

Montpellier, January 19th, 2016

Reply to Reviewers' comments on **Determinants of modelling choices for 1-D free-surface flow and erosion issues in hydrology: a review** by Bruno Cheviron and Roger Moussa.

This document details our responses to Reviewer#1 and #2. The line numbers noted **LXXX-YYY** refer to the change-tracking version of the revised manuscript. The new version of the manuscript has been uploaded separately.

Response to Reviewer #1

General Comment

Overall, I liked reading this critical review of existing models that can be used for the coupled modeling of water flow and sediment transport in hydrology. I'm not a big fan of review articles in general, but in this case – provided that review articles are allowed by the journal – I express appreciation for the work and limit myself to specific comments about the manuscript. The paper is clear and well written and the review of the methods presented (at least to my knowledge) is quite complete. The English usage is correct, and the presentation of good quality. I have only a few minor comments about the manuscript:

Response (R): We thank Reviewer #1 for his positive evaluation and encouragements. We totally agree with his comments and will introduce responses in the revised version as shown below.

- I see that while the "flow" parts are discussed based on specific equations, the "sediment transport" sections are a bit more qualitative, and there are no equations there. I'm wondering if this is a deliberate choice of the authors and if they could comment a bit on this choice, maybe even in the manuscript;

R: This was indeed a deliberate choice, searching for the determinants of modelling strategies in the refinement of the flow and erosion models, then in flow typologies, then in the dimensionless numbers used.

Regarding erosion, the default/starting hypothesis was that the complexity of erosion models roughly tended to match that of the flow models to which they were associated (Section 2). However, the search for determinants of erosion modelling goes through several other stages, as announced now in Section 2.1.2. An explanation will be added **L168-175**.

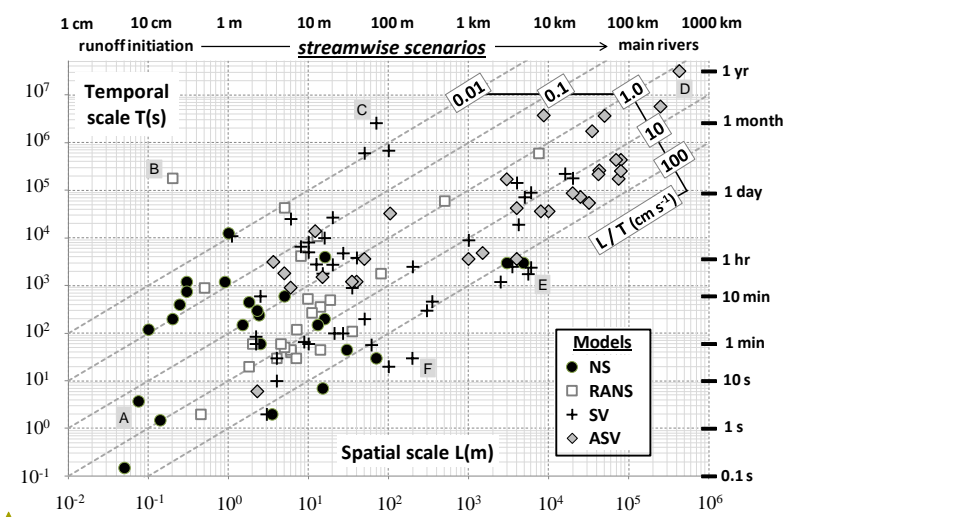
"On the one hand, this advocates the examination of erosion issues from the angle of decreasing refinements of the "flow and erosion" models seen as a whole (e.g. expecting the most complicated erosion processes to be out of reach of the simplest combined models). On the other hand, there might be a certain disconnection between the refinement of the flow model and that of the chosen friction and erosion models, so the determinants of modelling choices should also be sought elsewhere: in flow typologies dictated by friction and flow retardation processes but also in "erosion types", seen through a dimensionless descriptor (Section 3)."

- Most of the paper's figures are quite dense, and I suggest to comment on these plots more broadly to guide the reader across them;

R: We agree. To do so, we will focus on a few "textbook cases", i.e. 6 cases now explicitly referred to in Fig.2, 3, 6a and 7a, shown by letters A to F, detailed in the new Table 1 and in

the associated paragraph added in the corpus L507-523. The legends of the cited figures have been slightly modified to mention these textbook cases.

The new Fig.2 (L464) is hereunder and only the last sentence of its legend has been modified to mention the A to F sketches.



Mis en forme : Police :(Par défaut)
Arial, Couleur de police : Bleu

Figure 2 – How increasing (L, T) spatiotemporal scales of the flow domain tend to be associated with decreasing complexity in the choice of flow models, sorted here into four levels of refinement: Navier-Stokes (NS), Reynolds-Averaged Navier-Stokes (RANS), Saint-Venant (SV) or Approximations to Saint-Venant (ASV). A transverse analysis involves forming L/T ratios, searching for clues to model selection according to these "system evolution velocities" or governed by flow typologies that would exhibit specific L/T ratios. This figure was assembled from information available in the studies cited in Appendix A, selecting six textbook cases (sketches A to F, Table 1) for illustration.

The new Table 1 is the following.

Case	Context	Authors	Model refinement	Spatiotemporal scales					Flow typology [‡]	Dimensionless numbers [§]					
				L (m)	T (s)	H (m)	L/T (m s ⁻¹)	H/L [†] (-)		T*	Re	Fr	S (%)	Λ _z	θ
A	Film flow	Charpin & Myers (2005)	NS	0.075	3.75	0.003	0.02	0.04	O	5	300	0.11	10	8.0	-
B	Laminar dynamics	Charru et al. (2004)	RANS	0.2	1.8 10 ⁵	0.007	1.1 10 ⁻⁵	0.035	O	6428	50	0.02	<0.01	12.1	0.14
C	Pool-riffles	Rathburn & Wohl (2003)	SV	70	2.6 10 ⁶	0.47	3.5 10 ⁻³	6.7 10 ⁻³	B	7.8 10 ³	7.1 10 ³	0.69	1.1	5108	34.1
D	Amazon River	Trigg et al. (2009)	ASV	4.3 10 ³	3.15 10 ⁶	10	1.4 10 ⁻³	2.3 10 ⁻³	F	58.5	8 10 ⁵	0.05	<0.01	6600	-
E	Step-pools	Grant et al. (1990)	SV	5530	1755	0.87	3.15	1.5 10 ⁻⁴	Hg	1.0	2.7 10 ³	1.03	4.5	1.25	-
F	Step-pools	Chin (1999)	SV	197.25	30	0.50	6.58	0.025	Hg	1.21	4.0 10 ³	3.58	6.25	1.22	-

[†] See section 3.1.2 - H/L is the fineness ratio of the flow

[‡] See Section 3.2 - O: Overland, Hg: High-gradient, B: Bedforms, F: Fluvial

[§] See Section 3.3 - T*: dimensionless period, Re: Reynolds number, Fr: Froude number, S: slope, Λ_z inundation ratio, θ Shields number

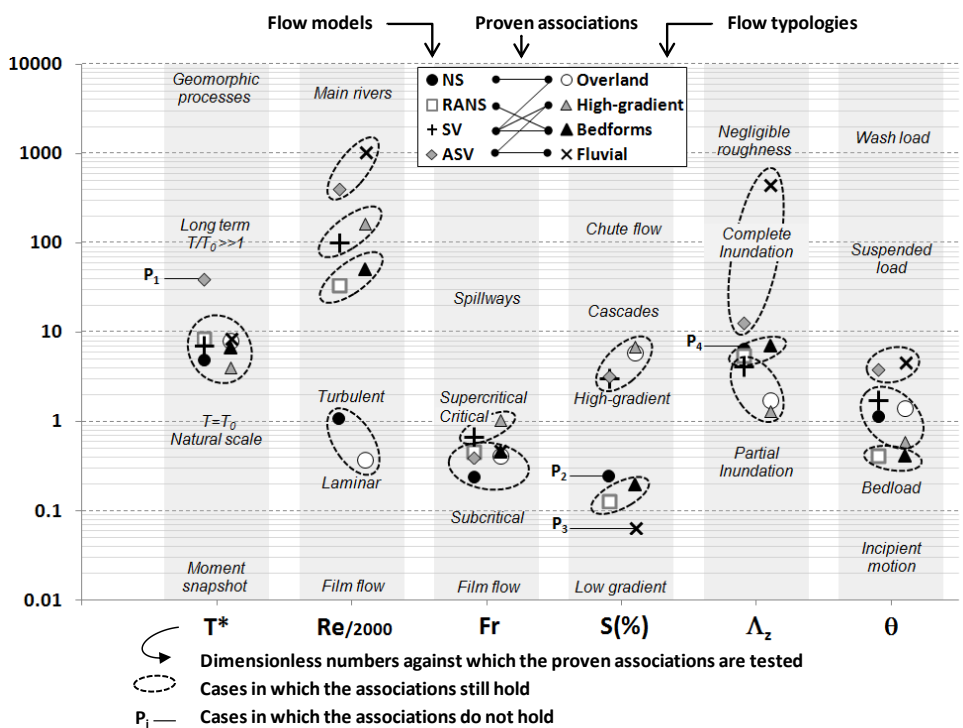
Table 1 - Six textbook cases representing an approximate envelope of all the tested cases in the L-T plane of Fig.2, where L is the spatial scale (length of the flow domain) and T the temporal scale (duration of the process studied). Spatiotemporal scales are the determinants of modelling choices discussed in Section 3.1. The additional influence of flow typology and dimensionless numbers are discussed in Sections 3.2 and 3.3.

The added paragraph is:

"To take a few examples and guide the reader through the arguments and the figures of this paper, Table 1 gathers the information available for the six textbook cases outlined by sketches A to F in Fig.2. The selected studies represent a wide variety of cases (drawing an approximate envelop of cases in the L-T plane of Fig.2) followed in the forthcoming stages of the analysis and associated figures in Section 3.1.2 (determinants of modelling choices in the L-H plane, Fig.3), Section 3.2 (determinants sought in flow typology, Fig.6a and 7a) and Section 3.3 (determinants sought in the values of dimensionless numbers attached to the flow)." (L507-513)

- In particular, the last figure of the paper is the most interesting result of the entire manuscript, and I suggest the authors to expand the description/comment on this very interesting result. I do not think the interpretation of this plot is trivial at all, so I believe its significance should be better emphasized in the paper text.

R: We agree. However, instead adding more elements in the text, we opted for some modifications of the old Fig.9 (now Fig.10, L865) to make it more self-explanatory, keeping its legend unchanged.



I do not have other specific comments, as the paper seems to be very accurate. The Figures are of good quality, referencing is appropriate and the discussion is clear and concise. Therefore, I congratulate the authors for the overall quality of the manuscript.

R: Thank you.

Response to Reviewer #2

The paper presents an interesting attempt to draw links between different modelling approaches and to find appropriate time and length scales for different types of models. The approach adopted in the paper intends to be a general approach considering very different types of flows, from runoff to flows in large rivers.

R: We thank Reviewer #2 for his positive comments. We totally agree with his comments, and in the revised version we will introduce responses to all points raised by Reviewer #2 as shown below.

However, it must be stressed that this generalization still remains in the field of hydrology, with a point of view that is not as general as it could be. In particular, the Navier-Stokes (NS) equations are mentioned, but without being considered in their general fluid mechanics framework. So, the NS model is presented as the most general one, which is certainly the case, but turbulence is not discussed. However, considering that the flow velocity is the sum of a mean velocity and a fluctuating component, the NS equations can be solved to resolve as many as possible of the turbulence scales in DNS type simulations, also in flows with significant water depths. These DNS simulations are not discussed here, and NS models always appear in the “runoff” range of applications, which is quite limiting. Of course, if one remembers that the general review concerns hydrological modelling, then it becomes acceptable. But if this is the intention of the authors, then it should be stated much more clearly in the objectives of the paper.

R: The paper is indeed turned towards applications in the field of hydraulics and hydrology. The word “hydrology” was mentioned in the title for disambiguation, especially for readers’ specialists in fluid mechanics who would expect the usual analytical framework. This “hydrological” option will be recalled for clarity, in one word or two, at several places in the introductory parts of the manuscript (abstract L21, introduction L85 and L142-143).

However, high-precision hydraulics (for example) requires the NS models and may involve various (turbulence) scales and flow structures. We have thus followed the recommendation to mention the possible context-dependent strategies (DNS, LES, RANS) to solve these equations, which hopefully restore a bit more genericity. An additional comment will be added (L196-205).

“There are many turbulence models (e.g. DNS-Direct Numerical Simulations, LES-Large Eddy Simulations and RANS-Reynolds-Averaged Navier-Stokes) suitable for free-surface flow modelling (Katopodes & Bradford 1999). Direct Numerical Simulations explicitly resolve all turbulence scales at the cost of more than Re^3 calculations (Härtel 1996) while Large Eddy Simulations (Smagorinsky 1963, Leonard 1974) filter out the smallest scales and resolve only the larger ones. The RANS equations (Smith & McLean 1977, Rödi 1988) do not resolve any scale but the stress terms used for their closure have proven useful for the modelling of near-bed turbulent patterns (see next subsection). The general trend is that improvements in efficiency of the algorithms have approximately kept pace with exponential improvements in computer power over the past 50 years (Moore 1965, Mavriplis 1998, Koomey et al. 2010) which tends to push the limitations of DNS and LES further away.”

In a similar way, it then appears quite strange to read the word “turbulence” only when RANS models are discussed.

R: The term will be added L159, Section 2.1.1: “(RANS: Reynolds 1895, for turbulent flows)”

Indeed, RANS models were developed because performing DNS simulations to resolve all turbulence scales is impossible in practice due to excessive computational cost. Current research tends to push this limitation of NS still further away because of increasing available computational power using e.g. parallel computing. This is also an issue that deserves to be discussed.

R: These points will be mentioned L196-200.

"There are many turbulence models (e.g. DNS-Direct Numerical Simulations, LES-Large Eddy Simulations and RANS-Reynolds-Averaged Navier-Stokes) suitable for free-surface flow modelling (Katopodes & Bradford 1999). Direct Numerical Simulations explicitly resolve all turbulence scales at the cost of more than Re^3 calculations (Härtel 1996) while Large Eddy Simulations (Smagorinsky 1963, Leonard 1974) filter out the smallest scales and resolve only the larger ones."

On the way erosion processes are handled, there can also be some debate. The references used by the authors are certainly pertinent in the field. However, the attempt of classifying the different approaches for erosion and grain movement at the same level as the NS, RANS, SV and ASV models is questionable. The distinction is not so clear, and a different classification, not directly linked to the flow models, but rather to the type of grain movement considered would maybe have been more appropriate.

R: We agree that erosion issues are not fully addressed with such a "parallel" strategy in terms of decreasing model refinements, which only provides a trend. Section 2.1.2 will be modified to clarify this point (L168-175) and to indicate that complementary indications on the determinants of modelling choices regarding erosion will be found in Section 3.

"or at the scale of the erodible bed asperities. On the one hand, this advocates the examination of erosion issues from the angle of decreasing refinements of the "flow and erosion" models seen as a whole (e.g. expecting the most complicated erosion processes to be out of reach of the simplest combined models). On the other hand, there might be a certain disconnection between the refinement of the flow model and that of the chosen friction and erosion models, so the determinants of modelling choices should also be sought elsewhere: in flow typologies dictated by friction and flow retardation processes but also in "erosion types", seen through a dimensionless descriptor (Section 3)."

These new lines (L168-175) mention a dimensionless descriptor for erosion (which will be the Shields number) which refers to phenomenologies that are not directly related to the NS, RANS, SV or ASV level, but rather to friction, bedforms and flow retardation processes as "proxys" for particle pick-up. What we intend to do is to introduce a new figure that shows a generalized Shields diagram. This would first offer an alternative to the reasoning in terms of refinement levels and second explicitly refer to the different erosion-transportation-deposition modes. This new Fig.9 comes at the end of Section 3.3.1 and is introduced by modifications in the text L826-834.

" This number seems appropriate for most erosion issues because it has been widely applied and debated in the literature (Coleman 1967, Ikeda 1982, Wiberg & Smith 1987, Zanke 2003, Lamb et al. 2008) and also because of its numerous possible adaptations (Neill 1968, Ouriemi et al. 2007, Miedema 2010) to various flow typologies and non-uniform or poorly-known bed conditions. An impressive review on the use of the Shields number to determine incipient motion conditions, over eight decades of experimental studies, may be found in Buffington & Montgomery (1997). Finally, Fig.9 provides a generalized Shields diagram that includes motion threshold criteria under the effects of high or low particle exposure (Miedema 2010) or for laminar flows, also indicating the conditions of significant suspension (Wright & Parker 2004)."

