Davies et al. 1997) to density currents (Parker et al. 1986), hyperconcentrated flows (Mulder & Alexander 2001) and up to debris flows (Bouchut et al. 2003, Bouchut & Westdickenberg 2004), the latter derived as mathematical generalisations of the well-known Savage & Hutter (1989, 1991) avalanche models over explicit, pronounced topographies. Moreover, the NS formalism offers the possibility to work on the energy equations: the erosive power and transport capacity of sediment-laden flows may be estimated from the energy of the flow, ~~examining~~ ~~debating the case of~~ turbulence damping (or not) with increasing sediment loads (Vanoni 1946, Hino 1963, Lyn et al. 1992, Mendoza

& Zhou 1997). The matter is not **completely** free from doubt today (Kneller & Buckee 2001) **though** the diagram proposed by Elghobashi (1991, 1994, p310) to describe the regimes of interactions between particles and turbulence seems rather widely accepted. For the most dilute suspensions ($c_d < 10^{-6}$) the sediment load is not supposed to have any influence on turbulence characteristics. For the intermediate case ($10^{-6} < c_d < 10^{-3}$) the sediment load is supposed to enhance turbulence only if the particle response time is at least two orders of magnitude greater than the Kolmogorov time scale, i.e. the characteristic time for the turbulent eddies to vanish: for the same sediment load and water viscosity, larger particles tend to enhance turbulence while smaller particles tend to damp it. For dense suspensions ($c_d > 10^{-3}$) ~~and~~ frictional drag, abrasion due to impacts of the travelling particles and increased flow viscosity have been described prone to enhance the detachment capacities of loaded flows (e.g. Alavian et al. 1992, Garcia & Parker 1993).

2.3 Reynolds-Averaged Navier-Stokes

2.3.1 Water flow

There are many turbulence models (e.g. DNS-Direct Numerical Simulations, LES-Large Eddy Simulations and RANS-Reynolds-Averaged Navier-Stokes) suitable for free-surface flow modelling (Katopodes & Bradford 1999). Direct Numerical Simulations explicitly resolve all turbulence scales at the cost of more than Re^3 calculations (Härtel 1996) while Large Eddy Simulations (Smagorinsky 1963, Leonard 1974) filter out the smallest scales and resolve only the larger ones. The RANS equations (Smith & McLean 1977, Rödi 1988) do not resolve any scale but the stress terms used for their closure have proven useful for the modelling of near-bed turbulent patterns. ~~Reynolds-Averaged Navier-Stokes (RANS) equations are a turbulence model, using~~ The RANS equations are time-averaged equations of fluid motion, less generic than the NS formalism. The hypothesis behind these equations is that instantaneous pressure (p), stresses (N , τ) and velocities (u , w) may be decomposed into time-averaged and randomly fluctuating turbulent parts (e.g. $u = \bar{u} + u'$) assuming the temporal

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average of any turbulent fluctuations is zero, ~~which finally yields~~. The RANS formulation usually arising from the NS equations is:

$$\rho \left(\frac{\partial \bar{u}^2}{\partial x} + \frac{\partial \bar{u} \bar{w}}{\partial z} \right) + \rho g \frac{\partial H}{\partial x} = \rho g S + \frac{\partial \bar{N}}{\partial x} - \frac{\partial \bar{\rho u'^2}}{\partial x} + \frac{\partial \bar{\tau}}{\partial z} - \frac{\partial \bar{\rho u' w'}}{\partial z} \quad (32)$$

where ~~\bar{u} [$L T^{-1}$] and \bar{w} [$L T^{-1}$] are the time-averaged local water velocities in x and z, H is the flow depth [L] and S is the bed slope [-]~~ the hydrostatic approximation has been used for the pressure term together with the hypothesis of small bed slopes. In the above, \bar{N} accounts for the viscous (laminar) pressure stresses, $\bar{\rho u'^2}$ is the normal stress due to turbulence, $\bar{\tau}$ becomes the viscous shear stress and $\bar{\rho u' w'}$ is the (turbulent) Reynolds stress.

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In this formulation, the "Reynolds stress" term τ is of crucial importance for free-surface flow, friction and erosion modelling, especially for shallow flows, first because it is the closure term ~~$\tau = -\bar{\rho u' w'}$~~ and second because the Reynolds stresses have been closely related, in magnitude and direction, to the size and arrangement of bed asperities. The combined analysis of the relative magnitude of the u' and w' terms has become the purpose of "quadrant analysis" (Kline et al. 1967, Raupach 1981, Kim et al. 1987) that identifies the four cases of outward interactions (quadrant I: $u' > 0$, $w' > 0$), ejections (quadrant II: $u' < 0$, $w' > 0$), inward interactions (quadrant III: $u' < 0$, $w' < 0$) and sweeps (quadrant IV: $u' > 0$, $w' < 0$). Depending on the submergence and geometry of bed asperities, the maximal Reynolds stresses, those with significant effects on flow structure, have most often been reported to occur near or just above the roughness crests (see Nikora et al. 2001, Pokrajac et al. 2007 and the review by Lamb et al. 2008a).

2.3.2 Morphodynamics ~~Erosion~~

Comparative reviews of RANS-level approaches to modelling sediment-laden two-phase flows within various two-way couplings have been performed by Bombardelli & Jha (2008) then Jha &

Bombardelli (2009), assessing the performances of "standard sediment transport models" (an advection-turbulent diffusion equation for the liquid-solid mixture), "partial two-fluid models" (distinct momentum conservation equations for the dispersed phase and the carrier phase, the latter seen as a liquid-solid mixture) and "complete two-fluid models" (general balance equations for both phases, inherited from the previous NS formulations) versus "Reynolds stress models" (expressing closure terms in function of the turbulent kinetic energy). The momentum balance in x for 1D approaches is the same for the dispersed phase in the complete and partial two-fluid models (Bombardelli & Jha 2008):

$$\rho_d \left(\frac{\partial c_d \bar{u}_d}{\partial t} + \frac{\partial c_d \bar{u}_d \bar{w}_d}{\partial z} \right) = \rho_d c_d g S - \frac{\partial \rho_d c_d \bar{u}'_d \bar{w}'_d}{\partial z} + F_D \quad (4)$$

where F_D [$\text{ML}^{-2}\text{T}^{-2}$] is the drag force term that allows two-way couplings, most often written as $F_D = 0.5 \rho_m C_D A (\bar{u}_c - \bar{u}_d)^2$ where ρ_m [ML^{-3}] is the density of the two-phase mixture, C_D (-) is the drag coefficient and A [L^2] is the cross-sectional area of the particles.

In their paper on movable river beds, Engelund & Fredsoe (1976) ~~judiciously~~ reformulated and exploited the existing hypotheses (Einstein & Banks 1950, Bagnold 1954, Fernandez Luque & van Beek 1976) of a partition between “tractive” destabilizing shear stresses and “dispersive” equalizing drags. The vertical concentration profiles of bedload and suspended load were calculated from incipient sediment motion conditions, relating stresses on the particles to the values and variations of near-bed velocities. One step further, the physical explanation, mathematical definition, point of application, main direction and erosive efficiency of the turbulent near-bed stresses have become an interesting feature ~~private hunting grounds~~ of the RANS models throughout the years (Nikora et al. 2001, Nino et al. 2003).

The maximal Reynolds stresses are located near the crests of the submerged bed asperities, where turbulent velocity fluctuations reach several times the average near-bed velocity values, which greatly enhances particle detachment (Raupach et al. 1991, Nikora & Goring 2000, Lamb et al. 2008a). Very

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few studies deal with the magnitude and point of application of the Reynolds stresses for partial inundation cases (Bayazit 1976, Dittrich & Koll 1997, Carollo et al. 2005) although turbulent flows between emergent obstacles often occur in natural settings. Particle detachment is generally attributed to “sweeps” (quadrant IV: $u'>0$, $w'<0$) (Sutherland 1967, Drake et al. 1988, Best 1992) or “outward interactions” ($u'>0$, $w'>0$) (Nelson et al. 1995, Papanicolaou et al. 2001) but depends on bed geometries and bed packing conditions. Finally, the RANS equations allow explicit calculations of shear stresses and particle-scale pick-up forces, thus incipient motion conditions (Nino et al. 2003, Afzalimehr et al. 2007). They may handle the movements of detached particles in weak transportation stages (Bounvilay 2003, Julien & Bounvilay 2013) down to near-laminar regimes (Charru et al. 2004).

2.4 Saint-Venant

2.4.1 Water flow

The Saint-Venant (SV) equations are obtained by depth-integrating the Navier–Stokes equations, neglecting thus the vertical velocities as well as vertical stratifications in the streamwise velocity (Stoker 1958, Johnson 1998, Whitham 1999). The SV equations also termed “shallow water equations” assume the $H \ll L$ hypothesis of shallow water which limits the admissible free-surface slope and implies a quasi-hydrostatic pressure distribution over the vertical. The integration process from NS to SV (Chow 1959, Abbott 1979) incorporates an explicit bottom friction term τ_0 that previously appeared only as a boundary condition in the NS and RANS equation:

$$\frac{\partial U}{\partial t} + U \frac{\partial U}{\partial x} + g \frac{\partial H}{\partial x} = gS + \frac{\tau_0}{\rho H} \quad (53)$$

Recent attempts have been made in the field of fluid mechanics to derive specific expressions for τ_0 (laminar flows: Gerbeau & Perthame 2001, macro-roughness: Roche 2006, thin flows: Devauchelle et al. 2007, turbulent flows: Marche 2007, multi-layer SV model: Audusse et al. 2008). However, the common practice in ~~hydraulics and~~ hydrology and hydraulics is rather to approximate steady-state

equilibrium between bottom friction τ_0 and the streamwise stress exerted at the bottom of a water column ($\tau_0 = \rho g H S_f$) to reach the popular formulation:

$$\frac{\partial U}{\partial t} + U \frac{\partial U}{\partial x} + g \frac{\partial H}{\partial x} = g(S - S_f) \quad (64)$$

(i) (ii) (iii) (iv) (v)

where (i) is the unsteadiness term, (ii) the convective acceleration term, (iii) the pressure gradient term, while (iv), (v) and (vi) form the diffusive wave approximation (later discussed).

In the above, S_f (-) is the “friction slope” whose expression depends on flow velocity and on the chosen friction law, often one of the Chézy, Darcy-Weisbach or Manning formulations (*e.g.*, $S_f = nU^2/8gH$ with Manning’s n friction coefficient). The derivation of the SV equations by Boussinesq (1877) involved a momentum correction coefficient β [-] in the advection term (King & Brater 1963, Chen 1992) to account for stratification effects in the vertical distribution of velocities, especially plausible in sediment-laden flows or in presence of density currents.

The SV equations may account for flows of variable widths and depths, for example in floodplains (Bates & De Roo 2000, Beltraos et al. 2012), rivers (Guinot & Cappelaere 2009), overland flow (Berger & Stockstill 1995, Ghavasieh et al. 2006, Kirstetter et al. 2016), overpressure in drainage systems (Henine et al. 2014), man-made channels (Zhou 1995, Sen & Garg 2002, Sau et al. 2010), vegetation flushing (Fovet et al. 2013), channel networks (Choi & Molinas 1993, Camacho & Lees 1999, Saleh et al. 2013), on benchmarks (Dimitriadis et al. 2016), interaction with subsurface (Pan et al., 2015), or natural settings (Moussa & Bocquillon 1996a, Wang & Chen 2003, Roux & Dartus 2006, Burguete et al. 2008, Bates et al. 2010), including these with curved boundaries (Sivakumaran & Yevjevich 1987). Discharge and cross-sectional area may conveniently be used instead of velocity and water depth, and the two equations describing mass and momentum in the Saint-Venant system now write (Sivapalan et al. 1997):

$$\frac{\partial A}{\partial t} + \frac{\partial Q}{\partial x} = q_a \quad (75)$$

$$\frac{1}{gA} \frac{\partial Q}{\partial t} + \frac{1}{gA} \frac{\partial}{\partial x} \left(\beta \frac{Q^2}{A} \right) + \frac{\partial H}{\partial x} + S_f - S = 0 \quad (86)$$

where A is the cross-sectional area [L^2], Q is the discharge [$L^3 T^{-1}$], q_a is the lateral flow per unit channel length [$L^2 T^{-1}$]. The magnitudes of the various terms in equations (5) and (6) are given in the literature (e.g. Henderson 1966, Kuchment 1972).

2.4.2 Morphodynamics ~~Erosion~~

In the hydro-morphodynamics ~~hydrology-erosion~~ community, the SV level is that of the *Concepts of mathematical modelling of sediment yield* by Bennett (1974). This landmark paper extended Exner's (1925) conservation of sediment mass, adding the possibility to handle different fluid and particle velocities, also accounting for particle dispersion *via* a diffusion term:-

$$\frac{\partial H c_d}{\partial t} + (1 - \phi_0) \frac{\partial z_0}{\partial t} + \frac{\partial H c_d U_d}{\partial x} = \frac{\partial}{\partial x} \left(H \eta_d \frac{\partial c_d}{\partial x} \right) \quad (9)$$

where ϕ_0 (-) is bed porosity, z_0 (-) is the bed level, U_d [$L T^{-1}$] is the spatial average of particle velocity over the cross section of the flow and η_d [$L^2 T^{-1}$] is a diffusivity coefficient. See for example Ancy & Heyman (2010) and Ballio et al. (2014) for the various possible formulations of the sediment continuity equation and associated numerical aspects, depending on the strength of the intended coupling with the carrier phase. The authors rather belong to the fluid mechanics-type of use of the SV equations, for hydro-environmental applications that necessitate taking maximum advantage of the level of details offered by (9), often by using SV-level formulations of the Exner equation in combination with RANS or NS-level flow models (e.g. Riberink 1987, Blom 2008, Sloff & Mosselman 2012).

Conversely ~~Unfortunately~~, in the field of hydrology, numerous ~~most~~ citing papers discard one or several terms from the Bennett (1974) equations ~~this term~~, typically taking particle velocity equal to

water velocity. The assumption seems false if transport occurs as bedload or saltation load, questionable for suspended load trapped into turbulent motions, exact only for very small particles borne by laminar flows. Although warning against the capability of first-order laws to “*represent the response of sediment load to changes in transport and detachment capacity*” (Bennett 1974, p.491), the author recommended the use of such a model (Foster and Meyer 1972). The proposed simplification writes $e/D_c = 1 - c/T_c$, where the net erosion rate (e) is normalised by the maximal detachment capacity (D_c) while sediment load (c) is normalised by the maximal transport capacity of the flow (T_c). An additional (uncertain) hypothesis was that of maximal detachment capacity for minimal sediment load, *i.e.*, clear water. See the controversial comments around the Wainwright et al. (2008) paper: the areas of disagreement revolve around the ability of models to handle unsteady flow conditions, to deal with suspended and/or bedload transport, to consider particles of different sizes and to stay valid over realistic ranges of sediment concentration.

Those questions directly address the possibilities of SV-level approaches. Higher-level models (NS, RANS) better address the dynamics of incipient motion (Dey & Papanicolaou 2008), especially in shallow laminar flows (Charpin & Myers 2005) or focusing on granular flows (Parker 1978a, b, Charru et al. 2004, Charru 2006). Refined models are also needed to explicitly handle specific particle velocities (Bounvilay 2003), to describe particle diffusion in secondary currents (Sharifi et al. 2009), to account for the spatial heterogeneity of “neither laminar nor turbulent” overland flows (Lajeunesse et al. 2010) or to introduce modifications in flow rheology (Sundaresan et al. 2003). On the other hand, many erosion controls have received attention within the SV or ASV formalisms, *i.e.* without explicit descriptions of particle-scale flow features: micro-scale variability (Risse et al. 1993, Kinnell et al. 2005), local sheltering effects (Nearing et al. 2007, Kim & Ivanov 2014), slope effects (Polyakov & Nearing 2003), particle-size effects (Van Rijn 1984a, Hairsine & Rose 1992a, Sander et al. 2007, Wainwright et al. 2008), flow stratification effects (van Maren 2007), the effects of hyperconcentrated flows (Hessel 2006). ~~and the B~~ bedload transport (*e.g.* Van Rijn 1984b, Julien & Simmons 1985,

Hairsine & Rose 1992b, Wainwright et al. 2008) has also motivated the search for dedicated formalisms. ~~have received much attention within the SV or ASV formalisms.~~

Whatever the liquid-solid coupling opted for, the SV level covers the widest variety of contexts, from overland erosion models (Simpson & Castelltort 2006, Nord & Esteves 2010, Stecca et al. 2015) to dam-break hydraulics over erodible beds (Cao et al. 2004) and the analysis of channel inception driven by the variations of the Froude number (Izumi & Parker 1995) or the impact of travelling particles (Sklar & Dietrich 2004, Lamb et al. 2008b). Sediment detachment and transport over plane beds (Williams 1970), rough beds (Afzalimehr & Anctil 1999, 2000, Gao & Abrahams 2004), channels (Villaret et al. 2013, 2016), step-pools (Lamarre & Roy 2008) or pool-riffle sequences (Sear 1996, Rathburn & Wohl 2003) have yielded often-cited studies, while sediment flushing in reservoirs (Campisano et al. 2004) and vegetation flushing in canals (Fovet et al. 2013) constitute more specific applications. Cited limitations of the SV approaches are their inability to explicitly describe the near-bed velocity fluctuations, especially the local accelerations responsible for particle entrainment but also the vertical gradients of the streamwise velocity, for bedload transport in the laminar layer. This lack of accuracy in the description of flow characteristics also endangers the possibility to predict the formation, transformation and migration of geometrical bed patterns, which in turn requires the full set of 3D (x, y, z) NS equations in several cases (Lagrée 2003, Charru 2006, Devauchelle et al. 2010).

There seems to exist a dedicated "NS-SV Morphodynamics" research lead that uses rather simple bedload transport formulae (Du Boys 1890, Meyer-Peter & Müller 1948, Einstein & Banks 1950, Bagnold 1966, Yalin 1977) to calculate sediment fluxes from excess bed shear stresses, in studies of long-term system evolutions. These low "system evolution velocities" appear under the "quasi-static" flow hypothesis: particle velocity may be neglected before water velocity, which allows neglecting the unsteadiness term in the momentum equation but on no account in the continuity equation (Exner law) that describes bed modifications (Parker 1976). Although derived for turbulent natural flows, ~~Moreover,~~ shear stresses may also be ~~are generally~~ calculated from near-bed laminar or near-laminar

velocity profiles, sometimes with the regularising hypothesis that detachment and transport occur just above the criterion for incipient motion (see the review by Lajeunesse et al 2010). Various applications address rivers with mobile bed and banks (Parker 1978a, b), focus on self-channelling (Métivier & Meunier 2003, Mangeney et al. 2007) and often resort to formulations at complexity levels between these of the NS and the SV approaches (Devauchelle et al. 2007, Lobkovsky et al. 2008).

2.5 Approximations to Saint-Venant

2.5.1 Water flow

When the full Saint-Venant equations are not needed or impossible to apply due to calculation time a lack of data, an option is to neglect one or several terms of the momentum equation (Ponce and Simons 1977, Romanowicz et al. 1988, Moussa & Bocquillon 1996a, Moussa & Bocquillon 2000, Rousseau et al. 2015). In most practical applications for flood routing, the unsteadiness (i) and convective acceleration (ii) terms in (4) may be neglected, suppressing the first two terms from (6). Combining the remaining terms in (5) and (6), we obtain the Diffusive Wave equation (Moussa, 1996):

$$\frac{\partial Q}{\partial t} + C \left(\frac{\partial Q}{\partial x} - q_a \right) - D \left(\frac{\partial^2 Q}{\partial x^2} - \frac{\partial q_a}{\partial x} \right) = 0 \quad (107)$$

where C [LT^{-1}] and D [L^2T^{-1}] are non-linear functions of the discharge Q (and consequently the flow depth H) known as the celerity and diffusivity, respectively.

In cases where the pressure-gradient term (iii) in (4) can also be neglected, the third term of (6) also vanishes and the Diffusive Wave becomes the Kinematic Wave equation, with $D=0$ in (7). The Diffusive Wave in the historic formulations (Cunge 1969, Akan & Yen 1981) or in more recent works (Rutschmann & Hager 1996, Wang et al. 2006, Wang et al. 2014, Cimorelli et al. 2015, Swain & Sahoo 2015) can thus be considered a higher order approximation than the Kinematic Wave

approximation (Katopodes 1982, Zoppou & O'Neill 1982, Daluz Vieira 1983, Ferrick 1985, Ponce 1990). Both have been largely studied (since Wooding 1965a,b, Singh 1975, Lane & Woolhiser 1977, Ponce 1991) until more recently (Szymkiewicz & Gasiorowski 2012, Yu & Duan 2014) and have proven very useful for canal control algorithms (Rodellar et al. 1993) or flood routing procedures, with lateral inflow (Fan & Li 2006), in rectangular channels (Keskin & Agiralioglu 1997), for real time forecast (Todini & Bossi 1986), in lowland catchments (Tiemeyer et al. 2007), for overland flows (Pearson 1989, Chua et al. 2008, 2010, 2011), on urban catchments (Gironás et al. 2009, Elga et al. 2015), for small catchments (Moussa et al. 2002, Chahinian et al. 2005, Charlier et al. 2007), for mountainous catchments (Moussa et al. 2007), for medium size catchments (Emmanuel et al. 2015) or tropical catchments (Charlier et al. 2009), at the largest scale of the Amazon basin (Trigg et al. 2009, Paiva et al. 2013), for anthropogenic hillslopes (Hallema & Moussa 2013), to address backwater effects (Munier et al. 2008), stormwater runoff on impervious surfaces (Singh 1975, Pearson 1989, Blandford & Meadows 1990, Parsons et al. 1997), stream-aquifer interactions (Perkins & Koussis 1996) or volume and mass conservation issues (Perumal & Price 2013). Given their "nominal" scales of application, the ASV models are sometimes fed by airborne (remote sensing) data acquisition (Jain & Singh 2005, Reddy et al. 2007). In addition, predictive uncertainties (Elhanafy et al. 2008) or the applicability of the kinematic and diffusive wave equations are the main scope of several studies (Liggett & Woolhiser 1967, Ponce & Simons 1977, Ponce et al. 1978, Moussa & Bocquillon 1996b, Bajracharya & Barry 1997), the evaluation of modelling strategies is that of Horritt & Bates (2002), while parameter estimation is addressed, among others, by Koussis et al. (1978).

2.5.2 Morphodynamics ~~Erosion~~

Whereas common practices in fluid mechanics and hydraulics are rather to seek context-specific strategies in morphodynamic ~~erosion~~ modelling, two simplifying and unifying trends, if not paradigms, have developed in the field of hydrology. The first one is the transport capacity concept (Foster & Meyer 1972) in which the erosive strength of the flow decreases with increasing suspended sediment

load, until a switch occurs from detachment- to transport-limited flows. The second one is the stream power concept (Bagnold 1956) that *slope times discharge* is the explicative quantity for erosion, with adaptations that mentioned unit stream power (*slope times velocity*, Yang 1974, Govers 1992) or fitted exponents to the slope and discharge terms (Julien & Simmons 1985).

However, in all cases where the volumetric concentration of the dispersed phase is difficult to know, a possible surrogate is the division of the sediment mixture into size fractions with specific erosion and transport properties (Einstein 1950, Egiazaroff 1957, Hirano 1970, Day 1980, Ribberink 1987) possibly expressed as specific travel distances (Kirkby 1991, 1992, Parsons et al. 2004, Wainwright et al. 2008). The latter presents the following formulation of sediment continuity:

$$\frac{\partial h_{s,\varphi}}{\partial t} + \frac{\partial q_{s,\varphi}}{\partial x} - \varepsilon_{\varphi} + d_{\varphi} = 0 \quad (11)$$

where the subscript φ represents "size- φ " sediments, $h_{s,\varphi}$ [L] is the equivalent depth of sediment transport per unit width of the flow, $q_{s,\varphi}$ [$L^2 T^{-1}$] is the unit discharge of sediment, ε_{φ} [$L T^{-1}$] is the rate of erosion of the surface and d_{φ} [$L T^{-1}$] is the rate of deposition. This equation is more general than the sediment continuity equation most often used in combination with ASV flow models,

$$\frac{\partial A c_d}{\partial t} + \frac{\partial Q c_d}{\partial x} - E = 0 \quad (12)$$

where E [$L^2 T^{-1}$] is the areal erosion rate.