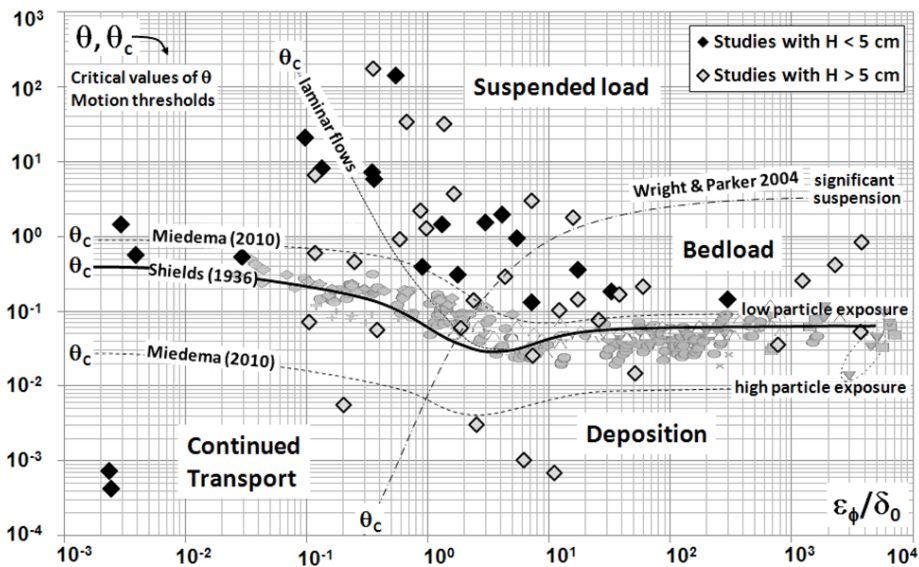


Figure 9 - Generalized dimensionless Shields diagram that summarizes the conditions and regimes of sediment transport or deposition, from the relative values of the Shields parameter (θ) and incipient motion criterion (θ_c). The X-axis bears the values of the ratio of particle size (ϵ_ϕ) on the depth of the laminar sublayer (δ_0). The diamonds refer to the studies cited in Appendix A that deal with erosion issues: black diamonds for studies in which flow depth is $H < 5$ cm, grey diamonds otherwise. Data in the background show the critical θ_c values reported in the wide Buffington & Montgomery (1993) review of incipient motion conditions for varied flow regimes, particle forms and exposures.

In particular, the authors mention that the SV framework offers a wide field for innovative research about sediment transport, which is certainly the case. But in these recent researches, many different types of sediment transport models are considered, depending also on the necessary level of simplification of the reality that is required. Indeed, the detailed composition of the soli to be eroded is not always known, or it is not possible to include that level of detail in the representation. So it is necessary to resort to averaging concepts, such as a representative grain diameter, then some factors to account for the non-uniformity of the grain-size distribution.

R: This is now explicitly mentioned in the responses to the previous comments.

The concentration of sediment in the flow could also be discussed: debris flows or mud flows are not handled in the same way as clear-water flows with sediment transport, and this distinction does not really appear here.

R: We fully agree again and your request incites us to reintroduce several elements that we had previously discarded from our working versions (as the Shields diagram) in an attempt to make the manuscript shorter.

Mis en forme : Police : (Par défaut)
Arial, Couleur de police : Bleu

In the discussion paper, we only mentioned hyperconcentrated flows and stratification (i.e. density) effects for sediment laden flows, not really addressing the effect of flow density (water+sediments mixture) on modelling options.

As far as we know, the trend is to use higher-level models when the water-sediment couplings become stronger. Again, the SV level allows many adaptations and strategies, but we feel there was a lack regarding the applications of the NS and RANS to dense, debris or avalanche flows, for example. A few lines on the subject were already present in Sections 2.2.2 and 2.4.2 but we will add some more literature elements in Section 2.2.2. [\(L216-220\)](#)

“Such couplings may be sorted by increasing sediment loads, from dispersed multiphase flows (Parker & Coleman 1986, Davies et al. 1997) to density currents (Parker et al. 1986), hyperconcentrated flows (Mulder & Alexander 2001) and up to debris flows (e.g. Bouchut et al. 2003, Bouchut & Westdickenberg 2004), the latter derived as mathematical generalizations of the well-known Savage & Hütter (1989, 1991) avalanche models over explicit, pronounced topographies.”

Minor comments and detailed suggestions for improvement will be submitted later as an attached file.

R: Thank you.

However, it must be stressed that this generalization still remains in the field of hydrology, with a point of view that is not as general as it could be.

R: Let us return to this phrase in the second comment of Reviewer#2 which pushed us to reconsider the conclusion of the paper, thus to formulate its concluding message in a quite different way. First, we split the conclusion in two and the previously existing part becomes Section 4.1 "Outcomes of this review" [\(L880\)](#). Second, the added part is Section 4.2 "Research challenges in hydrology and philosophy of modelling" [\(L957-1021\)](#) including a new Fig.11 that summarizes what has been done and what should be done in complement. Figure 11 finds itself at the tilting point between Sections 4.1 and 4.2.

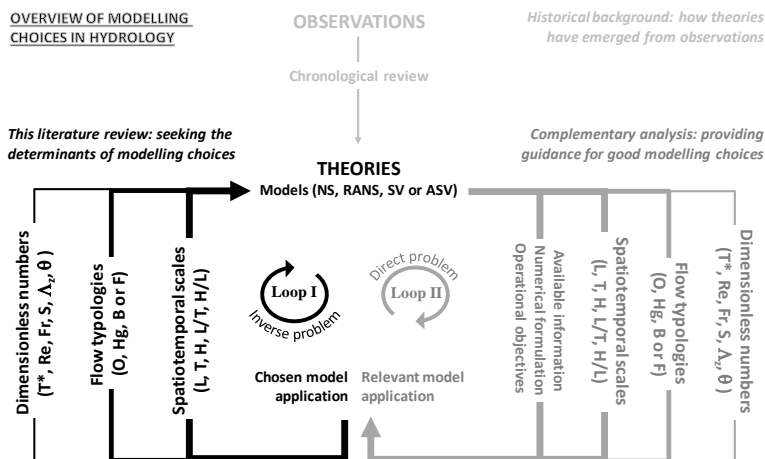


Figure 11 – This figure provides a simplified overview of the available modelling choices in hydrology, in three distinct colours associated with specific research purposes or disciplines, showing the position of the present review relative to the others. The pale grey section aims at understanding how the available flow models have emerged from observations and early formulations of the flow equations, focusing on their conditions of validity i.e. the successive hypotheses made during their derivation. The black section recalls the procedure followed in this review paper (Loop I, "inverse problem"). Literature sources are processed through a procedure that analyses how the spatiotemporal scales (spatial scale L, time scale T, flow depth H, L/T and H/L ratios), then flow typology (Overland O, High-gradient Hg, Bedforms B or Fluvial F) and dimensionless numbers (dimensionless period T^* , Reynolds number Re , Froude number Fr , bed slope S, inundation ratio \square_z , Shields parameter \square) determine the choice of a flow model (Navier-Stokes NS, Reynolds-Averaged Navier-Stokes RANS, Saint-Venant SV or approximations ASV). Suggested in medium grey on the right are the scope and principles of future research challenges that would address the "what should be done?" (Loop II, "direct problem") question in echo to the current "what has been done?" concern (Loop I).

On the one hand, the added section 4.2 discusses pending challenges and possible approaches quite specific to the fields of hydrology and hydraulics. On the other hand, it reintroduces very generic concepts and decision rules (hence the title "philosophy of modelling") in suggesting to select the approaches that respect the principle of parsimony.

"This review has sought the determinants of modelling choices in hydrology (Figure 11, Loop I) from the basis provided by literature sources, without any intention to provide recommendations. However, for most practical applications, the starting point is the definition of a scope and the endpoint is the evaluation of the objective function to evaluate the success or the failure of the chosen modelling strategy. A question thus arises on how to guide the modeller in the choice of an adequate model, in function of given, approximately known spatiotemporal scales, flow typology and dimensionless numbers (Figure 11, Loop II). According to the principle of parsimony, modellers should seek the simplest modelling strategy capable of (i) a realistic representation of the physical processes, (ii) matching the performances of more complex models and (iii) providing the right answers for the right reasons.

- (i) Throughout the last decades, an important change of the scope of free-surface flow modelling applications has taken place, with subsequent changes in the objective functions resorted to. The development of hydrological and hydraulic sciences has been directly linked to the progresses in understanding processes, in theoretical model development (e.g. computational facilities: numerical techniques, data assimilation, thorough model exploration, inverse calculus) and in data acquisition (new devices, remote sensing, LiDAR). "It may seem strange to end a review of modelling with an observation that future progress is very strongly linked to the acquisition of new data and to new experimental work but that, in our opinion, is the state of the science" (Hornberger & Boyer 1995).

- (ii) However, there remains an important need for research on classical free-surface flow (hydrological or hydraulic) modelling for engineering applications in predicting floods, designing water supply infrastructures and for water resources management, from the headwater catchment to the regional scale. More recently, free-surface flow modelling has become an indispensable tool for many interdisciplinary projects, such as predicting pollution and/or erosion incidents, the impact of anthropogenic and climate change on environmental variables such as water, soil, biology, ecology, or socio-economy and ecosystemic services. The direct consequence is a significant increase of the complexity of the objective function, from simple mono-site (e.g. one-point), mono-variable (e.g. the water depth) and mono-criterion (e.g. the error on peakflow) to complex multi-site (e.g. large number of points within a catchment), multi-variable (e.g. water depth, hydrograph, water table, concentrations,

ecological indicators, economic impact) and multi-criteria (e.g. errors on peakflow, volume, RMSE) objective functions.

- (iii) There is often a mismatch between model types, site data and objective functions. First, models were developed independently from the specificities of the study site and available data, prior to the definition of any objective function. In using free-surface flow models, the context of their original purpose and development is often lost, so that they may be applied to situations beyond their validity or capabilities. Second, site data are often collected independently of the objectives of the study. Third, the objective function must be specific to the application but also meet standard practices in evaluating model performance, in order to compare modelling results between sites and to communicate the results to other scientists or stakeholders. The known danger is to use flow and erosion equations outside their domains of validity (i.e., breaking the assumptions made during their derivation) then to rely on the calibration of model parameters as for technical compensations of theoretical flaws, at the risk of losing the physical sense of model parameters, creating equifinality and obtaining the "right results for the wrong reason" (Klemeš 1986). Choosing the right model for the right reason is crucial but the identification of the optimal data-model couple to reach a predefined objective is not straightforward. We need a framework to seek the optimum balance between the model, data and the objective function as a solution for a hydrological or hydraulic problem, on the basis of the principle of parsimony. The latter follows a famous quote often attributed to Einstein, that "everything should be made as simple as possible, but not simpler" which somehow originates in the philosophy of William of Ockham (1317) (*Numquam ponenda est pluralitas sine necessitate* [Plurality must never be posited without necessity]) or may even be traced back to Aristotle's (~350 BCE) *Analytica Posteriora* that already advocated demonstrations relying on the fewest possible number of conjectures, i.e. the dominant determinisms.

Finally, analytical procedures for free-surface flows and erosion issues necessitates a comprehensive analysis of the interplay between models (assumptions, accuracy, validity), data requirements and all contextual information available, encompassed in the "signature" of any given application: model refinement, spatiotemporal scales, flow typology and scale-independent description by dimensionless numbers. This review helps the modeller positioning his (or her) case study with respect to the modelling practices most encountered in the literature, without providing any recommendation. A complementary step and future research challenge is to decipher relevant modelling strategies from the available theoretical and practical material, resorting to the same objects, the previously defined signatures. Its purpose clearly is to address the "which model, for which scales and objectives?" question. A complete analytical framework, comprised of both loops, would provide references and guidelines for modelling strategies. Its normative structure in classifying theoretical knowledge (the mathematics world, equations and models) and contextual descriptions (real-life physical processes, scales and typologies) hopefully makes it also relevant for other Earth Sciences."

Additional references

Bouchut, F., Mangeney-Castelnau, A., Perthame, B., Vilotte, J.-P. (2003). A new model of Saint-Venant and Savage-Hutter type for gravity driven shallow water flows, *Comptes-Rendus de l'Académie des Sciences de Paris*, 336, pp.531-536.

Bouchut, F. and Westdickenberg, M. (2004). Gravity-driven shallow water models for arbitrary topography, *Communications in Mathematical Sciences*, 2, pp.359-389. Härtel, C. (1996). Turbulent flows: direct numerical simulation and large-eddy simulation, In: *Handbook of Computational Fluid Mechanics*, R. Peyret Ed., Elsevier, New York, pp.283-338.

Katopodes, N. D. and Bradford, S. F. (1999). Mechanics of overland flow, *Proceedings of the International Workshop on Numerical Modelling of Hydrodynamic Systems*, Zaragoza, Spain, 21-24 June 1999.

Klemes, V. (1986). Dilletantism in hydrology: Transition or destiny?, *Water Resources Research*, 22, pp.177-188.

Koomey, J., Berard, S., Sanchez, M. and Wong, H. (2010). Implications of historical trends in the electrical efficiency of computing, *IEEE Annals of the history of computing*, 33 (3), pp.46-54.

Leonard, A. (1974). Energy cascade in large-eddy simulation of turbulent channel flow, *Advances in Geophysics*, 18, pp.237-248.

Mavriplis, D. (1998). On Convergence Acceleration Techniques for Unstructured Meshes. Technical Report ICASE No. 98-44 and NASA/CR-1998-208732, Institute for Computer Applications in Science and Engineering, NASA Langley Research Center.

Moore, G. E. (1965). Cramming More Components onto Integrated Circuits, *Electronics*, pp. 114–117.

Rödi, W. (1988). Turbulence models and their application in hydraulics - a state of the art review, *International Association for Hydraulic Research*, Delft, The Netherlands, 47p.

Savage, S. B. and Hutter, K. (1989). The motion of a finite mass of granular material down a rough incline, *Journal of Fluid Mechanics*, 199, pp.177-215.

Savage, S. B. and Hutter, K. (1991). The dynamics of avalanches of granular materials from initiation to runout. Part I: Analysis, *Acta Mechanica*, 86, pp.201-223.

Smagorinsky, J. (1963). General circulation experiments with the primitive equations, *Monthly Weather Review*, 91 (3), pp.99-164.

Wright, S. and Parker, G. (2004). Flow Resistance and Suspended Load in Sand-Bed Rivers: Simplified Stratification Model, *Journal of Hydraulic Engineering*, 130 (8), pp.796-805.

1

2

3 **Determinants of modelling choices for 1-D free-surface flow and**
4 **erosion issues in hydrology: a review**

5

6

7

8 Bruno Cheviron^a, Roger Moussa^b

9

10

11 ^a IRSTEA, UMR G-EAU "Gestion de l'Eau, Acteurs et Usages", 361 rue Jean-François Breton, BP
12 5095, 34196 Montpellier Cedex 5, France.

13 bruno.cheviron@irstea.fr

14 ^b INRA, UMR LISAH, "Laboratoire d'étude des Interactions entre Sol - Agrosystème -
15 Hydrosystème", 2 Place Pierre Viala, 34060 Montpellier Cedex 1, France.

16 moussa@supagro.inra.fr

17

18 **Abstract**

19
20 This review paper investigates the determinants of modelling choices, for numerous applications of
21 1-D free-surface flow and erosion equations [in hydrology](#), across multiple spatiotemporal scales. We
22 aim to characterize each case study by its signature composed of model refinement (Navier-Stokes:
23 NS, Reynolds-Averaged Navier-Stokes: RANS, Saint-Venant: SV or Approximations of Saint-
24 Venant: ASV), spatiotemporal scales (domain length: L from 1 cm to 1000 km; temporal scale: T from
25 1 second to 1 year; Flow depth: H from 1 mm to 10 m), flow typology (Overland: O, High gradient:
26 Hg, Bedforms: B, Fluvial: F) and dimensionless numbers (Dimensionless time period T^* , Reynolds
27 number Re , Froude number Fr , Slope S , Inundation ratio Λ_z , Shields number θ). The determinants of
28 modelling choices are therefore sought in the interplay between flow characteristics, cross-scale and
29 scale-independent views. The influence of spatiotemporal scales on modelling choices is first
30 quantified through the expected correlation between increasing scales and decreasing model
31 refinements, identifying then flow typology a secondary but mattering determinant in the choice of
32 model refinement. This finding is confirmed by the discriminating values of several dimensionless
33 numbers, that prove preferential associations between model refinements and flow typologies. This
34 review is intended to help each modeller positioning his (her) choices with respect to the most frequent
35 practices, within a generic, normative procedure possibly enriched by the community for a larger,
36 comprehensive and updated image of modelling strategies.

37

38 **Keywords**

39 Free-surface flow, modelling strategy, cross-scale analysis, flow typology, dimensionless numbers.

40

41	Summary	
42	1 Introduction	4
43	2 Flow models	7
44	2.1 List of flow models.....	7
45	2.1.1 <i>Water flow</i>	7
46	2.1.2 <i>Erosion</i>	7
47	2.2 Navier-Stokes.....	8
48	2.2.1 <i>Water flow</i>	8
49	2.2.2 <i>Erosion</i>	9
50	2.3 Reynolds-Averaged Navier-Stokes.....	10
51	2.3.1 <i>Water flow</i>	10
52	2.3.2 <i>Erosion</i>	11
53	2.4 Saint-Venant.....	12
54	2.4.1 <i>Water flow</i>	12
55	2.4.2 <i>Erosion</i>	13
56	2.5 Approximations to Saint-Venant.....	16
57	2.5.1 <i>Water flow</i>	16
58	2.5.2 <i>Erosion</i>	17
59	3 Determinants of modelling choices	18
60	3.1 Spatiotemporal scales.....	19
61	3.1.1 <i>Influence of domain length (L) and time scale (T)</i>	19
62	3.1.2 <i>Influence of domain length (L) and flow depth (H)</i>	22
63	3.1.3 <i>Influence of domain length (L), time scale (T) and flow depth (H)</i>	24
64	3.2 Flow typology.....	24
65	3.2.1 <i>From friction laws and bed topography to flow characteristics</i>	24
66	3.2.2 <i>From flow characteristics to flow typologies</i>	27
67	3.2.3 <i>Influence of flow typologies on modelling choices</i>	30
68	3.3 Dimensionless numbers.....	35
69	3.3.1 <i>Contextual dimensionless numbers</i>	35
70	3.3.2 <i>Influence of the dimensionless numbers</i>	38
71	4 Conclusion	41
72	4.1 Outcomes of this review.....	41
73	4.2 Research challenges in hydrology and philosophy of modelling.....	44
74	Appendix A. References used in the Figures	47
75	Acknowledgements	49
76	References	50
77		

78 **1 Introduction**

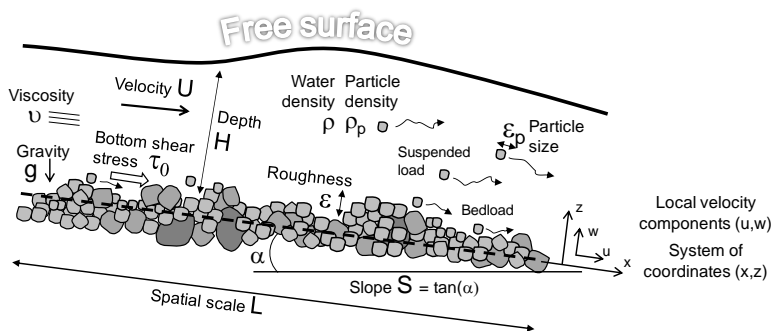
79 Free-surface flow models cover a wide range of environmental and engineering applications, across
80 multiple spatiotemporal scales, through successive flow aggregations over various bed topographies:
81 these govern both the qualitative (flow typology) and quantitative (dimensionless numbers) flow
82 characteristics. Each case study may thus be positioned along "streamwise scenarios" (from runoff
83 initiation to the main rivers) from unequivocal indications of the spatiotemporal scales, flow typology
84 and associated dimensionless numbers. This literature review investigates the determinants of choices
85 made for 1-D free-surface flow and erosion modelling in hydrology, seeking links between contextual
86 information (spatiotemporal scales, flow typologies, dimensionless numbers) and conceptual
87 descriptions (refinement of the flow equations or, equivalently, richness of the physical basis). The
88 entire set of descriptors, i.e. model refinement, spatiotemporal scales, flow typology and
89 dimensionless numbers, constitutes the signature of a study, which is the open normative procedure
90 designed to allow comparisons between studies and to be fed by the community.

91
92 For the sake of genericity, this review addresses a wide range of spatiotemporal scales, starting at
93 the smallest plot scales (spatial scale: domain length $L < 10$ m; time scale: duration of the process
94 $T < 10$ s; flow depth: $H < 1$ cm, Fig. 1), those of runoff genesis, overland flow hydraulics and detailed
95 particle-scale physics (Horton 1945, Emmett 1970, Feng & Michaelides 2002, Schmeeckle & Nelson
96 2003). The intermediate scales of catchment and hillslope processes are these expected to exhibit the
97 widest variety of flow typologies thus modelling strategies (Croke & Mockler 2001, Parsons et al.
98 2003, Aksoy & Kavvas 2005). The larger river basin scales ($L > 100$ km; $T > 10$ days; $H > 1$ m) are also
99 handled here, relevant for river flow modelling, flood prediction and water resources management
100 (Nash & Sutcliffe 1970, Rosgen 1994, Loucks & van Beek 2005) with regional surface-subsurface
101 interactions (De Marsily 1986), non-point pollution, fluvial sediment budgets and global
102 biogeochemical cycles (Walling 1983, Milliman & Syvitski 1992, Syvitski & Milliman 2007).

103

104 On the Earth's surface, flow aggregation in the streamwise direction occurs across several
 105 geomorphic thresholds (Kirkby 1980, Milliman & Sivitsky 1992, Church 2002, Paola et al. 2009),
 106 through a succession of flow typologies (Emmett 1970, Grant et al. 1990, Rosgen 1994, Montgomery
 107 & Buffington 1997). Flow aggregation in space and time is described, through the width function and
 108 geomorphological unit hydrograph concepts (Kirkby 1976, Robinson et al. 1995, Agnese et al. 1998),
 109 under the angle of connecting-scale hydrological and sedimentological pathways (see the review by
 110 Bracken et al. 2013) or debating the merits of similitude laws versus upscaling issues in the
 111 description of hydrological processes (Strahler 1956, Blöschl and Sivapalan 1995, Slaymaker 2006).
 112 An alternative consists in examining the scale matching between available data and modelling aims
 113 (Lilburne 2002). This raises technical (contextual) as well as strategic (conceptual) issues, handled
 114 here from an overview on the most popular modelling practices, confronting the theoretical refinement
 115 of flow models to the specific nominal scales of the processes at play.

116



117

118 **Figure 1 - Quantities most often used in the literature of free-surface flow and erosion modelling, with**
 119 **explicit reference to the (L, T, H) spatiotemporal scales of interest. This review is limited to 1D (x) spatial**
 120 **representations for simplicity, focusing on the streamwise (x) component of the mass and momentum**
 121 **conservation equations. The streamwise length (L) and velocity (U) suggest a natural time scale $T_0=L/U$**
 122 **for the propagation of information, waves or perturbations, to be compared with the time scales (T) opted**
 123 **for in the literature.**

124

125 Many papers or handbooks have summarised free-surface flow modelling and numerical
126 techniques in hydraulics (King & Brater 1963, Abbott 1979, Cunge et al. 1980, Carlier 1980, French
127 1985) or hydrology (Chow 1959, Kirkby 1978, Beven 2000) for various contexts, purposes and flow
128 typologies. Less works have discussed the concern of *ad hoc* friction laws (Leopold et al. 1960,
129 Gerbeau & Perthame 2001, Nikora et al. 2001, Roche 2006, Burguete et al. 2008), at the microscopic
130 or macroscopic scales (Richardson 1973, Jansons 1988, Priezjev & Troian 2006, Smith et al. 2007,
131 Powell 2014) although friction, flow retardation and energy dissipation processes are closely related to
132 bedforms, thus plausibly govern flow typologies then, possibly, modelling choices. Often outside any
133 focus on friction, numerous works have provided wide overviews on erosion modelling (Ritchie &
134 McHenry 1990, Lafren et al. 1991, Merritt et al. 2003, Aksoy and Kavvas 2005, Boardman 2006).
135 Erosion models that lean on the most sophisticated flow models calculate explicit particle detachment,
136 transport and deposition from velocity fields or flow energetics (Vanoni 1946, Hino 1963, Lyn et al.
137 1992, Mendoza & Zhou 1997) while reduced complexity models either assume the "transport
138 capacity" (Foster & Meyer 1972, Bennett 1974) or "transport distance" schools of thoughts (see details
139 in Wainwright et al. 2008).

140

141 This multidisciplinary review (hydrology, hydraulics, fluid mechanics and erosion science)
142 searches for the determinants of modelling choices. It focuses on hydrology but borrows from
143 hydraulics and fluid mechanics, also when addressing erosion issues. The methodology consists in
144 defining the "signature" of each case study as the chosen model refinement and the given flow
145 typology, spatiotemporal scales and dimensionless numbers, hypothesizing the conceptual element
146 (model refinement) is the consequence of the contextual elements. The paper is organized as follows:
147 section 2 sorts the flow equations into four levels of refinement, section 3 plots these refinements
148 versus the spatiotemporal scales of the studies, also depicting the influence of flow typologies and
149 dimensionless numbers. Section 4 discusses the results and future research leads. Some of the best
150 documented references among the cited literature have been gathered in Appendix A: most figures in
151 this manuscript were plotted from this database.

152

153 **2 Flow models**

154 **2.1 List of flow models**

155 *2.1.1 Water flow*

156 Free-surface flow equations in the literature may roughly be sorted into four levels of decreasing
157 refinement, from the richness of their physical basis. The choice made here includes the Navier-Stokes
158 equations (noted NS: Navier 1822, Stokes 1845), their average in time termed Reynolds-Averaged
159 Navier-Stokes equations (RANS: Reynolds 1895, [for turbulent flows](#)), the depth-averaged Saint-
160 Venant equations (SV: Saint-Venant 1871) and further approximations (referred to as ASV), among
161 which the Diffusive Wave (DWE: Hayami 1951) and Kinematic Wave Equations (KWE: Iwagaki
162 1955, Lighthill & Whitham 1955).

163 *2.1.2 Erosion*

164 The associated erosion equations (not shown) are based on a representation of detachment and
165 transport on hillslopes (Bennett 1974, Van Rijn 1984a, b, Wainwright et al. 2008), in streams (Einstein
166 1950) or through the channel network (Du Boys 1879, Exner 1925, Hjulström 1935, Shields 1936,
167 Bagnold 1956). Friction is the link between water flow and erosion issues in terms of physical
168 processes at play at the particle scale, or at the scale of the erodible bed asperities. [On the one hand,](#)
169 [this advocates the examination of erosion issues from the angle of decreasing refinements of the "flow](#)
170 [and erosion" models seen as a whole \(e.g. expecting the most complicated erosion processes to be out](#)
171 [of reach of the simplest combined models\). On the other hand, there might be a certain disconnection](#)
172 [between the refinement of the flow model and that of the chosen friction and erosion models, so the](#)
173 [determinants of modelling choices should also be sought elsewhere: in flow typologies dictated by](#)
174 [friction and flow retardation processes but also in "erosion types", seen through a dimensionless](#)
175 [descriptor \(Section 3\). ~~However, the scope here is not to review the choices made for friction~~](#)

176 ~~modelling; friction phenomena, with the associated flow retardation and energy dissipation processes,~~
177 ~~are rather considered for their influence on flow typologies, as discussed later in the manuscript.~~
178

179 **2.2 Navier-Stokes**

180 *2.2.1 Water flow*

181 The Navier-Stokes (NS) equations have suitable simplifications for the shallow water cases
182 ($L \gg H$) commonly used to describe free-surface flows. The three-dimensional fluid motion problem is
183 reduced here to a two-dimensional description, whose projection along the streamwise axis writes:

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + w \frac{\partial u}{\partial z} + \frac{1}{\rho} \frac{\partial p}{\partial x} = g_x + \frac{1}{\rho} \frac{\partial \tau}{\partial x} \quad (1)$$

184 where x is the longitudinal distance [L], z the vertical coordinate [L], t is time [T], u is the local water
185 velocity in x [LT^{-1}], ρ is water density [ML^{-3}], g_x is the projection of gravity g on x [LT^{-2}] and τ is the
186 tangential stress due to water [$ML^{-1}T^{-2}$] noted τ_0 on the bed in Fig. 1.

187

188 The Navier-Stokes equations stay valid throughout the full range of flow regimes, scales and
189 contexts. They are preferentially used where much complexity is needed, often when relevant
190 simplified flow descriptions could not be derived, for example for particle-scale applications (Chen &
191 Wu 2000, Wu & Lee 2001, Feng & Michaelides 2002), overland flow (Dunkerley 2003, 2004) or
192 flows over pronounced bedforms (Booker et al. 2001, Schmeeckle & Nelson 2003). A very wide
193 review of numerical methods and applications for the NS equations is provided by Gresho & Sani
194 (1998) and a benchmark of numerous solvers by Turek (1999).

195

196 There are many turbulence models (e.g. DNS-Direct Numerical Simulations, LES-Large Eddy
197 Simulations and RANS-Reynolds-Averaged Navier-Stokes) suitable for free-surface flow modelling
198 (Katopodes & Bradford 1999). Direct Numerical Simulations explicitly resolve all turbulence scales at
199 the cost of more than Re^3 calculations (Härtel 1996) while Large Eddy Simulations (Smagorinsky

200 | 1963, Leonard 1974) filter out the smallest scales and resolve only the larger ones. The RANS
201 | equations (Smith & McLean 1977, Rödi 1988) do not resolve any scale but the stress terms used for
202 | their closure have proven useful for the modelling of near-bed turbulent patterns (see next subsection).
203 | The general trend is that improvements in efficiency of the algorithms have approximately kept pace
204 | with exponential improvements in computer power over the past 50 years (Moore 1965, Mavriplis
205 | 1998, Koomey et al. 2010) which tends to push the limitations of DNS and LES further away.

206

207 | 2.2.2 Erosion

208 | Several types of practical applications dictate the use of high-level formalisms in the description of
209 | particle detachment and transport, typically to handle explicit bed geometries and alterations, for
210 | example jet scours and regressive erosion (Stein et al. 1993, Bennett et al. 2000, Alonso et al. 2002),
211 | diverging sediment fluxes in canals (Belaud & Paquier 2001) or incipient motion conditions,
212 | calculated from grain size, shape and weight (Stevenson et al. 2002). The NS formalism is also needed
213 | to describe strong water-sediment, *i.e.* couplings in which the solid phase exerts an influence on the
214 | liquid phase, acting upon velocity fields, flow rheology and erosive properties (Sundaresan et al.
215 | 2003), ~~Parker & Coleman 1986, Parker et al. 1986, Davies et al. 1997, Mulder & Alexander 2001~~.
216 | Such couplings may be sorted by increasing sediment loads, from dispersed multiphase flows (Parker
217 | & Coleman 1986, Davies et al. 1997) to density currents (Parker et al. 1986), hyperconcentrated flows
218 | (Mulder & Alexander 2001) and up to debris flows (Bouchut et al. 2003, Bouchut & Westdickenberg
219 | 2004), the latter derived as mathematical generalisations of the well-known Savage & Hütter (1989,
220 | 1991) avalanche models over explicit, pronounced topographies. Moreover, the NS formalism offers
221 | the possibility to work on the energy equations: the erosive power and transport capacity of sediment-
222 | laden flows may be estimated from the energy of the flow, debating the case of turbulence damping
223 | (or not) with increasing sediment loads (Vanoni 1946, Hino 1963, Lyn et al. 1992, Mendoza & Zhou
224 | 1997). The matter is not free from doubt today (Kneller & Buckee 2001) and frictional drag, abrasion

225 due to impacts of the travelling particles and increased flow viscosity have been described prone to
226 enhance the detachment capacities of loaded flows (Alavian et al. 1992, Garcia & Parker 1993).

227

228 **2.3 Reynolds-Averaged Navier-Stokes**

229 *2.3.1 Water flow*

230 The Reynolds-Averaged Navier–Stokes (RANS) equations are a turbulence model, using time-
231 averaged equations of fluid motion, less generic than the NS formalism. The hypothesis behind these
232 equations is that instantaneous pressure and velocities may be decomposed into time-averaged and
233 randomly fluctuating turbulent parts, which finally yields:

$$\frac{\partial \bar{u}}{\partial t} + \bar{u} \frac{\partial \bar{u}}{\partial x} + \bar{w} \frac{\partial \bar{u}}{\partial z} + g \frac{\partial H}{\partial x} = gS + \frac{1}{\rho} \frac{\partial \tau}{\partial z} \quad (2)$$

234 where \bar{u} [LT^{-1}] and \bar{w} [LT^{-1}] are the time-averaged local water velocities in x and z, H is the flow
235 depth [L] and S is the bed slope [-].