Many catchment-scale hydrology-erosion models (*e.g.* ANSWERS: Beasley et al. 1980, CREAMS: Knisel 1980, KINEROS: Smith et al. 1995, LISEM: De Roo et al. 1996, WEPP: Ascough et al. 1997, EUROSEM: Morgan et al. 1998, MAHLERAN: Wainwright et al. 2008, MHYDAS-Erosion: Gumiere et al. 2011b, Gregoretti et al. 2016, Hould-Gosselin et al. 2016) adopt the 1D Diffusive or Kinematic Wave Equations to route water fluxes, possibly through vegetated strips (Muñoz-Carpena et al. 1999), together with the simplest possible couplings between water and sediment fluxes (Aksoy & Kavvas 2005).

544 |
545 A known difficulty when embracing larger scales with simplified models is to describe the
546 spatially-distributed sources and sinks of sediments (Jetten et al. 1999, 2003) with or without explicit
547 descriptions of the permanent or temporary connectivity lines, for water and sediment movements
548 (Prosser & Rustomji 2000, Croke & Mockler 2001, Pickup & Marks 2001, Bracken et al. 2013). What
549 | tends to force reduced complexity approaches in [most catchment-scale](#) erosion models is the necessity
550 to handle distinct detachment, transport and deposition processes (from the very shallow diffuse flows
551 formed during runoff initiation to the regional-scale basin outlets) with only sparse data on flow
552 structure and soil characteristics (cohesion, distribution of particle sizes, bed packing). Parsons &
553 Abrahams (1992) have established how the agronomical, engineering and fluvial families of
554 approaches have converged into similar modelling techniques, especially on the subject of erosion in
555 overland flows (Prosser & Rustomji 2000). The ASV formalism also allows fitting bedload transport
556 formulae against mean discharge values as a surrogate to the overcomplicated explicit descriptions of
557 erosion figures in high-gradient streams with macro-roughness elements (Smart 1984, Aziz & Scott
558 1989, Weichert 2006, Chiari 2008). ASV-level couplings have also been applied to study the slope
559 independence of stream velocity in eroding rills (Gimenez & Govers 2001) and the appearance of bed
560 patterns in silt-laden rivers (van Maren 2007).

561

562 **3 Determinants of modelling choices**

563 This section aims at the construction of a signature for each case study, relating the "conceptual"
564 choice of a model refinement (Navier-Stokes: NS, Reynolds-Averaged Navier-Stokes: RANS, Saint-
565 Venant: SV or Approximations to Saint-Venant ASV) to the "contextual" descriptors, *i.e.* the
566 spatiotemporal scales (section 3.1), spatiotemporal scales and flow typologies (section 3.2),
567 spatiotemporal scales, flow typologies and dimensionless numbers (section 3.3). Figures 2, 3, 5, 6 and
568 | 7 in this section were drawn from the ~~158~~179 studies listed in Appendix A.

3.1 Spatiotemporal scales

3.1.1 Influence of domain length (L) and time scale (T)

A cross-disciplinary analysis of the cited literature indicates a clear correlation between the (L , T) spatiotemporal scales on one side and the chosen model refinement (NS, RANS, SV or ASV) with the (δL , δT) spatiotemporal subscales (data collection and/or numerical schemes) on the other side. In this (L , T) plane, Fig.-2a quantifies the expected trend that sophisticated (NS, RANS) models are required to represent rapidly-varying small-scale phenomena (lower left) while simplified approaches (ASV) pertain to increased durations and spatial extensions (upper right). The same pattern is visible in Fig.2b for the (δL , δT) subscales, reporting a strong correlation between the choice of a model and the size of the modelling subscales, for given (L , T) values. Typical scales of application may be identified for each model refinement: NS (10 cm< L <100 m, 10 s< T <1 hr), RANS (1 m< L <100 m, 10 s< T <1 hr), SV (10 m< L <20 km, 1 min< T <5 days) and ASV (10 m< L <1000 km, 30 min< T <1 yr). However, some studies consider larger spatial or temporal scales, for example Charru et al. (2004) for overland granular flows (RANS, L ~20 cm, T ~2 days) or Rathburn & Wohl (2003) for pool-riffle sequences (SV, L ~70 m, T ~30 days). Nevertheless, the existence of overlap regions suggests that the (L , T) spatiotemporal scales are not the only factor governing the choice of flow models.

The influence of flow typologies is discussed later in details but could the modelling choices be dictated by the scientific background of the modeller? A striking example is that of the SV models, responsible for the largest overlaps in Fig.-2. They may for example be used by physicists, as an upgraded alternative to the NS equations, in the field of environmental fluid mechanics (for limited scales). They may as well be convenient for soil scientists interested in high-resolution hydrology or for civil engineers who may need to cope with flow unsteadiness to handle morphodynamic erosion issues or to allow correct sizing of the man-made structures (for somewhat wider scales).

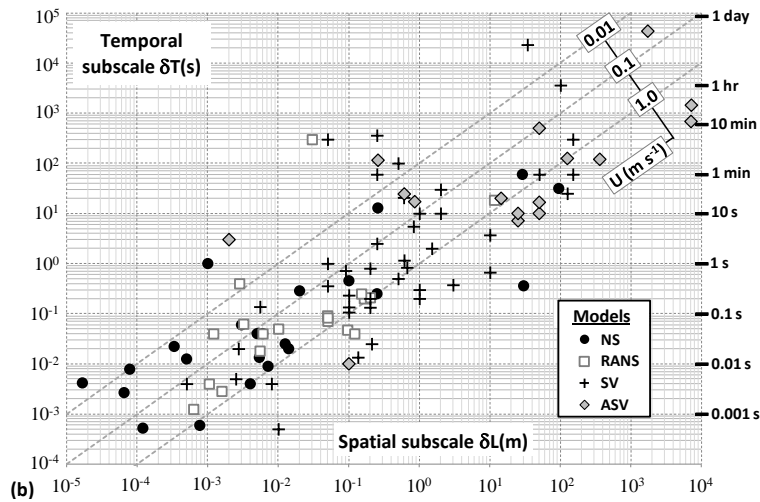
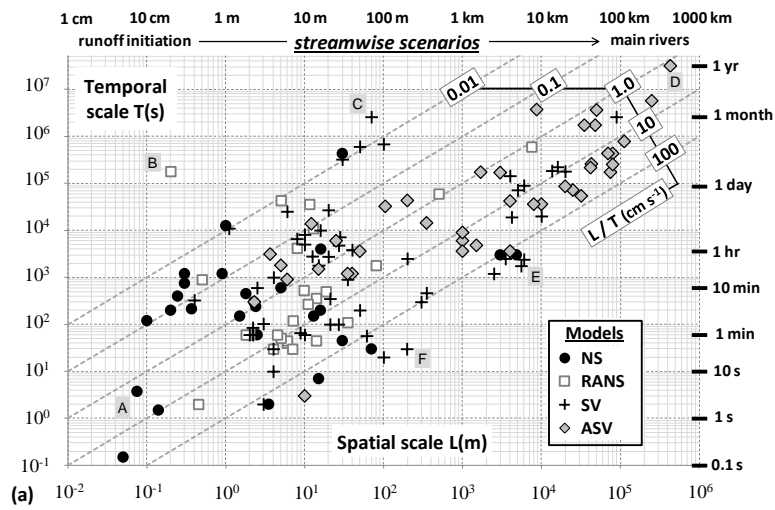


Figure 2 – How increasing (L, T) spatiotemporal scales (a) and $(\delta L, \delta T)$ subscales (b) of the flow domain tend to be associated with decreasing complexity in the choice of flow models, sorted here into four levels of refinement: Navier-Stokes (NS), Reynolds-Averaged Navier-Stokes (RANS), Saint-Venant (SV) or Approximations to Saint-Venant (ASV). A transverse analysis involves forming L/T ratios, searching for clues to model selection according to these "system evolution velocities" or governed by flow typologies that would exhibit specific L/T ratios (a). Unit values of the Courant number ($Cr=U\delta T/\delta L$) have been used to trace characteristic flow velocities of $U=0.01, 0.1$ and 1 m s^{-1} and the indicative numerical stability

criterion is $Cr \leq 1$: for given δL and U values, δT should lie behind the dotted line (b). Both ~~This figure plots~~
~~were~~~~was~~ assembled from information available in the studies cited in Appendix A, selecting six textbook
cases (sketches A to F, Table 1) for illustration (a).

Figure 2a bears another type of information than the trend to decreasing model refinement with
increasing spatiotemporal scales. As the x-ordinate indicates the spatial scale L and the y-ordinate the
time scale T , then the L/T ratio has dimensions of a velocity. However, this quantity should not be
interpreted as a flow velocity. It rather indicates which of the temporal (long-term, low L/T ratio) or
spatial (short-term, high L/T ratio) aspects are predominant in the study. Hence, the five dotted
diagonals ($L/T=10^{-4}$, 10^{-3} , 10^{-2} , 0.1 and 1 m s^{-1}) establish the numerical link between the spatial and
temporal scales of the cited experiments. They also show the dispersion with respect to the expected
(say "natural") correlation between increasing L and T values. ~~This dispersion contains a lot of~~
~~information.~~ Judging from the plotted literature, the lowest L/T ratios (e. g. 10^{-4} m s^{-1}) tend to indicate
systems with low "evolution velocities", possibly associated with long-term changes or effects (high T
values, low L values) obtained from repeated phenomena, multiple cycles and progressive
modifications. By contrast, high L/T ratios (e.g. 1 m s^{-1}) rather refer to single-event situations, more
associated with quick modifications of flow patterns or bed morphologies. Most applications find
themselves in the $10^{-2} < L/T < 10^2 \text{ cm s}^{-1}$ range, exhibiting no clear difference between the NS, RANS,
SV or ASV refinements. Conversely, this indicates that each level of refinement has been used to
model high or low system evolution velocities, sometimes by relying on specific (adapted or
upgraded) formulations of the systems of equation (see for example the hybrid NS-SV level of
refinement needed for detailed morphodynamics, especially to reproduce the long-term evolution of
bed topography).

If rules of thumb in problem dimensioning were to be drawn from Fig. 2a, geomorphological
concerns (dune migration, basin sedimentation, long-term bed modifications) probably require
stretching up the temporal scale so that low "system evolution velocities" would fall beneath $L/T=10^{-2}$
 m s^{-1} while event-based modelling (dam breaks, formative discharges, flash floods) should be able to

632 handle high "system evolution velocities" near or beyond $L/T=1 \text{ m s}^{-1}$. This "fixed-L, chosen-T"
633 description of system evolution and characteristic time scales also refers to Fig.-1 and Fig.2b in which
634 the choice of T and that of δT are somehow left at the modeller's discretion, as a degrees of freedom:
635 how different from T_0 should T be to allow long-enough observation and/or simulation periods?
636 These points are the subject of detailed investigations in the field of morphodynamics (Paola et al.
637 1992, Howard 1994, Van Heijst et al. 2001, Allen 2008, Paola et al. 2009). Indicators of "system
638 evolution velocities" with units of a velocity but different definitions may for example be found in
639 Sheets et al. (2002), who took the channel depth (H) divided by the average deposition rate to obtain a
640 relevant, characteristic time scale (T). For the same purpose, Wang et al. (2011) took the characteristic
641 bed roughness (ϵ) instead of channel depth. The objective is often to discriminate what Allen (2008)
642 called the "reactive" (high L/T) and "buffer" (low L/T) systems. With or without morphodynamic
643 erosion issues, a reasonable hypothesis here seems that the dispersion in L/T ratios arises from the
644 variety of flow contexts, which may necessitate different modelling strategies. In other terms, it is
645 deemed in this study that this secondary trend, associated with flow typologies, is also a determinant in
646 the choice of the flow model.