236

237 In this formulation, the "Reynolds stress" term τ is of crucial importance for free-surface flow,
238 friction and erosion modelling, especially for shallow flows, first because it is the closure term (
239 $\tau = -\rho \overline{u'w'}$) and second because the Reynolds stresses have been closely related, in magnitude and
240 direction, to the size and arrangement of bed asperities. The combined analysis of the relative
241 magnitude of the u' and w' terms has become the purpose of "quadrant analysis" (Kline et al. 1967,
242 Raupach 1981, Kim et al. 1987) that identifies the four cases of outward interactions (quadrant I: $u'>0$,
243 $w'>0$), ejections (quadrant II: $u'<0$, $w'>0$), inward interactions (quadrant III: $u'<0$, $w'<0$) and sweeps
244 (quadrant IV: $u'>0$, $w'<0$). Depending on the submergence and geometry of bed asperities, the
245 maximal Reynolds stresses, those with significant effects on flow structure, have most often been
246 reported to occur near or just above the roughness crests (see Nikora et al. 2001, Pokrajac et al. 2007
247 and the review by Lamb et al. 2008a).

248

249 **2.3.2 Erosion**

250 In their paper on movable river beds, Engelund & Fredsoe (1976) judiciously reformulated and
251 exploited the existing hypotheses (Einstein & Banks 1950, Bagnold 1954, Fernandez Luque & van
252 Beek 1976) of a partition between “tractive” destabilizing shear stresses and “dispersive” equalizing
253 drags. The vertical concentration profiles of bedload and suspended load were calculated from
254 incipient sediment motion conditions, relating stresses on the particles to the values and variations of
255 near-bed velocities. One step further, the physical explanation, mathematical definition, point of
256 application, main direction and erosive efficiency of the turbulent near-bed stresses have become
257 private hunting grounds of the RANS models throughout the years (Nikora et al. 2001, Nino et al.
258 2003).

259

260 The maximal Reynolds stresses are located near the crests of the submerged bed asperities, where
261 turbulent velocity fluctuations reach several times the average near-bed velocity values, which greatly
262 enhances particle detachment (Raupach et al. 1991, Nikora & Goring 2000, Lamb et al. 2008a). Very
263 few studies deal with the magnitude and point of application of the Reynolds stresses for partial
264 inundation cases (Bayazit 1976, Dittrich & Koll 1997, Carollo et al. 2005) although turbulent flows
265 between emergent obstacles often occur in natural settings. Particle detachment is generally attributed
266 to “sweeps” (quadrant IV: $u' > 0, w' < 0$) (Sutherland 1967, Drake et al. 1988, Best 1992) or “outward
267 interactions” ($u' > 0, w' > 0$) (Nelson et al. 1995, Papanicolaou et al. 2001) but depends on bed
268 geometries and bed packing conditions. Finally, the RANS equations allow explicit calculations of
269 shear stresses and particle-scale pick-up forces, thus incipient motion conditions (Nino et al. 2003,
270 Afzalimehr et al. 2007). They may handle the movements of detached particles in weak transportation
271 stages (Bounvilay 2003, Julien & Bounvilay 2013) down to near-laminar regimes (Charru et al. 2004).

272 **2.4 Saint-Venant**

273 **2.4.1 Water flow**

274 The Saint-Venant (SV) equations are obtained by depth-integrating the Navier–Stokes equations,
275 neglecting thus the vertical velocities as well as vertical stratifications in the streamwise velocity
276 (Stoker 1958, Johnson 1998, Whitham 1999). The integration process (Chow 1959, Abbott 1979)
277 incorporates an explicit bottom friction term τ_0 that previously appeared only as a boundary condition
278 in the NS and RANS equation:

$$\frac{\partial U}{\partial t} + U \frac{\partial U}{\partial x} + g \frac{\partial H}{\partial x} = gS + \frac{\tau_0}{\rho H} \quad (3)$$

279
280 Recent attempts have been made in the field of fluid mechanics to derive specific expressions for τ_0
281 (laminar flows: Gerbeau & Perthame 2001, macro-roughness: Roche 2006, thin flows: Devauchelle et
282 al. 2007, turbulent flows: Marche 2007, multi-layer SV model: Audusse et al. 2008). However, the
283 common practice in hydraulics and hydrology is rather to approximate steady-state equilibrium
284 between bottom friction τ_0 and the streamwise stress exerted at the bottom of a water column
285 ($\tau_0 = \rho g H S_f$) to reach the popular formulation:

$$\frac{\partial U}{\partial t} + U \frac{\partial U}{\partial x} + g \frac{\partial H}{\partial x} = g(S - S_f) \quad (4)$$

(i) (ii) (iii) (iv) (v)

286 where (i) is the unsteadiness term, (ii) the convective acceleration term, (iii) the pressure gradient
287 term, while (iv) and (v) form the diffusive wave approximation (later discussed).

288
289 In the above, $S_f (-)$ is the “friction slope” whose expression depends on flow velocity and on the
290 chosen friction law, often one of the Chézy, Darcy-Weisbach or Manning formulations (e.g.
291 $S_f = nU^2/8gH$ with Manning’s n friction coefficient). The derivation of the SV equations by Boussinesq
292 (1877) involved a momentum correction coefficient $\beta [-]$ in the advection term (King & Brater 1963,

293 Chen 1992) to account for stratification effects in the vertical distribution of velocities, especially
294 plausible in sediment-laden flows or in presence of density currents.

295

296 The SV equations may account for flows of variable widths and depths, for example in floodplains
297 (Bates & De Roo 2000, Beltaos et al. 2012), rivers (Guinot & Cappelaere 2009), overland flow
298 (Berger & Stockstill 1995, Ghavasieh et al. 2006), overpressure in drainage systems (Henine et al.
299 2014), man-made channels (Zhou 1995, Sen & Garg 2002, Sau et al. 2010), vegetation flushing (Fovet
300 et al. 2013), channel networks (Choi & Molinas 1993, Camacho & Lees 1999) or natural settings
301 (Moussa & Bocquillon 1996a, Wang & Chen 2003, Roux & Dartus 2006, Burguete et al. 2008, Bates
302 et al. 2010), including these with curved boundaries (Sivakumaran & Yevjevich 1987). Discharge and
303 cross-sectional area may conveniently be used instead of velocity and water depth, and the two
304 equations describing mass and momentum in the Saint-Venant system now write (Sivapalan et al.
305 1997):

$$\frac{\partial A}{\partial t} + \frac{\partial Q}{\partial x} = q_a \quad (5)$$

$$\frac{1}{gA} \frac{\partial Q}{\partial t} + \frac{1}{gA} \frac{\partial}{\partial x} \left(\beta \frac{Q^2}{A} \right) + \frac{\partial H}{\partial x} + S_f - S = 0 \quad (6)$$

306 where A is the cross-sectional area [L^2], Q is the discharge [L^3T^{-1}], q_a is the lateral flow per unit
307 channel length [L^2T^{-1}]. The magnitudes of the various terms in equations (5) and (6) are given in the
308 literature (e.g. Henderson 1966, Kuchment 1972).

309

310 2.4.2 Erosion

311 In the hydrology-erosion community, the SV level is that of the *Concepts of mathematical*
312 *modelling of sediment yield* by Bennett (1974). This landmark paper extended Exner's (1925)
313 conservation of sediment mass, adding the possibility to handle different fluid and particle velocities,
314 also accounting for particle dispersion *via* a diffusion term. Unfortunately, most citing papers discard
315 this term, taking particle velocity equal to water velocity. The assumption seems false if transport

316 occurs as bedload or saltation load, questionable for suspended load trapped into turbulent motions,
317 exact only for very small particles borne by laminar flows. Although warning against the capability of
318 first-order laws to “*represent the response of sediment load to changes in transport and detachment*
319 *capacity*” (Bennett 1974, p.491), the author recommended the use of such a model (Foster and Meyer
320 1972). The proposed simplification writes $e/D_c=1-c/T_c$, where the net erosion rate (e) is normalised by
321 the maximal detachment capacity (D_c) while sediment load (c) is normalised by the maximal transport
322 capacity of the flow (T_c). An additional (uncertain) hypothesis was that of maximal detachment
323 capacity for minimal sediment load, *i.e.*, clear water. See the controversial comments around the
324 Wainwright et al. (2008) paper: the areas of disagreement revolve around the ability of models to
325 handle unsteady flow conditions, to deal with suspended and/or bedload transport, to consider particles
326 of different sizes and to stay valid over realistic ranges of sediment concentration.

327

328 Those questions directly address the possibilities of SV-level approaches. Higher-level models
329 (NS, RANS) better address the dynamics of incipient motion (Dey & Papanicolaou 2008), especially
330 in shallow laminar flows (Charpin & Myers 2005) or focusing on granular flows (Parker 1978a, b,
331 Charru et al. 2004, Charru 2006). Refined models are also needed to explicitly handle specific particle
332 velocities (Bounvilay 2003), to describe particle diffusion in secondary currents (Sharifi et al. 2009),
333 to account for the spatial heterogeneity of “neither laminar nor turbulent” overland flows (Lajeunesse
334 et al. 2010) or to introduce modifications in flow rheology (Sundaresan et al. 2003). On the other
335 hand, slope effects (Polyakov & Nearing 2003), particle-size effects (Van Rijn 1984a, Hairsine &
336 Rose 1992a, Sander et al. 2007, Wainwright et al. 2008), flow stratification effects (van Maren 2007),
337 the effects of hyperconcentrated flows (Hessel 2006) and the bedload transport (Van Rijn 1984b,
338 Julien & Simmons 1985, Hairsine & Rose 1992b, Wainwright et al. 2008) have received much
339 attention within the SV or ASV formalisms.

340

341 Whatever the liquid-solid coupling opted for, the SV level covers the widest variety of contexts,
342 from overland erosion models (Simpson & Castellort 2006, Nord & Esteves 2010) to dam-break

343 hydraulics over erodible beds (Cao et al. 2004) and the analysis of channel inception driven by the
344 variations of the Froude number (Izumi & Parker 1995) or the impact of travelling particles (Sklar &
345 Dietrich 2004, Lamb et al. 2008b). Sediment detachment and transport over plane beds (Williams
346 1970), rough beds (Afzalimehr & Ancil 1999, 2000, Gao & Abrahams 2004), step-pools (Lamarre &
347 Roy 2008) or pool-riffle sequences (Sear 1996, Rathburn & Wohl 2003) have yielded often-cited
348 studies, while sediment flushing in reservoirs (Campisano et al. 2004) and vegetation flushing in
349 canals (Fovet et al. 2013) constitute more specific applications. Cited limitations of the SV approaches
350 are their inability to explicitly describe the near-bed velocity fluctuations, especially the local
351 accelerations responsible for particle entrainment but also the vertical gradients of the streamwise
352 velocity, for bedload transport in the laminar layer. This lack of accuracy in the description of flow
353 characteristics also endangers the possibility to predict the formation, transformation and migration of
354 geometrical bed patterns, which in turn requires the full set of 3D (x, y, z) NS equations in several
355 cases (Lagrée 2003, Charru 2006, Devauchelle et al. 2010).

356
357 There seems to exist a dedicated "NS-SV Morphodynamics" research lead that uses rather simple
358 bedload transport formulae (Du Boys 1890, Meyer-Peter & Müller 1948, Einstein & Banks 1950,
359 Bagnold 1966, Yalin 1977) to calculate sediment fluxes from excess bed shear stresses, in studies of
360 long-term system evolutions. These low "system evolution velocities" appear under the "quasi-static"
361 flow hypothesis: particle velocity may be neglected before water velocity, which allows neglecting the
362 unsteadiness term in the momentum equation but on no account in the continuity equation (Exner law)
363 that describes bed modifications (Parker 1976). Moreover, shear stresses are generally calculated from
364 near-bed laminar or near-laminar velocity profiles, sometimes with the regularising hypothesis that
365 detachment and transport occur just above the criterion for incipient motion (see the review by
366 Lajeunesse et al 2010). Various applications address rivers with mobile bed and banks (Parker 1978a,
367 b), focus on self-channelling (Métivier & Meunier 2003, Mangeney et al. 2007) and often resort to
368 formulations at complexity levels between these of the NS and the SV approaches (Devauchelle et al.
369 2007, Lobkovsky et al. 2008).

370

371 **2.5 Approximations to Saint-Venant**

372 **2.5.1 Water flow**

373 When the full Saint-Venant equations are not needed or impossible to apply due to a lack of data,
374 an option is to neglect one or several terms of the momentum equation (Ponce and Simons 1977,
375 Romanowicz et al. 1988, Moussa & Bocquillon 1996a, Moussa & Bocquillon 2000). In most practical
376 applications for flood routing, the unsteadiness (i) and convective acceleration (ii) terms in (4) may be
377 neglected, suppressing the first two terms from (6). Combining the remaining terms in (5) and (6), we
378 obtain the Diffusive Wave equation (Moussa, 1996):

$$\frac{\partial Q}{\partial t} + C \left(\frac{\partial Q}{\partial x} - q_a \right) - D \left(\frac{\partial^2 Q}{\partial x^2} - \frac{\partial q_a}{\partial x} \right) = 0 \quad (7)$$

379 where $C [LT^{-1}]$ and $D [L^2T^{-1}]$ are non-linear functions of the discharge Q (and consequently the flow
380 depth H) known as the celerity and diffusivity, respectively.

381

382 In cases where the pressure-gradient term (iii) in (4) can also be neglected, the third term of (6)
383 also vanishes and the Diffusive Wave becomes the Kinematic Wave equation, with $D=0$ in (7). The
384 Diffusive Wave (Cunge 1969, Akan & Yen 1981, Rutschmann & Hager 1996, Wang et al. 2006,
385 Wang et al. 2014) can thus be considered a higher order approximation than the Kinematic Wave
386 approximation (Katopodes 1982, Zoppou & O'Neill 1982, Daluz Vieira 1983, Ferrick 1985, Ponce
387 1990). Both have proven very useful for canal control algorithms (Rodellar et al. 1993) or flood
388 routing procedures, with lateral inflow (Fan & Li 2006), in rectangular channels (Keskin &
389 Agiraliloglu 1997), for real time forecast (Todini & Bossi 1986), in lowland catchments (Tiemeyer et
390 al. 2007), for small catchments (Moussa et al. 2002, Chahinian et al. 2005, Charlier et al. 2007), for
391 mountainous catchments (Moussa et al. 2007) or tropical catchments (Charlier et al. 2009), at the
392 largest scale of the Amazon basin (Trigg et al. 2009, Paiva et al. 2013), for anthropogenic hillslopes
393 (Hallema & Moussa 2013), to address backwater effects (Munier et al. 2008), stormwater runoff on

394 impervious surfaces (Blandford & Meadows 1990, Parsons et al. 1997), stream-aquifer interactions
395 (Perkins & Koussis 1996) or volume and mass conservation issues (Perumal & Price 2013). Given
396 their "nominal" scales of application, the ASV models are sometimes fed by airborne (remote sensing)
397 data acquisition (Jain & Singh 2005, Reddy et al. 2007). In addition, predictive uncertainties (Elhanafy
398 et al. 2008) or the applicability of the kinematic and diffusive wave equations are the main scope of
399 several studies (Liggett & Woolhiser 1967, Ponce & Simons 1977, Ponce et al. 1978, Moussa &
400 Bocquillon 1996b, Bajracharya & Barry 1997), the evaluation of modelling strategies is that of Horritt
401 & Bates (2002), while parameter estimation is addressed, among others, by Koussis et al. (1978).

402

403 2.5.2 Erosion

404 Whereas common practices in fluid mechanics and hydraulics are rather to seek context-specific
405 strategies in erosion modelling, two simplifying and unifying trends, if not paradigms, have developed
406 in the field of hydrology. The first one is the transport capacity concept (Foster & Meyer 1972) in
407 which the erosive strength of the flow decreases with increasing suspended sediment load, until a
408 switch occurs from detachment- to transport-limited flows. The second one is the stream power
409 concept (Bagnold 1956) that *slope times discharge* is the explicative quantity for erosion, with
410 adaptations that mentioned unit stream power (*slope times velocity*, Yang 1974, Govers 1992) or fitted
411 exponents to the slope and discharge terms (Julien & Simmons 1985). Many catchment-scale
412 hydrology-erosion models (e.g. ANSWERS: Beasley et al. 1980, CREAMS: Knisel 1980, KINEROS:
413 Smith et al. 1995, LISEM: De Roo et al. 1996, WEPP: Ascough et al. 1997, EUROSEM: Morgan et
414 al. 1998, MAHLERAN: Wainwright et al. 2008, MHYDAS-Erosion: Gumiere et al. 2011) adopt the
415 1D Diffusive or Kinematic Wave Equations to route water fluxes, possibly through vegetated strips
416 (Muñoz-Carpena et al. 1999), together with the simplest possible couplings between water and
417 sediment fluxes (Aksoy & Kavvas 2005).