647
648 To take a few examples and guide the reader through the arguments and the figures of this paper,
649 Table 1 gathers the information available for the six textbook cases outlined by sketches A to F in
650 Fig.2a. The selected studies represent a wide variety of cases (drawing an approximate envelop of
651 cases in the L-T plane of Fig.2a) followed in the forthcoming stages of the analysis and associated
652 figures in Section 3.1.2 (determinants of modelling choices in the L-H plane, Fig.3), Section 3.2
653 (determinants sought in flow typology, Fig.6a and 7a) and Section 3.3 (determinants sought in the
654 values of dimensionless numbers attached to the flow).

655

Case	Context	Authors	Model refinement	Spatiotemporal scales					Flow typology ²	Dimensionless numbers ³					
				L (m)	T (s)	H (m)	L/T (m s ⁻¹)	H/L ¹ (-)		T*	Re	Fr	S (%)	Λ_z	θ
A	Film flow	Charpin & Myers (2005)	NS	0.075	3.75	0.003	0.02	0.04	O	5	300	0.11	10	8.0	-
B	Laminar dynamics	Charru et al. (2004)	RANS	0.2	$1.8 \cdot 10^3$	0.007	$1.1 \cdot 10^{-3}$	0.035	O	6428	50	0.02	<0.01	12.1	0.14
C	Pool-riffles	Rathburn & Wohl (2003)	SV	70	$2.6 \cdot 10^6$	0.47	$3.5 \cdot 10^{-3}$	$6.7 \cdot 10^{-3}$	B	$7.8 \cdot 10^7$	$7.1 \cdot 10^7$	0.69	1.1	5108	34.1
D	Amazon River	Trigg et al. (2009)	ASV	$4.3 \cdot 10^3$	$3.15 \cdot 10^8$	10	$1.4 \cdot 10^{-3}$	$2.3 \cdot 10^{-3}$	F	58.5	$8 \cdot 10^7$	0.05	<0.01	6600	-

E	Step-pools	Grant et al. (1990)	SV	5530	1755	0.87	3.15	$1.5 \cdot 10^{-4}$	Hg	1.0	$2.7 \cdot 10^6$	1.03	4.5	1.25	-
F	Step-pools	Chin (1999)	SV	197.25	30	0.50	6.58	0.025	Hg	1.21	$4.0 \cdot 10^7$	3.58	6.25	1.22	-

† See section 3.1.2 - H/L is the fineness ratio of the flow comparing flow depth (H) to the length of the flow domain (L)

‡ See Section 3.2 - O: Overland, Hg: High-gradient, B: Bedforms, F: Fluvial

§ See Section 3.3 - T*: dimensionless period, Re: Reynolds number, Fr: Froude number, S: slope, Λ_z inundation ratio, θ Shields number

Table 1 - Six textbook cases representing an approximate envelope of all the tested cases in the L-T plane of Fig.2a, where L is the spatial scale (length of the flow domain) and T the temporal scale (duration of the process studied). Spatiotemporal scales are the determinants of modelling choices discussed in Section 3.1. The additional influence of flow typology and dimensionless numbers are discussed in Sections 3.2 and 3.3.

3.1.2 Influence of domain length (L) and flow depth (H)

The NS, RANS, SV and ASV equations are now positioned with respect to the spatial scale (L) and flow depth (H) of the reported experiments (Fig.-3), showing patterns and trends very similar to those of the (L, T) plane, though less pronounced. The global trend stays a decrease in refinement of the flow models from the smallest to the largest (L, H) values and typical scales of application may again be identified for each model refinement, NS ($10 \text{ cm} < L < 100 \text{ m}$, $1 \text{ mm} < H < 30 \text{ cm}$), RANS ($1 \text{ m} < L < 100 \text{ m}$, $5 \text{ cm} < H < 50 \text{ cm}$), SV ($10 \text{ m} < L < 20 \text{ km}$, $1 \text{ cm} < H < 2 \text{ m}$) and ASV ($10 \text{ m} < L < 1000 \text{ km}$, $10 \text{ cm} < H < 10 \text{ m}$). Some studies provide outliers for example Gejadze & Copeland (2006) for canal control purposes (NS, $L \sim 3 \text{ km}$, $H \sim 10 \text{ m}$) or Cassan et al. (2012) for flows in lined channels (RANS, $L \sim 50 \text{ cm}$, $H \sim 75 \text{ cm}$). In an overview, wider overlaps and more dispersion occur in the (L, H) than in the (L, T) plane, especially for low to medium scales: flow depth (H) seems less discriminating than the time scale (T) in the choice of a flow model.

The transverse analysis of H/L "fineness ratios" (dotted diagonals $H/L = 10^{-1}$, 10^{-2} , 10^{-3} , 10^{-4} and 10^{-5}) provides additional information, or rather a complementary reading grid on the information already plotted. First, only the NS and RANS models allow 2D (x, z) flow descriptions, which explains why these models have many of the largest H/L ratios (which, in most cases, stay within the $H < L$ shallow water hypothesis). Second, low H/L ratios provide justifications to discard 2D (x, z) descriptions at the benefit of 1D (x) descriptions within but also without the NS and RANS formalisms, so that the

second diagonal of Fig.-3 (roughly from the upper right to the lower left) also shows a decrease in model refinement, towards SV and ASV points.

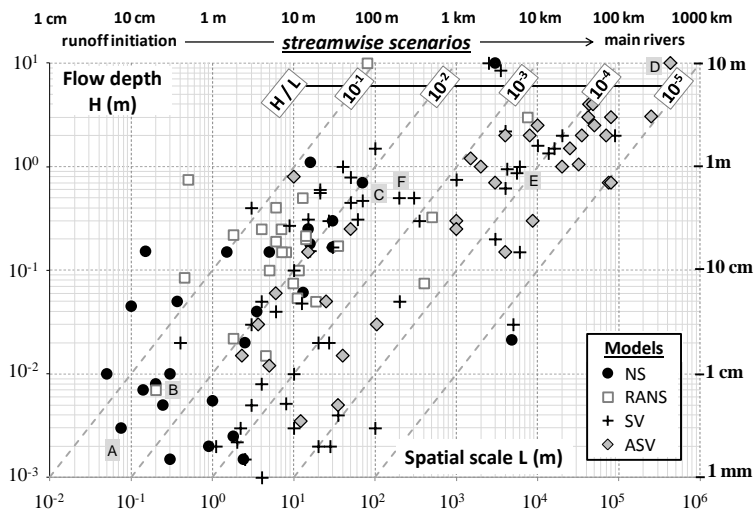


Figure 3 – How increasing (L, H) spatiotemporal scales of the flow domain tend to be associated with decreasing complexity in the choice of flow models, sorted here into four levels of refinement: Navier-Stokes (NS), Reynolds-Averaged Navier-Stokes (RANS), Saint-Venant (SV) and Approximations to Saint-Venant (ASV). A transverse analysis involves forming H/L ratios, searching for clues to model selection according to the "fineness" of the flow or governed by flow typologies that would exhibit specific H/L ratios. This figure was assembled from information available in the studies cited in Appendix A, selecting six textbook cases (sketches A to F, Table 1) for illustration.

3.1.3 Influence of domain length (L), time scale (T) and flow depth (H)

The links between model refinements (NS, RANS, SV or ASV) and spatiotemporal scales (L, T, H) were shown in the (L, T) and (L, H) planes (Fig.-2a and 3). There was first the expected correlation between increasing scales and decreasing model refinements. Then the transverse analyses involved re-examining the same dataset from the values of the L/T and H/L ratios, also seeking the determinants of modelling choices in the "system evolution velocity" (L/T) and "fineness" of the flow (H/L).

- The values of the L/T ratios indicate that modelling choices owe much to the long-term (low L/T) or short-term (high L/T) objectives associated with the target variables (velocity, discharge, particle transport, bed modifications) thus influencing the choice of T values. However, this choice is not totally free: it is likely constrained by flow characteristics and typologies.

- The values of the H/L ratios also indicate that flow typology (here, only its "fineness" is explicit) may be a mattering determinant for the choice of a modelling strategy. This idea is explored in far more details hereafter. The next section outlines the influence of friction, flow retardation and energy dissipation processes on flow typology. It advocates thus the definition of flow typologies from quantities related to the different types and/or magnitudes of flow retardation processes, provided these quantities are easily accessible (e.g. bed geometry, water depth, bed slope, size of the roughness elements).

3.2 Flow typology

3.2.1 From friction laws and bed topography to flow characteristics

Early insights on fluid friction and the definition of shear stress proportional to local velocity gradients came together with the action-reaction law (Newton 1687): friction exerted on the flow was of equal magnitude as the erosive drag, originally termed "critical tractive force" (Du Buat 1779) and held responsible for particle detachment. The friction laws mostly resorted to in present-day modelling do not often involve adaptations or generalisations of their famous empirical predecessors in civil engineering (Chézy 1775, Weisbach 1845, Darcy 1857, Manning 1871) even if practitioners and modellers are now confronted to far less controlled bed topographies and flow conditions, thus to a wider variety of flow typologies. The theoretical derivation (or justification) of contextually relevant friction laws seems therefore crucial, for water flow modelling at the microscopic (Richardson 1973, Jansons 1988, Priezjev & Troian 2006) or macroscopic scales (Smith et al. 2007, Powell 2014), and even more for **morphodynamic erosion** issues. In the literature, the modelling choices to account for

friction phenomena are most often correlated with the refinement of the flow models used (NS, RANS, SV, ASV) but also constrained by bed topographies and flow typologies in numerous cases.

Several studies at the NS level of refinement advocate the use of the "partial slip" (Navier 1827) condition or parented formulations in which the near-bed slip velocity is either proportional to the shear stress (Jäger & Mikelic 2001, Basson & Gerard-Varet 2008) or depends on it in a non-linear way (Achdou et al. 1998, Jäger & Mikelic 2003). Other works plead for "no-slip" conditions (Panton 1984, Casado & Diaz 2003, Myers 2003, Bucur et al. 2008, 2010) or suggest the separation of flow domains within or outside bed asperities, with a complete slip condition (non-zero tangential velocity) at the interface (Gerard-Varet & Masmoudi 2010). A wider consensus exists at the RANS level, calculating bottom friction as the local grain-scale values of the "Reynolds stresses" (Kline et al. 1967, Nezu & Nekagawa 1993, Keshavarzy & Ball 1997), which has proven especially relevant for flows in small streams over large asperities (Lawless & Robert 2001, Nikora et al. 2001, Pokrajac et al. 2007, Schmeeckle et al. 2007). However, he who can do more, can do less, and it is still possible to use the simplest empirical friction coefficients (Chézy, Manning) within sophisticated flow descriptions (NS: Lane et al. 1994, RANS: Métivier & Meunier 2003). In the literature, the SV level of refinement is a tilting point in complexity, that allows fundamental research, deriving *ad hoc* shear stress formulae from the local fluid-solid interactions (Gerbeau & Perthame 2001, Roche 2006, Devauchelle et al. 2007, Marche 2007) or applied research, adjusting parameter values in existing expressions, for specific contexts (*e.g.* boulder streams: Bathurst 1985, 2006, step-pool sequences: Zimmermann & Church 2001, irrigation channels: Hauke 2002, gravel-bed channels: Ferro 2003). This trend holds for most studies at the ASV level of refinement, though theoretical justifications of Manning's empirical formula were recently derived (Gioia & Bombardelli 2002) and a recent mathematical study of the diffusive wave equation (Alonso et al. 2008) introduces generalized friction laws for flows over non-negligible topographic obstacles. The event-based variability of the friction coefficient in ASV models has been investigated by Gaur & Mathur (2002).