418

419 A known difficulty when embracing larger scales with simplified models is to describe the
420 spatially-distributed sources and sinks of sediments (Jetten et al. 1999, 2003) with or without explicit
421 descriptions of the permanent or temporary connectivity lines, for water and sediment movements
422 (Prosser & Rustomji 2000, Croke & Mockler 2001, Pickup & Marks 2001, Bracken et al. 2013). What
423 tends to force reduced complexity approaches in erosion models is the necessity to handle distinct
424 detachment, transport and deposition processes (from the very shallow diffuse flows formed during
425 runoff initiation to the regional-scale basin outlets) with only sparse data on flow structure and soil
426 characteristics (cohesion, distribution of particle sizes, bed packing). Parsons & Abrahams (1992)
427 have established how the agronomical, engineering and fluvial families of approaches have converged
428 into similar modelling techniques, especially on the subject of erosion in overland flows (Prosser &
429 Rustomji 2000). The ASV formalism also allows fitting bedload transport formulae against mean
430 discharge values as a surrogate to the overcomplicated explicit descriptions of erosion figures in high-
431 gradient streams with macro-roughness elements (Smart 1984, Aziz & Scott 1989, Weichert 2006,
432 Chiari 2008). ASV-level couplings have also been applied to study the slope independence of stream
433 velocity in eroding rills (Gimenez & Govers 2001) and the appearance of bed patterns in silt-laden
434 rivers (van Maren 2007).

435

436 **3 Determinants of modelling choices**

437 This section aims at the construction of a signature for each case study, relating the "conceptual"
438 choice of a model refinement (Navier-Stokes: NS, Reynolds-Averaged Navier-Stokes: RANS, Saint-
439 Venant: SV or Approximations to Saint-Venant ASV) to the "contextual" descriptors, i.e. the
440 spatiotemporal scales (section 3.1), spatiotemporal scales and flow typologies (section 3.2),
441 spatiotemporal scales, flow typologies and dimensionless numbers (section 3.3). Figures 2, 3, 5, 6 and
442 7 in this section were drawn from the 158 studies listed in Appendix A.

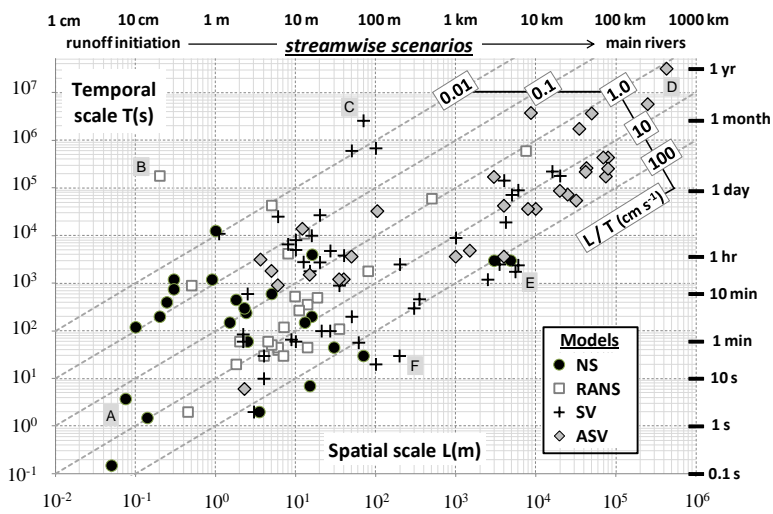
443 **3.1 Spatiotemporal scales**

444 **3.1.1 Influence of domain length (L) and time scale (T)**

445 A cross-disciplinary analysis of the cited literature indicates a clear correlation between the (L , T)
446 scales and the chosen model refinement (NS, RANS, SV or ASV). In this (L , T) plane, Fig. 2
447 quantifies the expected trend that sophisticated (NS, RANS) models are required to represent rapidly-
448 varying small-scale phenomena (lower left) while simplified approaches (ASV) pertain to increased
449 durations and spatial extensions (upper right). Typical scales of application may be identified for each
450 model refinement: NS ($10\text{ cm} < L < 100\text{ m}$, $10\text{ s} < T < 1\text{ hr}$), RANS ($1\text{ m} < L < 100\text{ m}$, $10\text{ s} < T < 1\text{ hr}$), SV
451 ($10\text{ m} < L < 20\text{ km}$, $1\text{ min} < T < 5\text{ days}$) and ASV ($10\text{ m} < L < 1000\text{ km}$, $30\text{ min} < T < 1\text{ yr}$). However, some
452 studies consider larger spatial or temporal scales, for example Charru et al. (2004) for overland
453 granular flows (RANS, $L \sim 20\text{ cm}$, $T \sim 2\text{ days}$) or Rathburn & Wohl (2003) for pool-riffle sequences
454 (SV, $L \sim 70\text{ m}$, $T \sim 30\text{ days}$). Nevertheless, the existence of overlap regions suggests that the (L , T)
455 spatiotemporal scales are not the only factor governing the choice of flow models.

456
457 The influence of flow typologies is discussed later in details but could the modelling choices be
458 dictated by the scientific background of the modeller? A striking example is that of the SV models,
459 responsible for the largest overlaps in Fig. 2. They may for example be used by physicists, as an
460 upgraded alternative to the NS equations, in the field of environmental fluid mechanics (for limited
461 scales). They may as well be convenient for soil scientists interested in high-resolution hydrology or
462 for civil engineers who may need to cope with flow unsteadiness to handle erosion issues or to allow
463 correct sizing of the man-made structures (for somewhat wider scales).

464



465
 466 **Figure 2 – How increasing (L, T) spatiotemporal scales of the flow domain tend to be associated with**
 467 **decreasing complexity in the choice of flow models, sorted here into four levels of refinement: Navier-**
 468 **Stokes (NS), Reynolds-Averaged Navier-Stokes (RANS), Saint-Venant (SV) or Approximations to Saint-**
 469 **Venant (ASV). A transverse analysis involves forming L/T ratios, searching for clues to model selection**
 470 **according to these "system evolution velocities" or governed by flow typologies that would exhibit specific**
 471 **L/T ratios. This figure was assembled from information available in the studies cited in Appendix A,**
 472 **selecting six textbook cases (sketches A to F, Table 1) for illustration.**

473
 474 Figure 2 bears another type of information than the trend to decreasing model refinement with
 475 increasing spatiotemporal scales. As the x-ordinate indicates the spatial scale L and the y-ordinate the
 476 time scale T, then the L/T ratio has dimensions of a velocity. However, this quantity should not be
 477 interpreted as a flow velocity. It rather indicates which of the temporal (long-term, low L/T ratio) or
 478 spatial (short-term, high L/T ratio) aspects are predominant in the study. Hence, the five dotted
 479 diagonals ($L/T=10^{-4}$, 10^{-3} , 10^{-2} , 0.1 and 1 m s^{-1}) establish the numerical link between the spatial and
 480 temporal scales of the cited experiments. They also show the dispersion with respect to the expected
 481 (say "natural") correlation between increasing L and T values. This dispersion contains a lot of
 482 information. Judging from the plotted literature, the lowest L/T ratios (e. g. 10^{-4} m s^{-1}) tend to indicate
 483 systems with low "evolution velocities", possibly associated with long-term changes or effects (high T

484 values, low L values) obtained from repeated phenomena, multiple cycles and progressive
485 modifications. By contrast, high L/T ratios (e.g. 1 m s^{-1}) rather refer to single-event situations, more
486 associated with quick modifications of flow patterns or bed morphologies.

487

488 If rules of thumb in problem dimensioning were to be drawn from Fig. 2, geomorphological
489 concerns (dune migration, basin sedimentation, long-term bed modifications) probably require
490 stretching up the temporal scale so that low "system evolution velocities" would fall beneath $L/T=10^{-2}$
491 m s^{-1} while event-based modelling (dam breaks, formative discharges, flash floods) should be able to
492 handle high "system evolution velocities" near or beyond $L/T=1 \text{ m s}^{-1}$. This "fixed-L, chosen-T"
493 description of system evolution and characteristic time scales also refers to Fig. 1 in which the choice
494 of T is somehow left at the modeller's discretion, as a degree of freedom: how different from T_0
495 should T be? These points are the subject of detailed investigations in the field of morphodynamics
496 (Paola et al. 1992, Howard 1994, Van Heijst et al. 2001, Allen 2008, Paola et al. 2009). Indicators of
497 "system evolution velocities" with units of a velocity but different definitions may for example be
498 found in Sheets et al. (2002), who took the channel depth (H) divided by the average deposition rate to
499 obtain a relevant, characteristic time scale (T). For the same purpose, Wang et al. (2011) took the
500 characteristic bed roughness (ϵ) instead of channel depth. The objective is often to discriminate what
501 Allen (2008) called the "reactive" (high L/T) and "buffer" (low L/T) systems. With or without erosion
502 issues, a reasonable hypothesis here seems that the dispersion in L/T ratios arises from the variety of
503 flow contexts, which may necessitate different modelling strategies. In other terms, it is deemed in this
504 study that this secondary trend, associated with flow typologies, is also a determinant in the choice of
505 the flow model.

506

507 To take a few examples and guide the reader through the arguments and the figures of this paper,
508 Table 1 gathers the information available for the six textbook cases outlined by sketches A to F in
509 Fig.2. The selected studies represent a wide variety of cases (drawing an approximate envelop of cases
510 in the L-T plane of Fig.2) followed in the forthcoming stages of the analysis and associated figures in

511 Section 3.1.2 (determinants of modelling choices in the L-H plane, Fig.3), Section 3.2 (determinants
 512 sought in flow typology, Fig.6a and 7a) and Section 3.3 (determinants sought in the values of
 513 dimensionless numbers attached to the flow).

514

Case	Context	Authors	Model refinement	Spatiotemporal scales					Flow typology [†]	Dimensionless numbers [‡]					
				L (m)	T (s)	H (m)	L/T (m s ⁻¹)	H/L [†] (-)		T*	Re	Fr	S (%)	Λ_z	θ
A	Film flow	Charpin & Myers (2005)	NS	0.075	3.75	0.003	0.02	0.04	O	5	300	0.11	10	8.0	-
B	Laminar dynamics	Charru et al. (2004)	RANS	0.2	$1.8 \cdot 10^3$	0.007	$1.1 \cdot 10^{-6}$	0.035	O	6428	50	0.02	<0.01	12.1	0.14
C	Pool-riffles	Rathburn & Wohl (2003)	SV	70	$2.6 \cdot 10^6$	0.47	$3.5 \cdot 10^{-3}$	$6.7 \cdot 10^{-3}$	B	$7.8 \cdot 10^3$	$7.1 \cdot 10^3$	0.69	1.1	5108	34.1
D	Amazon River	Trigg et al. (2009)	ASV	$4.3 \cdot 10^3$	$3.15 \cdot 10^8$	10	$1.4 \cdot 10^{-3}$	$2.3 \cdot 10^{-3}$	F	58.5	$8 \cdot 10^3$	0.05	<0.01	6600	-
E	Step-pools	Grant et al. (1990)	SV	5530	1755	0.87	3.15	$1.5 \cdot 10^{-4}$	Hg	1.0	$2.7 \cdot 10^3$	1.03	4.5	1.25	-
F	Step-pools	Chin (1999)	SV	197.25	30	0.50	6.58	0.025	Hg	1.21	$4.0 \cdot 10^3$	3.58	6.25	1.22	-

515

516

517

518

519

520

521

522

523

524

[†] See section 3.1.2 - H/L is the fineness ratio of the flow comparing flow depth (H) to the length of the flow domain (L)

[‡] See Section 3.2 - O: Overland, Hg: High-gradient, B: Bedforms, F: Fluvial

[§] See Section 3.3 - T*: dimensionless period, Re: Reynolds number, Fr: Froude number, S: slope, Λ_z inundation ratio, θ Shields number

Table 1 - Six textbook cases representing an approximate envelope of all the tested cases in the L-T plane of Fig.2, where L is the spatial scale (length of the flow domain) and T the temporal scale (duration of the process studied). Spatiotemporal scales are the determinants of modelling choices discussed in Section 3.1. The additional influence of flow typology and dimensionless numbers are discussed in Sections 3.2 and 3.3.

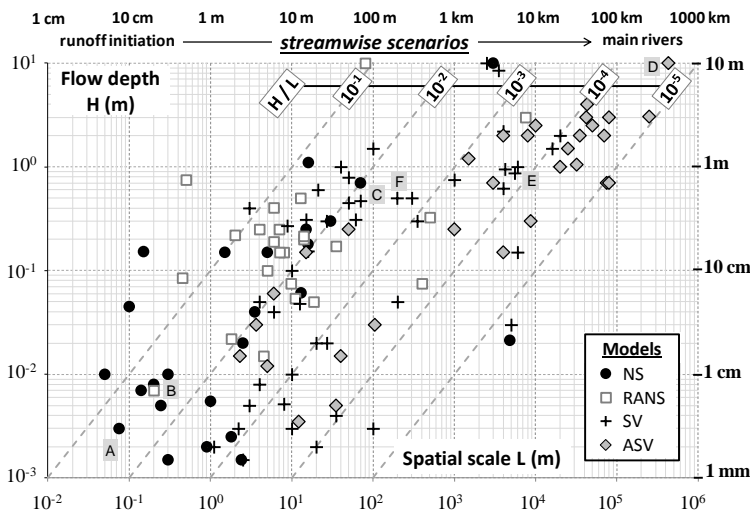
525

526 3.1.2 Influence of domain length (L) and flow depth (H)

527 The NS, RANS, SV and ASV equations are now positioned with respect to the spatial scale (L) and
 528 flow depth (H) of the reported experiments (Fig. 3), showing patterns and trends very similar to those
 529 of the (L, T) plane, though less pronounced. The global trend stays a decrease in refinement of the
 530 flow models from the smallest to the largest (L, H) values and typical scales of application may again
 531 be identified for each model refinement, NS (10 cm<L<100 m, 1 mm<H<30 cm), RANS
 532 (1 m<L<100 m, 5 cm<H<50 cm), SV (10 m<L<20 km, 1 cm<H<2 m) and ASV (10 m<L<1000 km,
 533 10 cm<H<10 m). Some studies provide outliers for example Gejadze & Copeland (2006) for canal
 534 control purposes (NS, L~3 km, H~10 m) or Cassan et al. (2012) for flows in lined channels (RANS,
 535 L~50 cm, H~75 cm). In an overview, wider overlaps and more dispersion occur in the (L, H) than in
 536 the (L, T) plane, especially for low to medium scales: flow depth (H) seems less discriminating than
 537 the time scale (T) in the choice of a flow model.

538

539 The transverse analysis of H/L "fineness ratios" (dotted diagonals $H/L=10^{-1}$, 10^{-2} , 10^{-3} , 10^{-4} and 10^{-5}) provides additional information, or rather a complementary reading grid on the information already
 540 plotted. First, only the NS and RANS models allow 2D (x, z) flow descriptions, which explains why
 541 these models have many of the largest H/L ratios (which, in most cases, stay within the $H \ll L$ shallow
 542 water hypothesis). Second, low H/L ratios provide justifications to discard 2D (x, z) descriptions at the
 543 benefit of 1D (x) descriptions within but also without the NS and RANS formalisms, so that the
 544 second diagonal of Fig. 3 (roughly from the upper right to the lower left) also shows a decrease in
 545 model refinement, towards SV and ASV points.
 546



547
 548 **Figure 3 – How increasing (L, H) spatiotemporal scales of the flow domain tend to be associated with**
 549 **decreasing complexity in the choice of flow models, sorted here into four levels of refinement: Navier-**
 550 **Stokes (NS), Reynolds-Averaged Navier-Stokes (RANS), Saint-Venant (SV) and Approximations to Saint-**
 551 **Venant (ASV). A transverse analysis involves forming H/L ratios, searching for clues to model selection**
 552 **according to the "fineness" of the flow or governed by flow typologies that would exhibit specific H/L**
 553 **ratios. This figure was assembled from information available in the studies cited in Appendix A, selecting**
 554 **six textbook cases (sketches A to F, Table 1) for illustration.**

555

556 **3.1.3 Influence of domain length (L), time scale (T) and flow depth (H)**

557 The links between model refinements (NS, RANS, SV or ASV) and spatiotemporal scales (L , T , H)
558 were shown in the (L , T) and (L , H) planes (Fig. 2 and 3). There was first the expected correlation
559 between increasing scales and decreasing model refinements. Then the transverse analyses involved
560 re-examining the same dataset from the values of the L/T and H/L ratios, also seeking the
561 determinants of modelling choices in the "system evolution velocity" (L/T) and "fineness" of the flow
562 (H/L).

563 - The values of the L/T ratios indicate that modelling choices owe much to the long-term (low L/T)
564 or short-term (high L/T) objectives associated with the target variables (velocity, discharge, particle
565 transport, bed modifications) thus influencing the choice of T values. However, this choice is not
566 totally free: it is likely constrained by flow characteristics and typologies.

567 - The values of the H/L ratios also indicate that flow typology (here, only its "fineness" is explicit)
568 may be a mattering determinant for the choice of a modelling strategy. This idea is explored in far
569 more details hereafter. The next section outlines the influence of friction, flow retardation and energy
570 dissipation processes on flow typology. It advocates thus the definition of flow typologies from
571 quantities related to the different types and/or magnitudes of flow retardation processes, provided
572 these quantities are easily accessible (e.g. bed geometry, water depth, bed slope, size of the roughness
573 elements).

574

575 **3.2 Flow typology**

576 **3.2.1 From friction laws and bed topography to flow characteristics**

577 Early insights on fluid friction and the definition of shear stress proportional to local velocity
578 gradients came together with the action-reaction law (Newton 1687): friction exerted on the flow was
579 of equal magnitude as the erosive drag, originally termed "critical tractive force" (Du Buat 1779) and
580 held responsible for particle detachment. The friction laws mostly resorted to in present-day modelling

581 do not often involve adaptations or generalisations of their famous empirical predecessors in civil
582 engineering (Chézy 1775, Weisbach 1845, Darcy 1857, Manning 1871) even if practitioners and
583 modellers are now confronted to far less controlled bed topographies and flow conditions, thus to a
584 wider variety of flow typologies. The theoretical derivation (or justification) of contextually relevant
585 friction laws seems therefore crucial, for water flow modelling at the microscopic (Richardson 1973,
586 Jansons 1988, Priezjev & Troian 2006) or macroscopic scales (Smith et al. 2007, Powell 2014), and
587 even more for erosion issues. In the literature, the modelling choices to account for friction
588 phenomena are most often correlated with the refinement of the flow models used (NS, RANS, SV,
589 ASV) but also constrained by bed topographies and flow typologies in numerous cases.

590

591 Several studies at the NS level of refinement advocate the use of the "partial slip" (Navier 1827)
592 condition or parented formulations in which the near-bed slip velocity is either proportional to the
593 shear stress (Jäger & Mikelic 2001, Basson & Gerard-Varet 2008) or depends on it in a non-linear way
594 (Achdou et al. 1998, Jäger & Mikelic 2003). Other works plead for "no-slip" conditions (Panton 1984,
595 Casado & Diaz 2003, Myers 2003, Bucur et al. 2008, 2010) or suggest the separation of flow domains
596 within or outside bed asperities, with a complete slip condition (non-zero tangential velocity) at the
597 interface (Gerard-Varet & Masmoudi 2010). A wider consensus exists at the RANS level, calculating
598 bottom friction as the local grain-scale values of the "Reynolds stresses" (Kline et al. 1967, Nezu &
599 Nekagawa 1993, Keshavarzy & Ball 1997), which has proven especially relevant for flows in small
600 streams over large asperities (Lawless & Robert 2001, Nikora et al. 2001, Pokrajac et al. 2007,
601 Schmeeckle et al. 2007). However, he who can do more, can do less, and it is still possible to use the
602 simplest empirical friction coefficients (Chézy, Manning) within sophisticated flow descriptions (NS:
603 Lane et al. 1994, RANS: Métivier & Meunier 2003). In the literature, the SV level of refinement is a
604 tilting point in complexity, that allows fundamental research, deriving ad hoc shear stress formulae
605 from the local fluid-solid interactions (Gerbeau & Perthame 2001, Roche 2006, Devauchelle et al.
606 2007, Marche 2007) or applied research, adjusting parameter values in existing expressions, for
607 specific contexts (e.g. boulder streams: Bathurst 1985, 2006, step-pool sequences: Zimmermann &

608 Church 2001, irrigation channels: Hauke 2002, gravel-bed channels: Ferro 2003). This trend holds for
609 most studies at the ASV level of refinement, though theoretical justifications of Manning's empirical
610 formula were recently derived (Gioia & Bombardelli 2002) and a recent mathematical study of the
611 diffusive wave equation (Alonso et al. 2008) introduces generalized friction laws for flows over non-
612 negligible topographic obstacles. The event-based variability of the friction coefficient in ASV models
613 has been investigated by Gaur & Mathur (2002).

614

615 If not decided from the level of refinement of the flow model, the friction coefficient (f) is chosen
616 in accordance with flow typology and bed topography, the former often described by the Reynolds
617 number (Re), the latter by the inundation ratio ($\Lambda_z = H/\epsilon$ where ϵ is the size of bed asperities, to which
618 flow depth H is compared). Such arguments were already present in the works of Keulegan (1938) and
619 Moody (1944) on flow retardation in open-channel and pipe flows, relating values of the friction
620 coefficient to the relative roughness ($\epsilon/H = 1/\Lambda_z$) of the flow, across several flow regimes (laminar,
621 transitional, turbulent) but only for small relative roughness (high inundation ratios). The existence of
622 implicit relations between f , Re and Λ_z has somehow triggered the search for contextual alternatives to
623 the sole f - Re relation for turbulent flows. Progressively lower inundation ratios were investigated
624 (Smith et al. 2007) until the real cases of emergent obstacles received attention (Bayazit 1976,
625 Abrahams & Parsons 1994, Bathurst 2006, Meile 2007, Mügler et al. 2010) including for non-
626 submerged vegetation (Prosser et al. 1995, Nepf 1999, Järvelä 2005, Nikora et al. 2008). For site-
627 specific friction laws, the default f - Re relation is sometimes complemented by f - Fr trends (Grant 1997,
628 Gimenez et al. 2004, Tatard et al. 2008) or f - Λ_z relations (Peyras et al. 1992, Chin 1999, Chartrand &
629 Whiting 2000, Church & Zimmermann 2007) in steep bed morphologies, where Fr is the Froude
630 number (Froude 1868).

631

632 Knowledge gained on flow retardation processes lead to the identification of key dimensionless
633 groups, to be included in any comprehensive analysis, formed from the "obvious", available elements
634 of bed geometry previously mentioned (Julien & Simons 1985, Lawrence 2000, Ferro 2003, Yager et

635 al. 2007). In numerous practical cases though, explicit bed geometries cannot be handled by the flow
636 models. A crucial surrogate becomes then to include as many geometrical effects as possible in the
637 chosen friction laws, for example these obtained from composite roughness experiments (Schlichting
638 1936, Colebrook & White 1937, Einstein & Banks 1950). A crucial advance was due to Smith &
639 McLean (1977) who attributed distinct retardation effects to bed particles, particle aggregates and
640 bedforms, corresponding to “grain spill”, “obstructions” and “long-wave form resistance” in the
641 subsequent literature. From then on, friction forces exerted by multiple roughness elements or scales
642 have often been described as additive-by-default, in shallow overland flows (Rauws 1980, Abrahams
643 et al. 1986), gravel-bed streams (Bathurst 1985, Lawless & Robert 2001, Ferro 2003), natural step-
644 pool formations (Chin & Wohl 2005, Canovaro & Solari 2007, Church & Zimmermann 2007) and
645 man-made spillways or weirs (Peyras et al. 1992, Chinnarasri & Wongwise 2006).

646

647 3.2.2 *From flow characteristics to flow typologies*

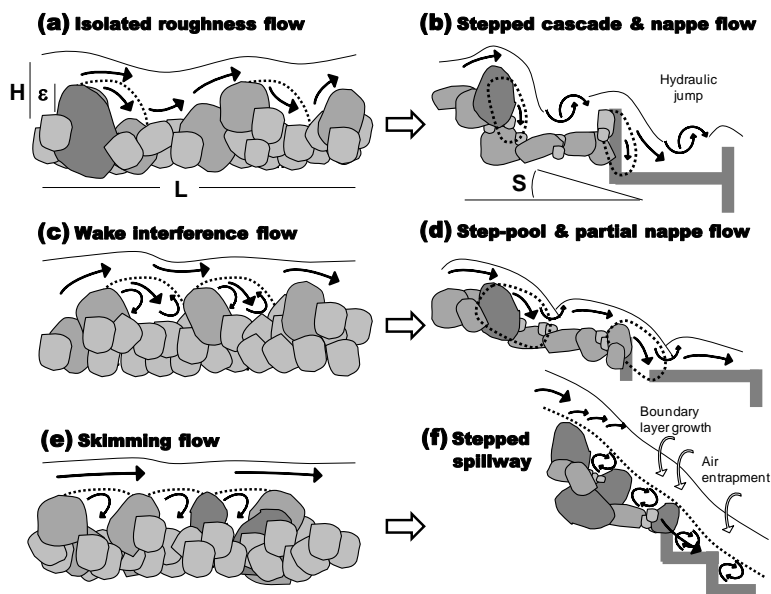
648 Several authors have put forward the existence of a scale-independent link between bed geometry,
649 flow retardation and flow structure, through the existence of three distinct flow regimes, from
650 geometrical arguments: "isolated roughness", "wake interference" and "skimming" flow (Morris 1955,
651 1959, Leopold et al. 1960, Fig. 4a, c and e). These flow descriptions were later applied in very
652 different contexts (Abrahams & Parsons 1994, Chanson 1994a, Papanicolaou et al. 2001,
653 Zimmermann & Church 2001), which suggests that analogies in energy dissipation and flow
654 retardation may exist across scales, from similar geometries and flow characteristics. This makes the
655 description somewhat generic, possibly used to constitute a set of flow typologies.

656

657 In Fig. 4a, the isolated roughness flow is laminar or weakly turbulent and the shade (streamline
658 diversion) of an obstacle does not reach the next. This setting ensures maximum energy dissipation,
659 which also holds for stepped cascades of natural or man-made nature in Fig. 4b: "nappe flows" loose
660 strength through energy-consuming fully-developed hydraulic jumps, isolated behind the major

661 obstacles (Peyras et al. 1992, Chanson 1994b, Wu & Rajaratnam 1996, 1998). In Fig. 4c the wake-
 662 interference flow is transitional or turbulent. The drag reduction and partial sheltering between
 663 obstacles depend on their spatial distribution and arrangements, as in Fig. 4d that shows "partial nappe
 664 flow" in relatively flat step-pool formations, with incomplete hydraulic jumps between obstacles of
 665 irregular sizes and spacing (Wu & Rajaratnam 1996, 1998, Chanson 2001). In Fig. 4e, the turbulent
 666 skimming flow exhibits a coherent stream cushioned by the recirculating fluid trapped between
 667 obstacles and responsible for friction losses. Similar characteristics appear in Fig. 4f, for submerged
 668 cascades or large discharges on stepped spillways. Air entrapment begins where the boundary layer
 669 reaches the free surface and flow aeration triggers subscale energy dissipation (Rajaratnam 1990,
 670 Chanson 1994b).

671



672

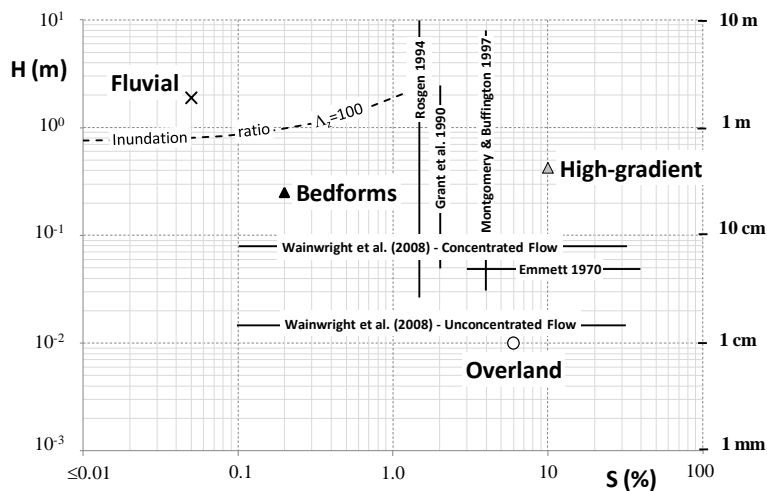
673 **Figure 4 – Analogies in flow characteristics, retardation processes and energy dissipation structures for**
 674 **very different flow typologies: streams (a, c, e) and high-gradient natural or man-made stepped flows (b,**
 675 **d, f). The combined values of flow depth (H), slope (S) and inundation ratio ($\Lambda_x=H/\epsilon$, where ϵ is the**
 676 **roughness size) appear as strong geometrical controls over flow characteristics and typologies.**

677

678

679 At this point, our set of flow typologies should be obtained from the geometrical arguments
 680 available in Fig. 4 (water depth H , bed slope S , inundation ratio $\Lambda_z=H/\epsilon$). The simplest way to proceed
 681 is to work in the (S, H) plane, then to add a criterion on Λ_z if the values of S and H are not
 682 discriminating enough. The first two flow typologies (Overland flow, noted O, and High-gradient
 683 flow, noted Hg) may be identified by a single criterion on H only ($H < H_{LIM}$, Emmett 1970, Wainwright
 684 et al. 2008) or on S only ($S > S_{LIM}$, Grant et al. 1990, Rosgen 1994, Montgomery & Buffington 1997).
 685 At least two flow typologies remained to be distinguished, Fluvial flows (F) and flows over significant
 686 bedforms (e.g. rough plane bed, dune-ripples or pool riffles, as suggested by Montgomery &
 687 Buffington 1997), referred to as Bedforms (B) in the following. Though Fluvial flows are expected to
 688 have the highest flow depths, an additional criterion on Λ_z may be used to make the difference
 689 between these last two typologies. Figure 5 positions the selected (O, Hg, B, F) flow typologies in the
 690 (S, H) plane.

691



692

693 **Figure 5 – Median position of the studies belonging to the "Overland", "High-gradient", "Bedforms" and**
 694 **"Fluvial" flow typologies, plotted on the (S, H) plane, also tracing an approximate**
 695 **additional criterion on the inundation ratio ($\Lambda_z=H/\epsilon$, where ϵ is the size of the bed asperities) to separate**

696 **the Fluvial and Bedforms types of flow. This figure was assembled from information available in the**
697 **studies cited in Appendix A.**

698

699 Moreover, there is a strong link between Fig. 4 and 5, which tends to ensure the genericity (if not
700 uniqueness) of the selected set of typologies. The Overland typology corresponds to Fig. 4a or c, the
701 Bedforms typology likely appears in Fig. 4c, the Fluvial typology in Fig. 4 and the High-gradient
702 typology in Fig. 4b, d or f. In coherence with Fig. 5, an increase in bed slope changes the Bedforms
703 and Fluvial typologies into the High-gradient typology, while an increase in both water depth and bed
704 slope is needed to do the same from the Overland typology.

705

706 *3.2.3 Influence of flow typologies on modelling choices*

707 Figures 6 and 7 provide a comprehensive picture of the most used associations between models
708 (NS, RANS, SV or ASV), scales (L, T, H) and flow typologies (O, Hg, B or F) just added to the
709 analysis. These figures seem to indicate preferential [NS, O], [RANS, B] and [SV, Hg] associations, in
710 addition to the obvious [ASV, F] pair. The (L, H) plot of Fig. 6 seems more discriminating than the (L,
711 T) plot of Fig. 7 though identical trends appear.

712

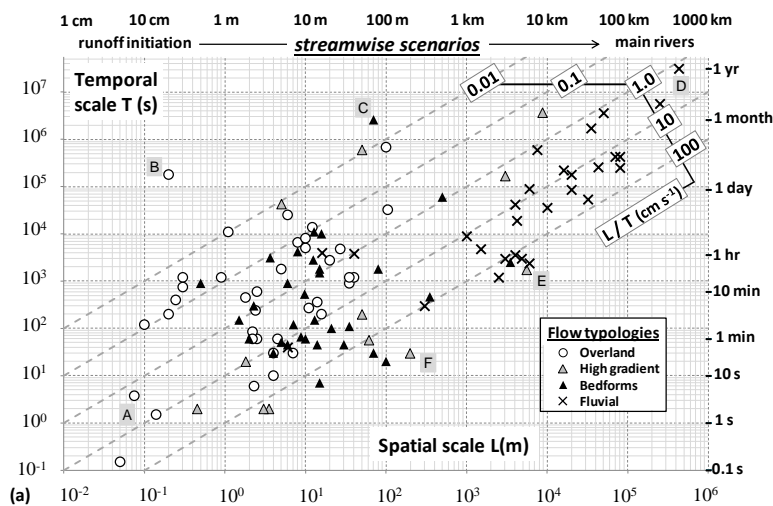
713 The [NS, O] association arises from the fact that several Overland studies involve very shallow
714 laminar flows and low sediment transport rates, best handled by adapted formulations of the NS
715 equations (nearly at the SV level), made suitable for low "system evolution velocities" ($L/T \approx 0.01 \text{ m s}^{-1}$,
716 Fig. 6). At somewhat larger spatial scales, the widely-used and multipurpose SV model has rather
717 low median $L/T \approx 0.02 \text{ m s}^{-1}$ values, mainly because many of its applications concern laminar flow
718 modelling and granular transport, as an alternative to the NS system or in formulations at complexity
719 levels intermediate between the NS and SV descriptions. These are clues that the [SV, O] association
720 may also be of special interest, despite the closest median positions of the NS and O points in the (L,
721 T) and (L, H) plots.