If not decided from the level of refinement of the flow model, the friction coefficient (f) is chosen in accordance with flow typology and bed topography, the former often described by the Reynolds number (Re), the latter by the inundation ratio ($\Lambda_z = H/\varepsilon$ where ε is the size of bed asperities, to which flow depth H is compared). Such arguments were already present in the works of Keulegan (1938) and Moody (1944) on flow retardation in open-channel and pipe flows, relating values of the friction coefficient to the relative roughness ($\varepsilon/H = 1/\Lambda_z$) of the flow, across several flow regimes (laminar, transitional, turbulent) but only for small relative roughness (high inundation ratios). The existence of implicit relations between f , Re and Λ_z has somehow triggered the search for contextual alternatives to the sole f - Re relation for turbulent flows. Progressively lower inundation ratios were investigated (Smith et al. 2007) until the real cases of emergent obstacles received attention (Bayazit 1976, Abrahams & Parsons 1994, Bathurst 2006, Meile 2007, Mügler et al. 2010) including for non-submerged vegetation (Prosser et al. 1995, Nepf 1999, Järvelä 2005, Nikora et al. 2008). For site-specific friction laws, the default f - Re relation is sometimes complemented by f - Fr trends (Grant 1997, Gimenez et al. 2004, Tatard et al. 2008) or f - Λ_z relations (Peyras et al. 1992, Chin 1999, Chartrand & Whiting 2000, Church & Zimmermann 2007) in steep bed morphologies, where Fr is the Froude number (Froude 1868).

Knowledge gained on flow retardation processes lead to the identification of key dimensionless groups, to be included in any comprehensive analysis, formed from the "obvious", available elements of bed geometry previously mentioned (Julien & Simons 1985, Lawrence 2000, Ferro 2003, Yager et al. 2007). In numerous practical cases though, explicit bed geometries cannot be handled by the flow models. A crucial surrogate becomes then to include as many geometrical effects as possible in the chosen friction laws, for example these obtained from composite roughness experiments (Schlichting 1936, Colebrook & White 1937, Einstein & Banks 1950). A crucial advance was due to Smith & McLean (1977) who attributed distinct retardation effects to bed particles, particle aggregates and bedforms, corresponding to "grain spill", "obstructions" and "long-wave form resistance" in the subsequent literature. From then on, friction forces exerted by multiple roughness elements or scales

have often been described as additive-by-default, in shallow overland flows (Rauws 1980, Abrahams et al. 1986), gravel-bed streams (Bathurst 1985, Lawless & Robert 2001, Ferro 2003), natural step-pool formations (Chin & Wohl 2005, Canovaro & Solari 2007, Church & Zimmermann 2007) and man-made spillways or weirs (Peyras et al. 1992, Chinnarasri & Wongwise 2006).

3.2.2 From flow characteristics to flow typologies

Several authors have put forward the existence of a scale-independent link between bed geometry, flow retardation and flow structure, through the existence of three distinct flow regimes, from geometrical arguments: "isolated roughness", "wake interference" and "skimming" flow (Morris 1955, 1959, Leopold et al. 1960, Fig. 4a, c and e). These flow descriptions were later applied in very different contexts (Abrahams & Parsons 1994, Chanson 1994a, Papanicolaou et al. 2001, Zimmermann & Church 2001), which suggests that analogies in energy dissipation and flow retardation may exist across scales, from similar geometries and flow characteristics. This makes the description somewhat generic, possibly used to constitute a set of flow typologies.

In Fig. 4a, the isolated roughness flow is laminar or weakly turbulent and the shade (streamline diversion) of an obstacle does not reach the next. This setting ensures maximum energy dissipation, which also holds for stepped cascades of natural or man-made nature in Fig. 4b: "nappe flows" loose strength through energy-consuming fully-developed hydraulic jumps, isolated behind the major obstacles (Peyras et al. 1992, Chanson 1994b, Wu & Rajaratnam 1996, 1998). In Fig. 4c the wake-interference flow is transitional or turbulent. The drag reduction and partial sheltering between obstacles depend on their spatial distribution and arrangements, as in Fig. 4d that shows "partial nappe flow" in relatively flat step-pool formations, with incomplete hydraulic jumps between obstacles of irregular sizes and spacing (Wu & Rajaratnam 1996, 1998, Chanson 2001). In Fig. 4e, the turbulent skimming flow exhibits a coherent stream cushioned by the recirculating fluid trapped between obstacles and responsible for friction losses. Similar characteristics appear in Fig. 4f, for submerged

cascades or large discharges on stepped spillways. Air entrapment begins where the boundary layer reaches the free surface and flow aeration triggers subscale energy dissipation (Rajaratnam 1990, Chanson 1994b).

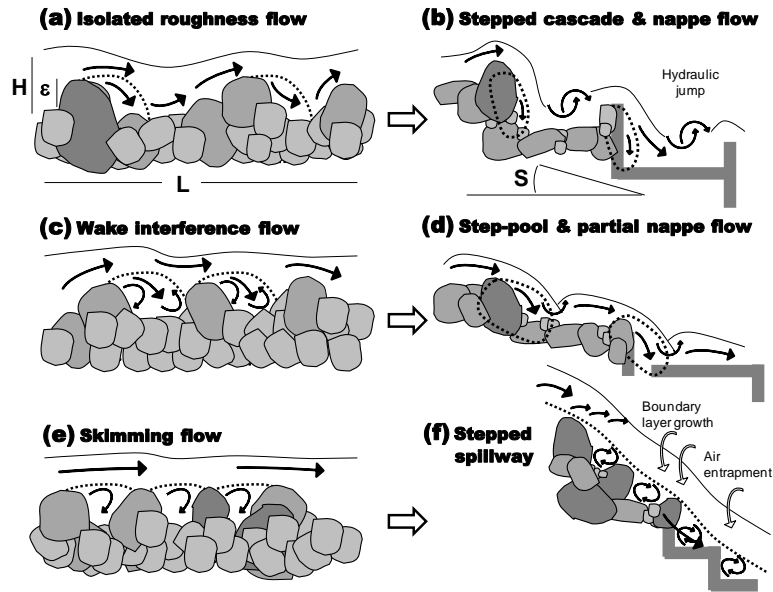


Figure 4 – Analogies in flow characteristics, retardation processes and energy dissipation structures for very different flow typologies: streams (a, c, e) and high-gradient natural or man-made stepped flows (b, d, f). The combined values of flow depth (H), slope (S) and inundation ratio ($\Lambda_z = H/\epsilon$, where ϵ is the roughness size) appear as strong geometrical controls over flow characteristics and typologies. The very small inundation ratios ($\Lambda_z < 1$) typical of overland flows in hydrology (flows through emergent obstacles, including vegetation) correspond to ϵ values larger than H values (tortuous flows are best seen in the top views of Fig.8).

At this point, our set of flow typologies should be obtained from the geometrical arguments available in Fig.-4 (water depth H, bed slope S, inundation ratio $\Lambda_z = H/\epsilon$). The simplest way to proceed is to work in the (S, H) plane, then to indicate the values of Λ_z if the values of S and H are not discriminating enough for each pair of (S, H) values. The first two flow typologies (Overland flow, noted O, and High-gradient flow, noted Hg) may be identified by a single criterion on H only

($H < H_{LIM}$, Emmett 1970, Wainwright et al. 2008) or on S only ($S > S_{LIM}$, Grant et al. 1990, Rosgen 1994, Montgomery & Buffington 1997). At least two flow typologies remained to be distinguished, Fluvial flows (F) and flows over significant bedforms (*e.g.* rough plane bed, dune-ripples or pool riffles, as suggested by Montgomery & Buffington 1997), referred to as Bedforms (B) in the following. Though Fluvial flows are expected to have the highest flow depths, an additional criterion on Λ_z may be used to make the difference between these last two typologies. Figure 5 positions the selected (O, Hg, B, F) flow typologies in the (S , H) plane.

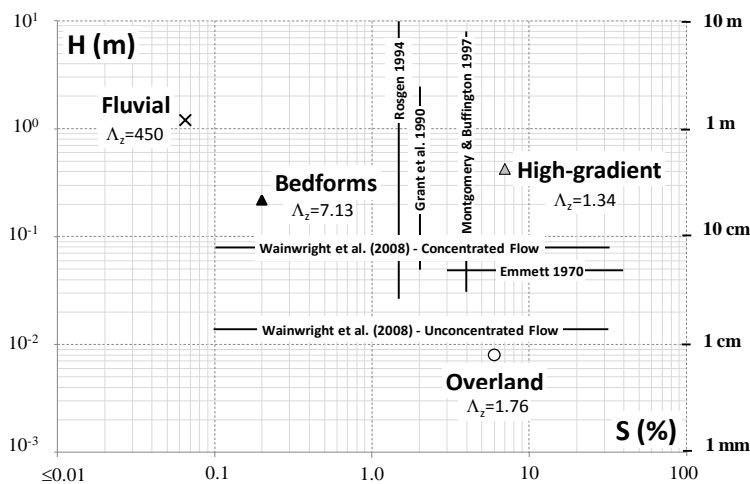


Figure 5 – Median position of the studies belonging to the "Overland", "High-gradient", "Bedforms" and "Fluvial" flow typologies, plotted on the (S : slope, H : water depth) plane, with indication of the associated also tracing an approximate additional criterion on the inundation ratio ($\Lambda_z = H/\epsilon$), where ϵ is the size of the bed asperities) to separate the Fluvial and Bedforms types of flow. This figure was assembled from information available in the studies cited in Appendix A.

Moreover, there is a strong link between Fig.-4 and 5, which tends to ensure the genericity (if not uniqueness) of the selected set of typologies. The Overland typology corresponds to Fig.-4a or c, the Bedforms typology likely appears in Fig.-4c, the Fluvial typology in Fig.-4 and the High-gradient typology in Fig.-4b, d or f. In coherence with Fig.-5, an increase in bed slope changes the Bedforms

and Fluvial typologies into the High-gradient typology, while an increase in both water depth and bed slope is needed to do the same from the Overland typology.

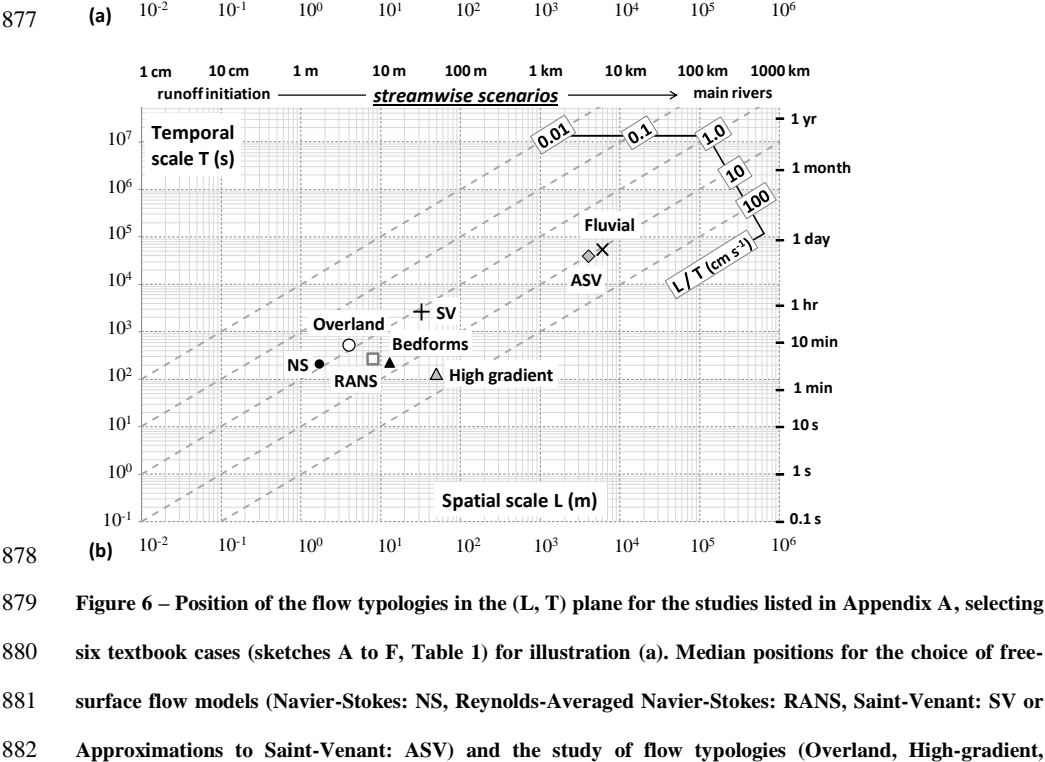
3.2.3 Influence of flow typologies on modelling choices

Figures 6 and 7 provide a comprehensive picture of the most used associations between models (NS, RANS, SV or ASV), scales (L, T, H) and flow typologies (O, Hg, B or F) just added to the analysis. These figures seem to indicate preferential [NS, O], [RANS, B] and [SV, Hg] associations, in addition to the obvious [ASV, F] pair. The (L, H) plot of Fig.-7b6 seems more discriminating than the (L, T) plot of Fig.-6b7 though similaridentical trends appear.