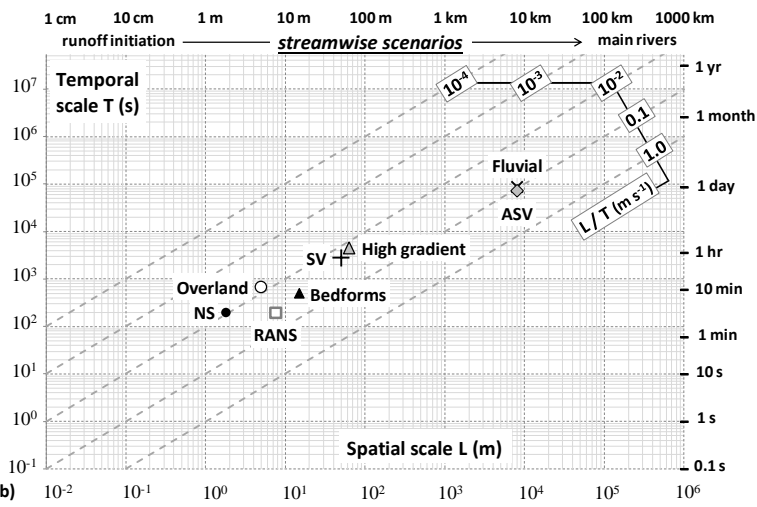
722

723 The RANS model (median $L/T \approx 0.07 \text{ m s}^{-1}$) and the ASV models (median $L/T \approx 0.1 \text{ m s}^{-1}$) tend to
724 involve higher "system evolution velocities". The former typically targets the description of numerous
725 short-term, high-frequency events (quadrant analysis for fluctuations in near-bed velocity, particle
726 pick-up by turbulent bursts). The latter is often associated with Fluvial flows: low H/L ratios with high
727 enough H and Λ_z values with weak friction, often resulting in very turbulent, high-velocity flow.
728 Moreover, studies handling erosion issues within the ASV formalism often hypothesize particle
729 transport to occur as suspended load only, equating particle and flow velocities, thus typically not
730 extending the time scale of the study to address the long-term, low velocity bedload transport involved
731 in morphodynamics, for example.

Mis en forme : Retrait : Première ligne : 0.5 cm

732

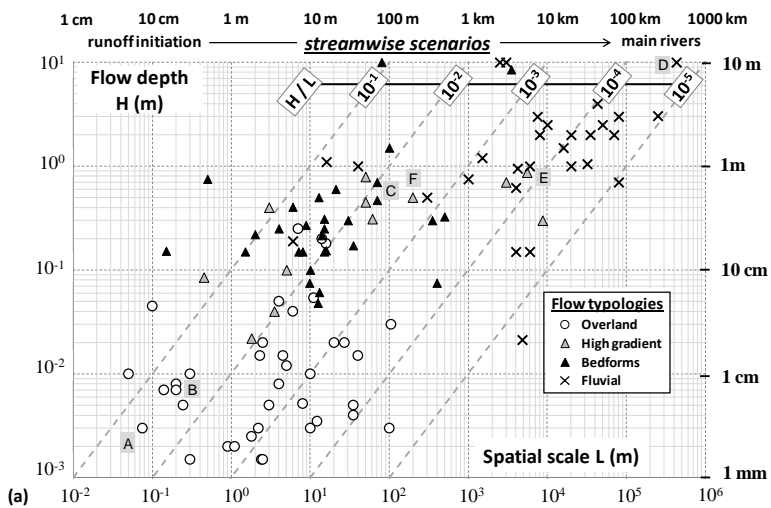




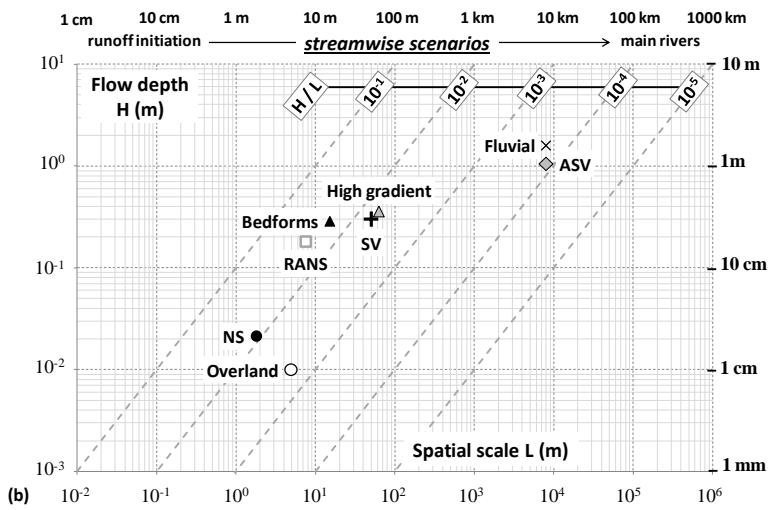
733 (b)

734 Figure 6 – Position of the flow typologies in the (L, T) plane for the studies listed in Appendix A, selecting
 735 six textbook cases (sketches A to F, Table 1) for illustration (a). Median positions for the choice of free-
 736 surface flow models (Navier-Stokes: NS, Reynolds-Averaged Navier-Stokes: RANS, Saint-Venant: SV or
 737 Approximations to Saint-Venant: ASV) and the study of flow typologies (Overland, High-gradient,
 738 Bedforms or Fluvial) across scales in the (L, T) plane (b). A transverse analysis involves forming L/T
 739 ratios, searching for clues to model selection according to these "system evolution velocities" or governed
 740 by flow typologies that would exhibit specific L/T ratios.

741



742 (a)



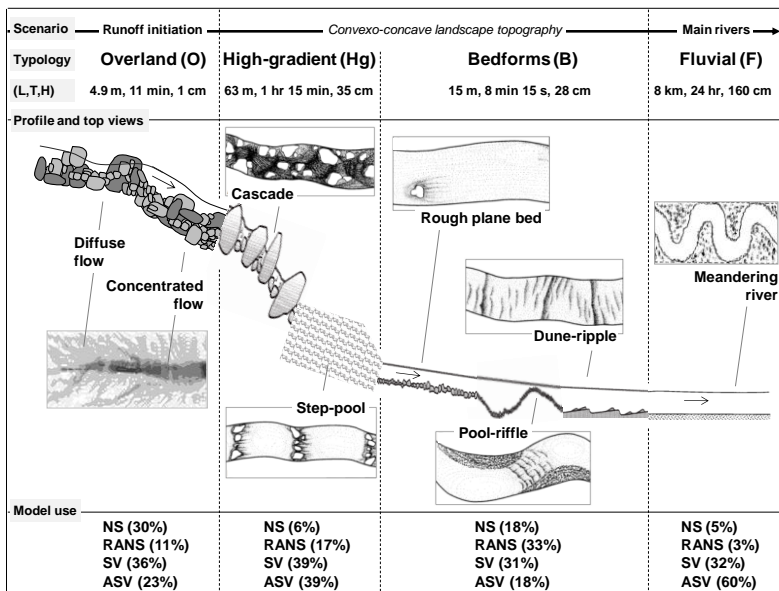
743 (b) 10^{-2} 10^{-1} 10^0 10^1 10^2 10^3 10^4 10^5 10^6

744 Figure 7 – Position of the flow typologies in the (L, H) plane for the studies listed in Appendix A, selecting
 745 six textbook cases (sketches A to F, Table 1) for illustration (a). Median positions for the choice of free-
 746 surface flow models (Navier-Stokes: NS, Reynolds-Averaged Navier-Stokes: RANS, Saint-Venant: SV or
 747 Approximations to Saint-Venant: ASV) and the study of flow typologies (Overland, High-gradient,
 748 Bedforms or Fluvial) across scales in the (L, H) plane (b). A transverse analysis involves forming H/L
 749 ratios, searching for clues to model selection according to these "finenesses" of the flow domain or
 750 governed by flow typologies that would exhibit specific H/L ratios.

751
 752

753 Several principles of organization between flow typologies may be inferred from reference studies
 754 (Grant et al. 1990, Montgomery & Buffington 1997, Church 2002) that discuss their succession in
 755 space (along longitudinal profiles) but also in time (which flow typologies are "experienced" by the
 756 flowing water during its course and which are the associated time scales). Plausible "streamwise
 757 scenarios" may therefore be assembled (Fig. 8), routing flow aggregations across increasing
 758 spatiotemporal scales and through several flow typologies, from the narrow-scale upland flows (runoff
 759 initiation) to the regional scales of the main rivers.

760



761

762 **Figure 8 – Streamwise scenario for a convexo-concave landscape topography, from runoff initiation to the**
 763 **main rivers, across flow typologies (Overland O, High-gradient Hg, Bedforms B or Fluvial F) and**
 764 **spatiotemporal scales (L, T, H). The indicated L, T and H values are the median values for the spatial**
 765 **scale, time scale and water depth, respectively, from the literature cited in Appendix A (Fig. 6 and 7). All**
 766 **sketches and drawings for the High-gradient and Bedforms typologies were taken from Montgomery &**
 767 **Buffington (1997). The top view for Overland flow is from Tatard et al. (2008) and that of a meandering**
 768 **river from Rosgen (1994). The "model use" panel indicates the model refinement most used (Navier-**
 769 **Stokes NS, Reynolds-Averaged Navier-Stokes RANS, Saint-Venant SV or Approximations to Saint-**
 770 **Venant ASV) to describe a given flow typology in the literature.**

771

772 **3.3 Dimensionless numbers**

773 **3.3.1 Contextual dimensionless numbers**

774 An angle of attack for the establishment of modelling strategies is provided by dimensional
775 analysis, to delineate the domains of validity of the selected flow models (NS, RANS, SV or ASV),
776 across their multiple spatiotemporal scales of application but in a powerful scale-independent analysis.
777 Justifications for the use of dimensionless numbers may be sought in the developments of similitude
778 laws (Fourier 1822, Rayleigh 1877, Bertrand 1878, Vaschy 1892, Riabouchinsky 1911), later extended
779 to dimensional analysis, providing guidance for the sizing of experimental facilities used in reduced-
780 scale modelling as well as more general arguments for the choice of adequate sets of dimensionless
781 quantities (Buckingham's 1914 π -theorem, Bridgman 1922, Langhaar 1951, Bridgman 1963,
782 Barenblatt 1987). Throughout history, the establishment of dimensionless numbers has led to the
783 recognition of contextually dominant terms in the flow equations, rendering them prone to dedicated
784 simplifications, provided these would not be used outside their conditions of validity, following
785 successive hypotheses made during their derivation.

786

787 From a wide overview of free-surface flow and erosion studies, a few dimensionless numbers stood
788 out and will be used in the procedure presented in the following. Some have already been mentioned
789 (Reynolds number Re , Froude number Fr) and some others have even been used to define flow
790 typologies (bed slope S , inundation ratio Λ_z). As all dimensionless numbers aim to describe flow
791 typology, the introduction of two more dimensionless numbers may be seen as an attempt to re-
792 examine the influence of flow typologies on modelling choices, from a different, more complete
793 perspective (especially if the dimensionless numbers not used in the definition of flow typologies
794 prove discriminating for the modelling choices).

795 - The dimensionless period $T^*=T/T_0$ handles temporal aspects by comparing the chosen time scale
796 (T) to the natural time scale (T_0) of the system, the latter obtained from the spatial scale of the system

797 and the average flow velocity as $T_0=L/U$ (Fig. 1). This dimensionless group or equivalent formulations
798 are used to model wave celerity in flood propagation issues (Ponce & Simons 1977, Moussa &
799 Bocquillon 1996a, Julien 2010) or to quantify the long characteristic times ($T^*\gg 1$) of basin-scale
800 sedimentation. In the latter, particle transport (and significant bed modifications) typically involve
801 lower velocities (and larger time scales) than these of water flow (Paola et al. 1992, Howard 1994,
802 Van Heijst et al. 2001) and the chosen T value witnesses this discrepancy.

803 - The Reynolds number $Re=UH/\nu$ compares flow inertia (velocity U times depth H) with the
804 adverse action of (kinematic) viscosity (ν [$L T^{-2}$]). In natural setting, over very rough boundaries, fully
805 turbulent flows are often reported for $Re>2000$, while the onset of turbulence within transitional
806 regimes occurs at $Re\sim 500$. Laminar overland flows, especially thin film flows, may have Re values as
807 low as $Re<100$.

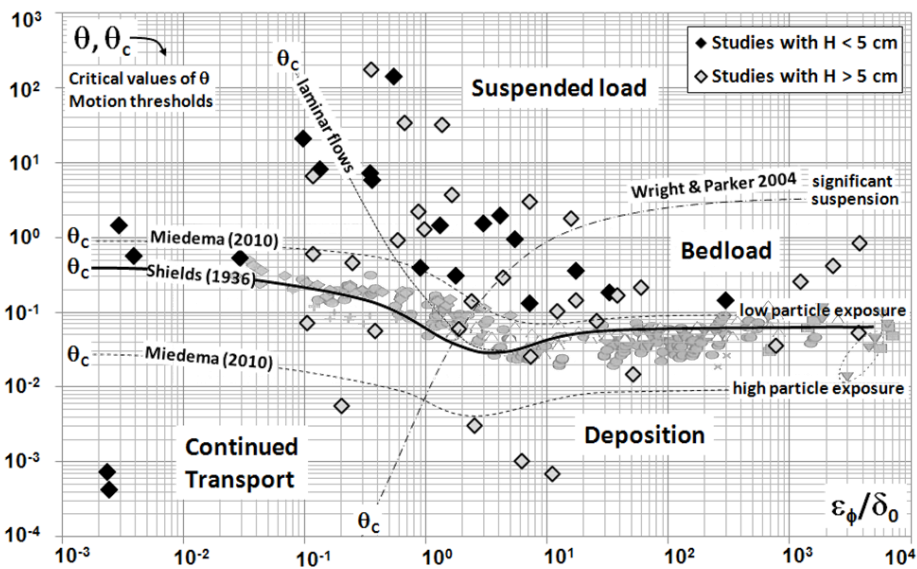
808 - The Froude number $Fr=U/(gH)^{0.5}$ denotes the influence of gravity (g) on fluid motion.
809 Supercritical $Fr>1$ values indicate torrential flows, accelerated by pressure effects, in which waves
810 propagate only downstream, also compatible with the appearance of localised energy dissipation
811 patterns (white waters, hydraulic jumps). Subcritical $Fr<1$ values indicate tranquil flows with
812 downstream controls.

813 - Topographical effects on flow phenomenology are almost always explicitly accounted for through
814 the average bed slope S , typically ranging from nearly zero ($S<0.01\%$) for large rivers to extremely
815 high values ($S\approx 100\%$) for gabion weirs, chutes or very steep cascades.

816 - Topography also appears through the inundation ratio $\Lambda_z=H/\varepsilon$ which allows a direct, model-
817 independent analysis of friction phenomena (Lawrence 1997, 2000, Ferguson 2007, Smith et al. 2007)
818 possibly dealing with large-size obstacles and form-induced stresses (Kramer & Papanicolaou 2005,
819 Manes et al. 2007, Cooper et al. 2013). The encountered values of Λ_z are very high for rivers flowing
820 on smooth, cohesive, fine-grained beds ($\Lambda_z>100$) and very low for all types of flows between emergent
821 obstacles ($\Lambda_z<1$).

822 - The dimensionless Shields number $\theta = \tau_0 / g \varepsilon_p (\rho_p - \rho)$ compares the drag force exerted on bed
 823 particles to their immersed weight, where ε_p and ρ_p account for the size and density of erodible
 824 particles.. The ratio between the current θ and the critical θ_c values indicates local flow conditions of
 825 deposition ($\theta < \theta_c$), incipient motion ($\theta \approx \theta_c$), transportation as bedload ($\theta > \theta_c$) or into suspension ($\theta \gg \theta_c$)
 826 (Shields 1936). This number seems appropriate for most erosion issues because it has been widely
 827 applied and debated in the literature (Coleman 1967, Ikeda 1982, Wiberg & Smith 1987, Zanke 2003,
 828 Lamb et al. 2008) and also because of its numerous possible adaptations (Neill 1968, ~~Parker et al.~~
 829 ~~2003~~, Ouriemi et al. 2007, Miedema 2010) to various flow typologies and non-uniform or poorly-
 830 known bed conditions. An impressive review on the use of the Shields number to determine incipient
 831 motion conditions, over eight decades of experimental studies, may be found in Buffington &
 832 Montgomery (1997). Finally, Fig.9 provides a generalized Shields diagram that includes motion
 833 threshold criteria under the effects of high or low particle exposure (Miedema 2010) or for laminar
 834 flows, also indicating the conditions of significant suspension (Wright & Parker 2004).