The [NS, O] association arises from the fact that several Overland studies involve very shallow laminar flows and low sediment transport rates, best handled by adapted formulations of the NS equations (nearly at the SV level), made suitable for low "system evolution velocities" ($L/T \approx 0.01 \text{ m s}^{-1}$, Fig.-6). At somewhat larger spatial scales, the widely-used and multipurpose SV model has rather low median $L/T \approx 0.02 \text{ m s}^{-1}$ values, mainly because many of its applications concern laminar flow modelling and granular transport, as an alternative to the NS system or in formulations at complexity levels intermediate between the NS and SV descriptions. These are clues that the [SV, O] association may also be of special interest, despite the closest median positions of the NS and O points in the (L, T) and (L, H) plots.

The RANS model (median $L/T \approx 0.07 \text{ m s}^{-1}$) and the ASV models (median $L/T \approx 0.1 \text{ m s}^{-1}$) tend to involve higher "system evolution velocities". The former typically targets the description of numerous short-term, high-frequency events (quadrant analysis for fluctuations in near-bed velocity, particle pick-up by turbulent bursts). The latter is often associated with Fluvial flows: low H/L ratios with high enough H and Λ_z values with weak friction, often resulting in very turbulent, high-velocity flow. Moreover, studies handling morphodynamicerosion issues within the ASV formalism often

873 hypothesize particle transport to occur as suspended load only, equating particle and flow velocities,
 874 thus typically not extending the time scale of the study to address the long-term, low velocity bedload
 875 transport involved in morphodynamics, for example.
 876



Bedforms or Fluvial) across scales in the (L, T) plane (b). A transverse analysis involves forming L/T ratios, searching for clues to model selection according to these "system evolution velocities" or governed by flow typologies that would exhibit specific L/T ratios.

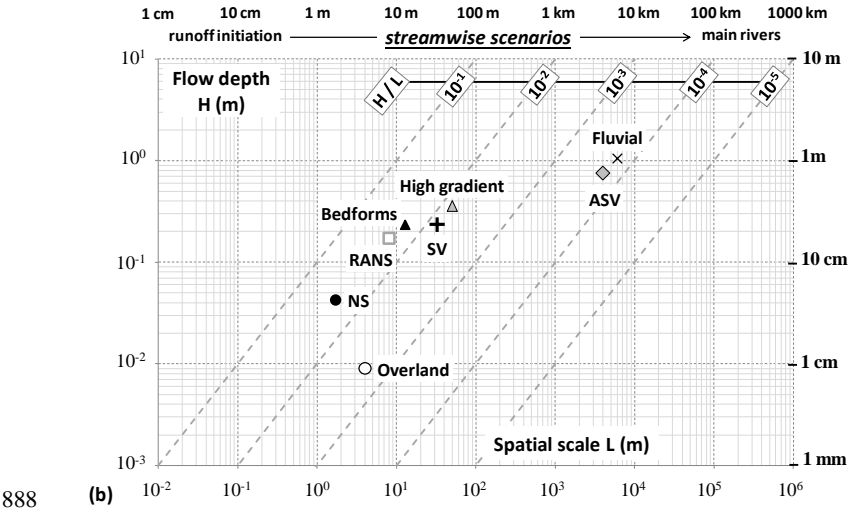
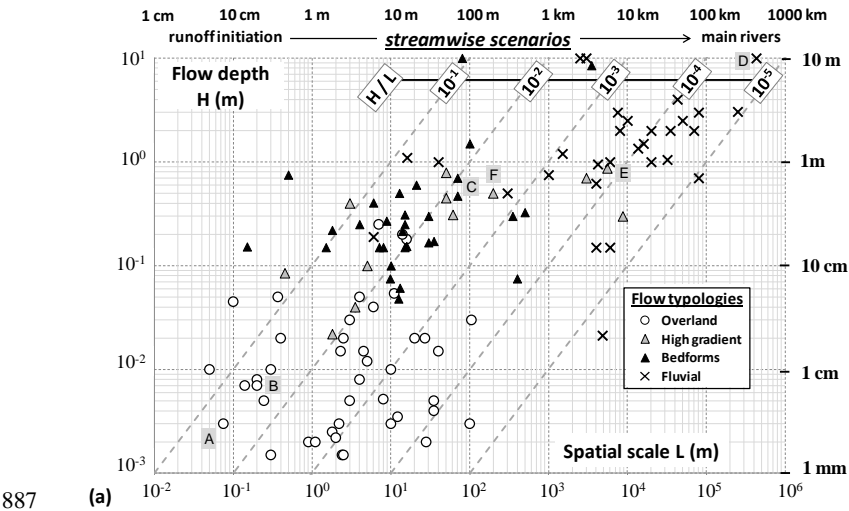


Figure 7 – Position of the flow typologies in the (L, H) plane for the studies listed in Appendix A, selecting six textbook cases (sketches A to F, Table 1) for illustration (a). Median positions for the choice of free-surface flow models (Navier-Stokes: NS, Reynolds-Averaged Navier-Stokes: RANS, Saint-Venant: SV or

893 **Approximations to Saint-Venant: ASV) and the study of flow typologies (Overland, High-gradient,**
894 **Bedforms or Fluvial) across scales in the (L, H) plane (b). A transverse analysis involves forming H/L**
895 **ratios, searching for clues to model selection according to these "finenesses" of the flow domain or**
896 **governed by flow typologies that would exhibit specific H/L ratios.**

897

898

Several principles of organization between flow typologies may be inferred from reference studies (Grant et al. 1990, Montgomery & Buffington 1997, Church 2002) that discuss their succession in space (along longitudinal profiles) but also in time (which flow typologies are "experienced" by the flowing water during its course and which are the associated time scales). Plausible "streamwise scenarios" may therefore be assembled (Fig.-8), routing flow aggregations across increasing spatiotemporal scales and through several flow typologies, from the narrow-scale upland flows (runoff initiation) to the regional scales of the main rivers.

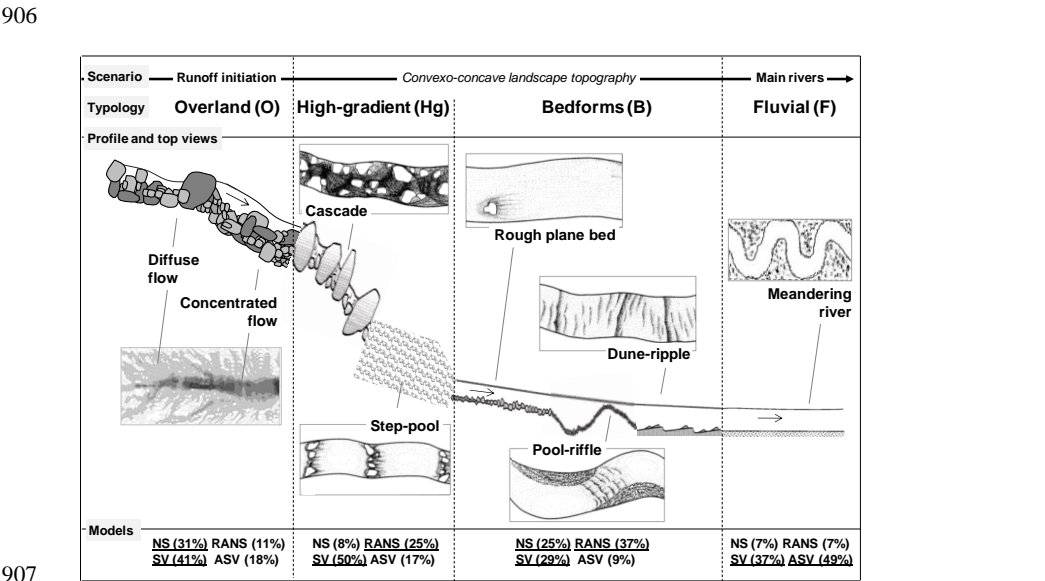


Figure 8 – Streamwise scenario for a convexo-concave landscape topography, from runoff initiation to the main rivers, across flow typologies (Overland O, High-gradient Hg, Bedforms B or Fluvial F) and spatiotemporal scales (L, T, H). The indicated L, T and H values are the median values for the spatial scale, time scale and water depth, respectively, from the literature cited in Appendix A (Fig. 6 and 7). All sketches and drawings for the High-gradient and Bedforms typologies were taken from Montgomery & Buffington (1997). The top view for Overland flow is from Tatard et al. (2008) and that of a meandering river from Rosgen (1994). The "Models use" panel indicates the model refinements most used (Navier-Stokes NS, Reynolds-Averaged Navier-Stokes RANS, Saint-Venant SV or Approximations to Saint-Venant ASV) to describe a given flow typology in the literature cited in Appendix A.

3.3 Dimensionless numbers

3.3.1 Contextual dimensionless numbers

Complementary indications on ~~An angle of attack for the establishment of~~ modelling strategies ~~are~~ is provided by dimensional analysis, to delineate the domains of validity of the selected flow models (NS, RANS, SV or ASV), across their multiple spatiotemporal scales of application but in a powerful scale-independent analysis. Justifications for the use of dimensionless numbers may be sought in the developments of similitude laws (Fourier 1822, Rayleigh 1877, Bertrand 1878, Vaschy 1892, Riabouchinsky 1911), later extended to dimensional analysis, providing guidance for the sizing of experimental facilities used in reduced-scale modelling as well as more general arguments for the choice of adequate sets of dimensionless quantities (Buckingham's 1914 π -theorem, Bridgman 1922, Langhaar 1951, Bridgman 1963, Barenblatt 1987). Throughout history, the establishment of dimensionless numbers has led to the recognition of contextually dominant terms in the flow equations, rendering them prone to dedicated simplifications, provided these would not be used outside their conditions of validity, following successive hypotheses made during their derivation. On the one hand the dimensionless numbers arise in the non-dimensionalisation of the systems of governing equations, being an inherent feature of the model. On the other hand only the selected dimensionless numbers appear in the non-dimensional formulation of the equations, from appropriate arrangements of their terms, and this choice indicates which are the physical processes of interest for the modeller. Finally, not all dimensionless numbers can be made explicit in the simplest mathematical models (especially the ASV models) but their values can always be calculated, thus correlated (or not) with the use of one or the other of the flow models.

From a wide overview of free-surface flow and ~~morphodynamic~~~~erosion~~ studies, a few dimensionless numbers stood out and will be used in the procedure presented in the following. Some have already been mentioned (Reynolds number Re , Froude number Fr) and some others have even been used to define flow typologies (bed slope S , inundation ratio Λ_z). As all dimensionless numbers

aim to describe flow typology, the introduction of two more dimensionless numbers may be seen as an attempt to re-examine the influence of flow typologies on modelling choices, from a different, more complete perspective (especially if the dimensionless numbers not used in the definition of flow typologies prove discriminating for the modelling choices).

- The dimensionless period $T^*=T/T_0$ handles temporal aspects by comparing the chosen time scale (T) to the natural time scale (T_0) of the system, the latter obtained from the spatial scale of the system and the average flow velocity as $T_0=L/U$ (Fig. 1). This dimensionless group or equivalent formulations are used to model wave celerity in flood propagation issues (Ponce & Simons 1977, Moussa & Bocquillon 1996a, Julien 2010) or to quantify the long characteristic times ($T^*\gg 1$) of basin-scale sedimentation. In the latter, particle transport (and significant bed modifications) typically involve lower velocities (and larger time scales) than these of water flow (Lyn 1987, Paola et al. 1992, Howard 1994, Van Heijst et al. 2001) and the chosen T value witnesses this discrepancy.

- The Reynolds number $Re=UH/\nu$ compares flow inertia (velocity U times depth H) with the adverse action of (kinematic) viscosity (ν [$L T^{-2}$]). In natural setting, over very rough boundaries, fully turbulent flows are often reported for $Re>2000$, while the onset of turbulence within transitional regimes occurs at $Re\sim 500$. Laminar overland flows, especially thin film flows, may have Re values as low as $Re<100$.

- The Froude number $Fr=U/(gH)^{0.5}$ denotes the influence of gravity (g) on fluid motion. Supercritical $Fr>1$ values indicate torrential flows, for example flows accelerated by pressure effects, in which waves propagate only downstream, also compatible with the appearance of localised energy dissipation patterns (white waters, hydraulic jumps). Subcritical $Fr<1$ values indicate tranquil flows with downstream controls. However, the presence of a movable bed makes the identification of sub- and supercritical regimes less obvious, as additional phenomena come into play (Lyn 1987, Lyn & Altinakar 2002).

- Topographical effects on flow phenomenology are almost always explicitly accounted for through the average bed slope S, typically ranging from nearly zero ($S<0.01\%$) for large rivers to extremely high values ($S\approx 100\%$) for gabion weirs, chutes or very steep cascades.

- Topography also appears through the inundation ratio $\Lambda_z = H/\varepsilon$ which allows a direct, model-independent analysis of friction phenomena (Lawrence 1997, 2000, Ferguson 2007, Smith et al. 2007) possibly dealing with large-size obstacles and form-induced stresses (Kramer & Papanicolaou 2005, Manes et al. 2007, Cooper et al. 2013). The encountered values of Λ_z are very high for rivers flowing on smooth, cohesive, fine-grained beds ($\Lambda_z > 100$) and very low for all types of flows between emergent obstacles ($\Lambda_z < 1$, Ferro 2003, Hogarth et al. 2005, Canovaro & Solari 2007, Ferguson 2007, Lamb et al. 2008) including flow through vegetation (see Järvelä 2004, Holden et al. 2008, Gumiere et al. 2011a, Kim et al. 2012, Nepf 2012).

- The dimensionless Shields number $\theta = \tau_0 / g \varepsilon_p (\rho_p - \rho)$ compares the drag force exerted on bed particles to their immersed weight, where ε_p and ρ_p account for the size and density of erodible particles.. The ratio between the current θ and the critical θ_c values indicates local flow conditions of deposition ($\theta < \theta_c$), incipient motion ($\theta \approx \theta_c$), transportation as bedload ($\theta > \theta_c$) or into suspension ($\theta \gg \theta_c$) (Shields 1936). This number seems appropriate for most morphodynamic erosion issues because it has been widely applied and debated in the literature (Coleman 1967, Ikeda 1982, Wiberg & Smith 1987, Zanke 2003, Lamb et al. 2008) and also because of its numerous possible adaptations (Neill 1968, Ouriemi et al. 2007, Miedema 2010) to various flow typologies and non-uniform or poorly-known bed conditions. An impressive review on the use of the Shields number to determine incipient motion conditions, over eight decades of experimental studies, may be found in Buffington & Montgomery (1997). Finally, Fig.9 provides a generalized Shields diagram that includes motion threshold criteria under the effects of high or low particle exposure (Miedema 2010) or for laminar flows, also indicating the conditions of significant suspension (Wright & Parker 2004). To search for additional indications, the points in Fig.9 have been sorted by flow depths with the arbitrary $H=5$ cm threshold. Other case classifications may be relevant, for example identify the hydrological and hydraulic contexts.

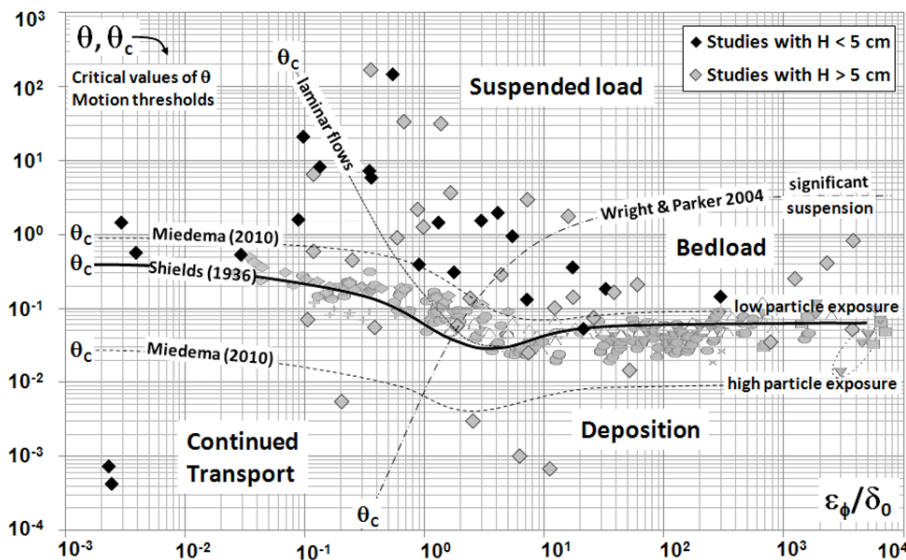


Figure 9 - Generalized dimensionless Shields diagram that summarizes the conditions and regimes of sediment transport or deposition, from the relative values of the Shields parameter (θ) and incipient motion criterion (θ_c). The X-axis bears the values of the ratio of particle size (ϵ_ϕ) on the depth of the laminar sublayer (δ_0). The diamonds refer to the studies cited in Appendix A that deal with morphodynamic issues: black diamonds for studies in which flow depth is $H < 5$ cm, grey diamonds otherwise. Data in the background show the critical θ_c values reported in the wide Buffington & Montgomery (1993) review of incipient motion conditions for varied flow regimes, particle forms and exposures.

3.3.2 Influence of the dimensionless numbers

As the purpose here is to re-examine the influence of flow typologies from the angle-point of view of the dimensionless numbers, the chosen representation (Fig.-10) discards the (L, T, H) spatiotemporal scales. It first recalls the preferential associations between models and flow typologies (see the "Models-use" panel of Fig.-8) by tracing connecting dotted lines between flow typologies and the models most used to handle them, in the legend of Fig.-10. It then examines whether these

1011 associations still hold, for each of the six dimensionless numbers, by plotting and comparing the
1012 median values of T^* , Re , Fr , S , Λ_z and θ for model uses (NS, RANS, SV or ASV) and flow typologies
1013 (O, Hg, B, F). The dotted ellipses are "confirmations" (*e.g.* no additional information may likely be
1014 obtained from Re , Fr and θ). Conversely, the presence of "non-associated" points (P_1 for T^* , P_2 and P_3
1015 for S , ~~P_4 for Λ_z~~) signals ~~something new: an influence not yet accounted for.~~ cases in which the
1016 determinants of modelling strategies should be thought altogether in spatiotemporal scales, flow
1017 typologies and the values of certain dimensionless numbers.

1018
1019 For example, the isolated P_1 point indicates the expected [ASV, -F] association does not appear on
1020 the T^* values, as the ASV applications exhibit higher median T^* values than the F typologies. The
1021 suggested interpretation is that large (L, T, H) scales and Fluvial flows likely trigger the use of the
1022 ASV model, though the necessity to handle large dimensionless periods makes the typological
1023 argument less conclusive. The P_2 and P_3 points also -indicate the break of the ~~NS-O and~~ [ASV, -F]
1024 associations when examined from the angle of the bed slopes. This reinforces the use of bed slopes in
1025 the search for determinants of modelling choices, either in the definition of flow typologies in the (S,
1026 H) plane or as such. ~~The P_4 point indicates the break of the NS-O association when considering the~~
1027 ~~values of the inundation ratio, with the same conclusion as above.~~

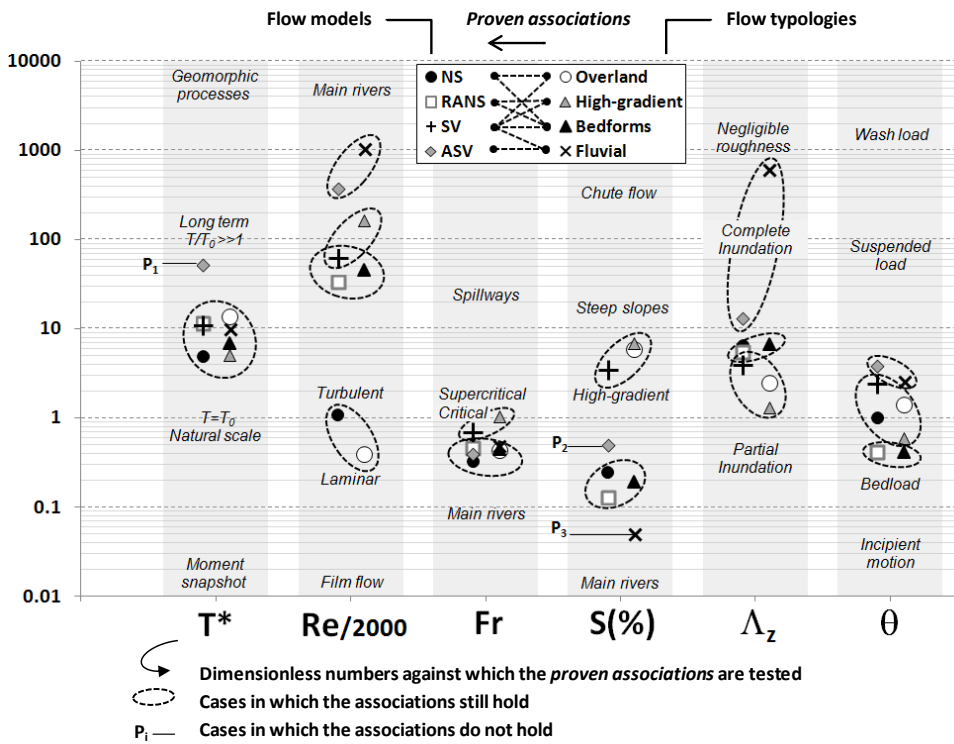


Figure 10 - Comparative overview of the median values of the six selected dimensionless numbers (dimensionless period $T^*=T/T_0$, ratio of the chosen time scale on the "natural" time scale of the flow, Reynolds number Re , Froude number Fr , slope S , inundation ratio Λ_z and Shields parameter θ) obtained for the use of systems of equations (NS, RANS, SV and ASV) and the description of flow typologies (O, Hg, B and F) in the cited literature. The expected associations are indicated by dotted connecting lines in the legend box. The confirmed associations are indicated by dotted ellipses. Broken associations (isolated points P_1 to P_4) are discussed in the text. The typical and extreme ranges of the mentioned dimensionless numbers have been added for indication. This figure was assembled from information available in the studies cited in Appendix A.

4 Conclusion

4.1 Outcomes of this review

In a free opinion on the use of models in hydrology, De Marsily (1994) elegantly argued that the modelling of observable phenomena should obey “*serious working constraints, well-known from classical tragedy: unity of place, unity of time, unity of action*”. This review paper investigates how known spatial scales, temporal scales and flow typologies constrain the choice of a modelling strategy. A normative procedure was built to facilitate the search for determinants of the modelling choices in the cited literature.

- Each free surface flow model was placed in one of the NS, RANS, SV or ASV categories, whose decreasing levels of refinement account for "Navier-Stokes", "Reynolds-Averaged Navier-Stokes", "Saint-Venant" or "Approximations to Saint-Venant" types of approaches.

- The explored (L, T, H) spatiotemporal scales cover multiple orders of magnitude in the streamwise direction ($5\text{ cm} < L < 1000\text{ km}$), the time duration ($0.1\text{ s} < T < 1\text{ yr}$) and flow depth ($1\text{ mm} < H < 10\text{ m}$) while the modelling subscales (δL , δT) used for data collection and/or the size of the calculation grid are in the $0.01\text{ mm} < \delta L < 10\text{ km}$ and $0.001\text{ s} < \delta T < 1\text{ day}$ intervals.

- This study also encompasses a wide variety of free-surface flows, reduced to four typologies from arguments on bed geometry, friction, flow retardation and energy dissipation processes. These typologies are Overland flow (O: diffuse or concentrated), High-gradient flow (Hg: cascades, step-pools), flows over significant Bedforms (B: rough plane beds, dune ripples, pool riffles) and Fluvial flows (F: rivers, canals). Overland flows have the shallowest depths, High-gradient flows the highest bed slopes, Fluvial flows have high flow depths and negligible bed roughness while Bedforms flows may have any flow depth, over pronounced, non-negligible bedforms.

- In addition to the spatiotemporal scales and flow typologies, the determinants of modelling choices are also sought in a series of six popular dimensionless numbers: the dimensionless period (T^*), Reynolds and Froude numbers (Re, Fr), the bed slope (S), the inundation ratio ($\Lambda_z = H/\varepsilon$ where ε is the size of bed asperities) and the Shields number (θ) that compares drag forces to particle weight.

1067

1068 In summary, each case-study may be defined by its signature, comprised of the *chosen* model (NS,
1069 RANS, SV or ASV) and modelling subscales (δL , δT), ~~the~~ versus given spatiotemporal scales (L, T,
1070 H), flow typology (O, H, B or F) and dimensionless numbers (T^* , Re, Fr, S, Λ_z , θ). Though non-
1071 unique, this signature is a generic and normative classification of studies interested in free-surface
1072 flow modelling, with or without morphodynamic ~~erosion~~ issues.

1073 - The present review first illustrated the expected dominant trend of decreasing model refinement
1074 with increasing (L, T, H) spatiotemporal scales and (δL , δT) subscales. It appeared then that model
1075 uses could also be sorted by their L/T and H/L ratios, though less clearly, which nevertheless provided
1076 indications that the spatiotemporal scales were not the only determinant of modelling choices. This
1077 result suggested that flow typologies (reduced here to the L/T "system evolution velocity" and H/L
1078 "fineness of the flow") were also influential factors.

1079 - A more exhaustive set of flow typologies was then derived from simple geometrical arguments,
1080 combining criteria on S, H and Λ_z , represented in the (S, H) plane. This allowed quantifying the
1081 median scales associated with studies interested in the Overland (O), Bedforms (B), High-gradient
1082 (Hg) and Fluvial (F) typologies, sorted here by increasing spatiotemporal scales. Then came the
1083 identification of preferential associations between flow models, scales and typologies: [NS, O] or [SV,
1084 O], [NS, B], [RANS, B] or [SV, B], [RANS, Hg] or [SV, Hg] ~~or [ASV, Hg]~~ and [ASV, F]. ~~for~~
1085 ~~increasing spatiotemporal scales.~~

1086 - The final step was to re-examine the previous associations from the values of the dimensionless
1087 numbers, thought here as more detailed, scale-independent descriptors of flow typologies. Several
1088 associations were confirmed by the median values of the associated dimensionless numbers but ~~the~~ T^*
1089 (dimensionless period) and S (bed slope) ~~and Λ_z (inundation ratio)~~ introduced additional information.,
1090 *i.e.* correcting trends.

1091

1092 All arguments prevailing in the identification and sorting of flow models, scales, typologies and
1093 dimensionless numbers may easily be debated and adapted, within the hydro-morphodynamics ~~logy~~

erosion community or for other research purposes. For example, multiple flow models, scales, typologies and dimensionless numbers also intervene in the fields of pesticide fate modelling and groundwater contamination issues, so the same procedure could be applied. Finally, this procedure offers the possibility to enrich the database of signatures if each modeller records their conceptual choices (flow models) in the proposed reading grid, together with the contextual elements (scales, typologies, dimensionless numbers) handled, for present and past studies. This would first help forming a comprehensive database view of modelling choices, thus seeking guidance from "what has been done in similar cases", which however does not provide any critical analysis. Complementary investigations could certainly address the question of "what should be done", this time deciding the "model" part of the signatures from recommendations based on the scales, typologies and dimensionless numbers, as well as from additional elements, typically the modelling objectives (Fig. 11).

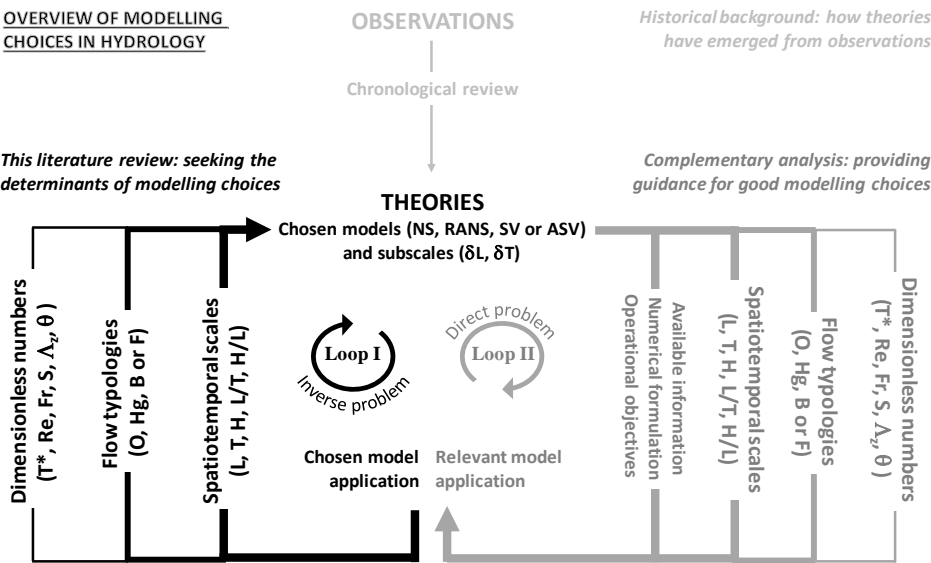


Figure 11 – This figure provides a simplified overview of the available modelling choices in hydrology, in three distinct colours associated with specific research purposes or disciplines, showing the position of the present review relative to the others. The pale grey section aims at understanding how the available flow models have emerged from

observations and early formulations of the flow equations, focusing on their conditions of validity *i.e.* the successive hypotheses made during their derivation. The black section recalls the procedure followed in this review paper (Loop I, "inverse problem"). Literature sources are processed through a procedure that analyses how the spatiotemporal scales (spatial scale L , time scale T , flow depth H , L/T and H/L ratios), then flow typology (Overland O, High-gradient Hg, Bedforms B or Fluvial F) and dimensionless numbers (dimensionless period T^* , Reynolds number Re , Froude number Fr , bed slope S , inundation ratio Λ_z , Shields parameter θ) determine the choice of a flow model (Navier-Stokes NS, Reynolds-Averaged Navier-Stokes RANS, Saint-Venant SV or approximations ASV) and that of data collection and/or modelling subscales (δL , δT). Suggested in medium grey on the right are the scope and principles of future research challenges that would address the "what should be done?" (Loop II, "direct problem") question in echo to the current "what has been done?" concern (Loop I).

4.2 Research challenges ~~in hydrology~~ and philosophy of modelling

This review has sought the determinants of modelling choices in hydrology (Fig. ~~ure~~-11, Loop I) from the basis provided by literature sources, without any intention to provide recommendations regarding appropriate (both relevant and cost-effective) modelling strategies. However, for most practical applications, the starting point is the definition of a scope and the endpoint is the evaluation of the objective function to evaluate the success or the failure of the chosen modelling strategy. A question thus arises on how to guide the modeller in the choice of an adequate model, in function of given, approximately known spatiotemporal scales, flow typology and dimensionless numbers (Fig. ~~ure~~-11, Loop II). According to the principle of parsimony, modellers should seek the simplest modelling strategy capable of (i) a realistic representation of the physical processes, (ii) matching the performances of more complex models and (iii) providing the right answers for the right reasons.

- (i) Throughout the last decades, an important change of the scope of free-surface flow modelling applications has taken place, with subsequent changes in the objective functions resorted to. The development of hydrological and hydraulic sciences has been directly linked to the progresses in understanding processes, in theoretical model development (*e.g.* computational facilities: numerical techniques, data assimilation, thorough model exploration, inverse calculus) and in data acquisition

1139 (new devices, remote sensing, LiDAR). *"It may seem strange to end a review of modelling with an*
1140 *observation that future progress is very strongly linked to the acquisition of new data and to new*
1141 *experimental work but that, in our opinion, is the state of the science"* (Hornberger & Boyer 1995).

1142 - (ii) However, there remains an important need for research on classical free-surface flow
1143 (hydrological or hydraulic) modelling for engineering applications in predicting floods, designing
1144 water supply infrastructures and for water resources management, from the headwater catchment to
1145 the regional scale. More recently, free-surface flow modelling has become an indispensable tool for
1146 many interdisciplinary projects, such as predicting pollution and/or erosion incidents, the impact of
1147 anthropogenic and climate change on environmental variables such as water, soil, biology, ecology, or
1148 socio-economy and ecosystemic services. The direct consequence is a significant increase of the
1149 complexity of the objective function, from simple mono-site (*e.g.* one-point), mono-variable (*e.g.* the
1150 water depth) and mono-criterion (*e.g.* the error on peakflow) to complex multi-site (*e.g.* large number
1151 of points within a catchment), multi-variable (*e.g.* water depth, hydrograph, water table,
1152 concentrations, ecological indicators, economic impact) and multi-criteria (*e.g.* errors on peakflow,
1153 volume, RMSE) objective functions.

1154 - (iii) There is often a mismatch between model types, site data and objective functions. First,
1155 models were developed independently from the specificities of the study site and available data, prior
1156 to the definition of any objective function. In using free-surface flow models, the context of their
1157 original purpose and development is often lost, so that they may be applied to situations beyond their
1158 validity or capabilities. Second, site data are often collected independently of the objectives of the
1159 study. Third, the objective function must be specific to the application but also meet standard practices
1160 in evaluating model performance, in order to compare modelling results between sites and to
1161 communicate the results to other scientists or stakeholders. The known danger is to use flow and
1162 | morphodynamic~~erosion~~ equations outside their domains of validity (*i.e.*, breaking the assumptions
1163 | made during their derivation) then to rely on the calibration of model parameters as for technical
1164 compensations of theoretical flaws, at the risk of losing the physical sense of model parameters,

1165 creating equifinality and obtaining the “*right results for the wrong reason*” (Klemeš 1986). Choosing
1166 the right model for the right reason is crucial but the identification of the optimal data-model couple to
1167 reach a predefined objective is not straightforward. We need a framework to seek the optimum balance
1168 between the model, data and the objective function as a solution for a hydrological or hydraulic
1169 problem, on the basis of the principle of parsimony. The latter follows a famous quote often attributed
1170 to Einstein, that “*everything should be made as simple as possible, but not simpler*” which somehow
1171 originates in the philosophy of William of Ockham (1317) (*Numquam ponenda est pluralitas sine*
1172 *necessitate [Plurality must never be posited without necessity]*) or may even be traced back to
1173 Aristotle's (~350 BCE) *Analytica Posteriora* that already advocated demonstrations relying on the
1174 fewest possible number of conjectures, *i.e.* the dominant determinisms.

1175 Finally, analytical procedures for free-surface flows and morphodynamic erosion issues necessitates
1176 a comprehensive analysis of the interplay between models (assumptions, accuracy, validity), data
1177 requirements and all contextual information available, encompassed in the “signature” of any given
1178 application: model refinement, spatiotemporal scales, flow typology and scale-independent description
1179 by dimensionless numbers. This review helps the modeller positioning his (or her) case study with
1180 respect to the modelling practices most encountered in the literature, without providing any
1181 recommendation. A complementary step and future research challenge is to decipher relevant
1182 modelling strategies from the available theoretical and practical material, resorting to the same objects,
1183 the previously defined signatures. Its purpose clearly is to address the “*which model, for which scales*
1184 *and objectives?*” question. A complete analytical framework, comprised of both loops, would provide
1185 references and guidelines for modelling strategies. Its normative structure in classifying theoretical
1186 knowledge (the mathematics world, equations and models) and contextual descriptions (real-life
1187 physical processes, scales and typologies) hopefully makes it also relevant for other Earth Sciences.

1188

1189

Appendix A. References used in the Figures.

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