835



836

Mis en forme : Couleur de police : Automatique

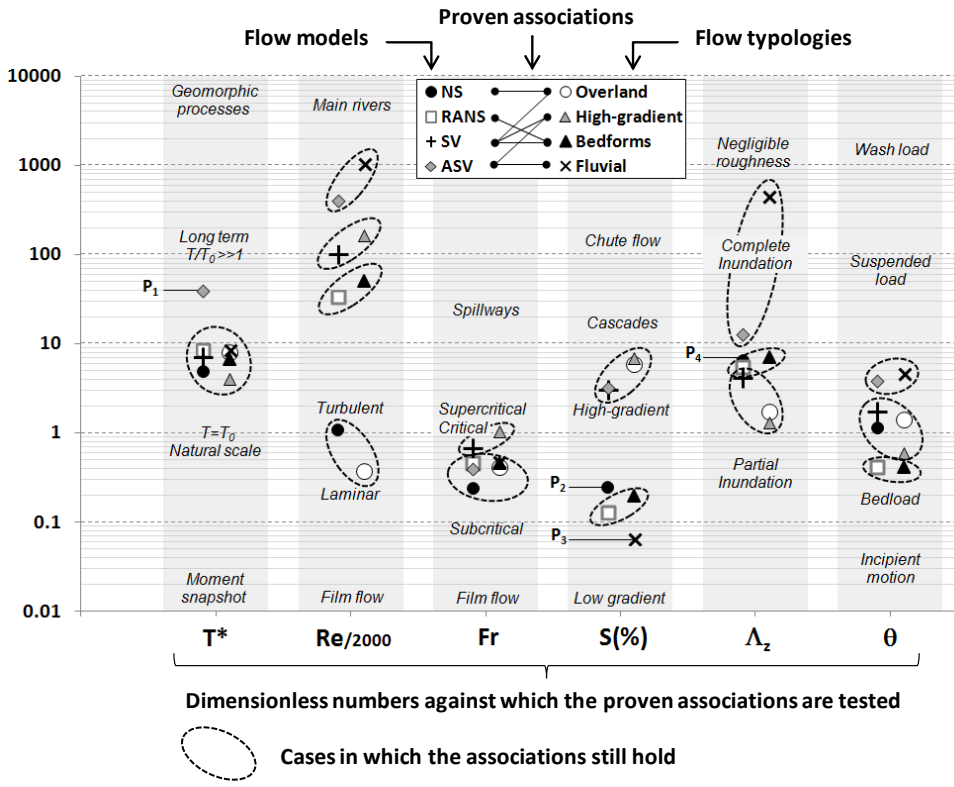
837 **Figure 9 - Generalized dimensionless Shields diagram that summarizes the conditions and regimes of**
838 **sediment transport or deposition, from the relative values of the Shields parameter (θ) and incipient**
839 **motion criterion (θ_c). The X-axis bears the values of the ratio of particle size (ϵ_p) on the depth of the**
840 **laminar sublayer (δ_0). The diamonds refer to the studies cited in Appendix A that deal with erosion issues:**
841 **black diamonds for studies in which flow depth is $H < 5$ cm, grey diamonds otherwise. Data in the**
842 **background show the critical θ_c values reported in the wide Buffington & Montgomery (1993) review of**
843 **incipient motion conditions for varied flow regimes, particle forms and exposures.**
844

845 3.3.2 Influence of the dimensionless numbers

846 As the purpose here is to re-examine the influence of flow typologies from the angle of the
847 dimensionless numbers, the chosen representation (Fig. 109) discards the (L, T, H) spatiotemporal
848 scales. It first recalls the preferential associations between models and flow typologies (see the "model
849 use" panel of Fig. 8) by tracing connecting dotted lines between flow typologies and the models most
850 used to handle them, in the legend of Fig. 109. It then examines whether these associations still hold,
851 for each of the six dimensionless numbers, by plotting and comparing the median values of T^* , Re, Fr,
852 S, Λ_z and θ for model uses (NS, RANS, SV or ASV) and flow typologies (O, Hg, B, F). The dotted
853 ellipses are "confirmations" (e.g. no additional information may likely be obtained from Re, Fr and θ).
854 Conversely, the presence of "non-associated" points (P_1 for T^* , P_2 and P_3 for S, P_4 for Λ_z) signals
855 something new: an influence not yet accounted for.
856

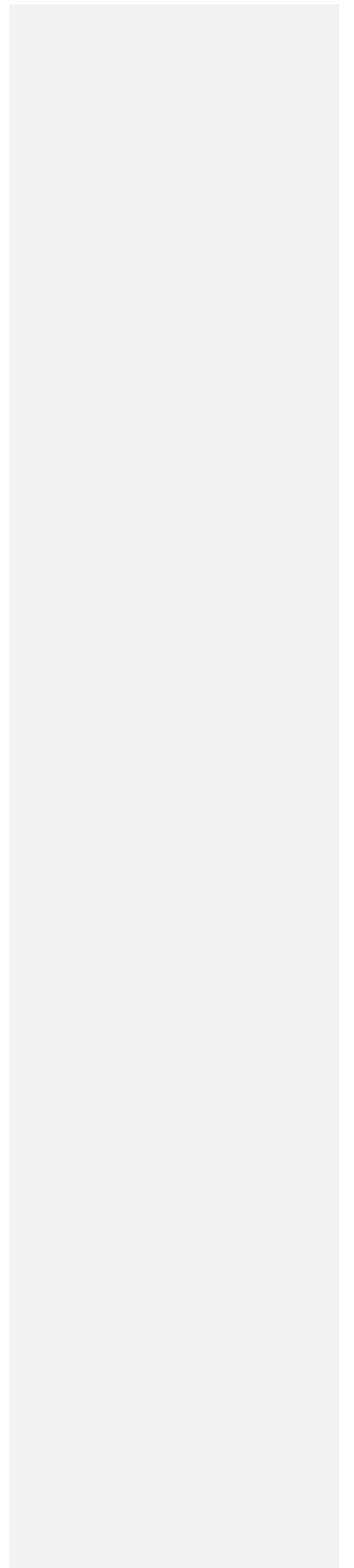
857 For example, the isolated P_1 point indicates the expected ASV-F association does not appear on the
858 T^* values, as the ASV applications exhibit higher median T^* values than the F typologies. The
859 suggested interpretation is that large (L, T, H) scales and Fluvial flows likely trigger the use of the
860 ASV model, though the necessity to handle large dimensionless periods makes the typological
861 argument less conclusive. The P_2 and P_3 points indicate the break of the NS-O and ASV-F associations
862 when examined from the angle of the bed slopes. This reinforces the use of bed slopes in the search for
863 determinants of modelling choices, either in the definition of flow typologies in the (S, H) plane or as

864 such. The P_4 point indicates the break of the NS-O association when considering the values of the
 865 inundation ratio, with the same conclusion as above.



866
 867 **Figure 109** - Comparative overview of the median values of the six selected dimensionless numbers
 868 (dimensionless period $T^*=T/T_0$, ratio of the chosen time scale on the "natural" time scale of the flow,
 869 Reynolds number Re , Froude number Fr , slope S , inundation ratio Λ_z and Shields parameter θ) obtained
 870 for the use of systems of equations (NS, RANS, SV and ASV) and the description of flow typologies (O,
 871 Hg, B and F) in the cited literature. The expected associations are indicated by dotted connecting lines in
 872 the legend box. The confirmed associations are indicated by dotted ellipses. Broken associations (isolated
 873 points P_1 to P_4) are discussed in the text. The typical and extreme ranges of the mentioned dimensionless
 874 numbers have been added for indication. This figure was assembled from information available in the
 875 studies cited in Appendix A.

876
 877



879 4 Conclusion

880 4.1 Outcomes of this review

881 In a free opinion on the use of models in hydrology, De Marsily (1994) elegantly argued that the
882 modelling of observable phenomena should obey “*serious working constraints, well-known from*
883 *classical tragedy: unity of place, unity of time, unity of action*”. This review paper investigates how
884 known spatial scales, temporal scales and flow typologies constrain the choice of a modelling strategy.
885 A normative procedure was built to facilitate the search for determinants of the modelling choices in
886 the cited literature.

887 - Each free surface flow model was placed in one of the NS, RANS, SV or ASV categories, whose
888 decreasing levels of refinement account for "Navier-Stokes", "Reynolds-Averaged Navier-Stokes",
889 "Saint-Venant" or "Approximations to Saint-Venant" types of approaches.

890 - The explored (L, T, H) spatiotemporal scales cover multiple orders of magnitude in the
891 streamwise direction ($5\text{ cm} < L < 1000\text{ km}$), the time duration ($0.1\text{ s} < T < 1\text{ yr}$) and flow depth (1
892 $\text{mm} < H < 10\text{ m}$).

893 - This study also encompasses a wide variety of free-surface flows, reduced to four typologies from
894 arguments on bed geometry, friction, flow retardation and energy dissipation processes. These
895 typologies are Overland flow (O: diffuse or concentrated), High-gradient flow (Hg: cascades, step-
896 pools), flows over significant Bedforms (B: rough plane beds, dune ripples, pool riffles) and Fluvial
897 flows (F: rivers, canals). Overland flows have the shallowest depths, High-gradient flows the highest
898 bed slopes, Fluvial flows have high flow depths and negligible bed roughness while Bedforms flows
899 may have any flow depth, over pronounced, non-negligible bedforms.

900 - In addition to the spatiotemporal scales and flow typologies, the determinants of modelling
901 choices are also sought in a series of six popular dimensionless numbers: the dimensionless period
902 (T^*), Reynolds and Froude numbers (Re , Fr), the bed slope (S), the inundation ratio ($\Lambda_z = H/\varepsilon$ where ε
903 is the size of bed asperities) and the Shields number (θ) that compares drag forces to particle weight.

904

905 In summary, each case-study may be defined by its signature, comprised of the chosen model (NS,
906 RANS, SV or ASV), the given spatiotemporal scales (L, T, H), flow typology (O, H, B or F) and
907 dimensionless numbers (T^* , Re, Fr, S, Λ_z , θ). Though non-unique, this signature is a generic and
908 normative classification of studies interested in free-surface flow modelling, with or without erosion
909 issues.

910 - The present review first illustrated the expected dominant trend of decreasing model refinement
911 with increasing (L, T, H) spatiotemporal scales. It appeared then that model uses could also be sorted
912 by their L/T and H/L ratios, though less clearly, which nevertheless provided indications that the
913 spatiotemporal scales were not the only determinant of modelling choices. This result suggested that
914 flow typologies (reduced here to the L/T "system evolution velocity" and H/L "fineness of the flow")
915 were also influential factors.

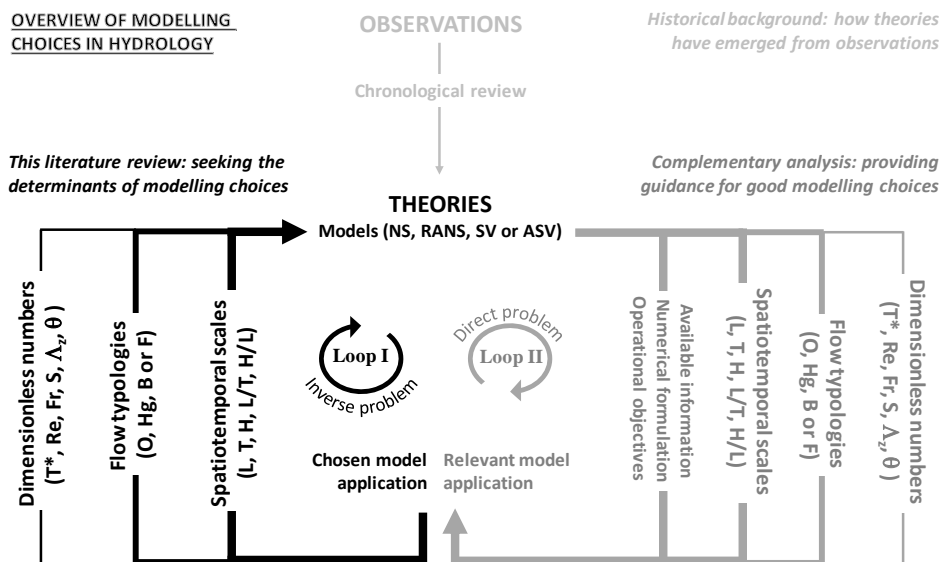
916 - A more exhaustive set of flow typologies was then derived from simple geometrical arguments,
917 combining criteria on S, H and Λ_z , represented in the (S, H) plane. This allowed quantifying the
918 median scales associated with studies interested in the Overland (O), Bedforms (B), High-gradient
919 (Hg) and Fluvial (F) typologies, sorted here by increasing spatiotemporal scales. Then came the
920 identification of preferential associations between flow models, scales and typologies: [NS, O] or [SV,
921 O], [RANS, B] or [SV, B], [SV, Hg] or [ASV, Hg] and [ASV, F] for increasing spatiotemporal scales.

922 - The final step was to re-examine the previous associations from the values of the dimensionless
923 numbers, thought here as more detailed, scale-independent descriptors of flow typologies. Several
924 associations were confirmed by the median values of the associated dimensionless numbers but the T^*
925 (dimensionless period), S (bed slope) and Λ_z (inundation ratio) introduced additional information., i.e.
926 correcting trends.

927
928 All arguments prevailing in the identification and sorting of flow models, scales, typologies and
929 dimensionless numbers may easily be debated and adapted, within the hydrology-erosion community
930 or for other research purposes. For example, multiple flow models, scales, typologies and
931 dimensionless numbers also intervene in the fields of pesticide fate modelling and groundwater

932 contamination issues, so the same procedure could be applied. Finally, this procedure offers the
 933 possibility to enrich the database of signatures if each modeller records his (or her) conceptual choices
 934 (flow models) in the proposed reading grid, together with the contextual elements (scales, typologies,
 935 dimensionless numbers) handled, for present and past studies. This would first help forming a
 936 comprehensive view of modelling choices, thus seeking guidance from "what has been done in similar
 937 cases", which however does not provide any critical analysis. Complementary investigations could
 938 certainly address the question of "what should be done", this time deciding the "model" part of the
 939 signatures from recommendations based on the scales, typologies and dimensionless numbers, as well
 940 as from additional elements, typically the modelling objectives (Figure 11).

941



942

943

944 **Figure 11 – This figure provides a simplified overview of the available modelling choices in hydrology, in three**
 945 **distinct colours associated with specific research purposes or disciplines, showing the position of the present review**
 946 **relative to the others. The pale grey section aims at understanding how the available flow models have emerged from**
 947 **observations and early formulations of the flow equations, focusing on their conditions of validity i.e. the successive**
 948 **hypotheses made during their derivation. The black section recalls the procedure followed in this review paper (Loop**
 949 **I, "inverse problem"). Literature sources are processed through a procedure that analyses how the spatiotemporal**

950 scales (spatial scale L , time scale T , flow depth H , L/T and H/L ratios), then flow typology (Overland O , High-gradient
951 Hg , Bedforms B or Fluvial F) and dimensionless numbers (dimensionless period T^* , Reynolds number Re , Froude
952 number Fr , bed slope S , inundation ratio A_2 , Shields parameter θ) determine the choice of a flow model (Navier-
953 Stokes NS , Reynolds-Averaged Navier-Stokes $RANS$, Saint-Venant SV or approximations ASV). Suggested in
954 medium grey on the right are the scope and principles of future research challenges that would address the "*what*
955 *should be done?*" (Loop II, "direct problem") question in echo to the current "*what has been done?*" concern (Loop I).

956

957 **4.2 Research challenges in hydrology and philosophy of modelling**

958 This review has sought the determinants of modelling choices in hydrology (Figure 11, Loop I)
959 from the basis provided by literature sources, without any intention to provide recommendations.
960 However, for most practical applications, the starting point is the definition of a scope and the
961 endpoint is the evaluation of the objective function to evaluate the success or the failure of the chosen
962 modelling strategy. A question thus arises on how to guide the modeller in the choice of an adequate
963 model, in function of given, approximately known spatiotemporal scales, flow typology and
964 dimensionless numbers (Figure 11, Loop II). According to the principle of parsimony, modellers
965 should seek the simplest modelling strategy capable of (i) a realistic representation of the physical
966 processes, (ii) matching the performances of more complex models and (iii) providing the right
967 answers for the right reasons.

968 - (i) Throughout the last decades, an important change of the scope of free-surface flow modelling
969 applications has taken place, with subsequent changes in the objective functions resorted to. The
970 development of hydrological and hydraulic sciences has been directly linked to the progresses in
971 understanding processes, in theoretical model development (e.g. computational facilities: numerical
972 techniques, data assimilation, thorough model exploration, inverse calculus) and in data acquisition
973 (new devices, remote sensing, LiDAR). "*It may seem strange to end a review of modelling with an*
974 *observation that future progress is very strongly linked to the acquisition of new data and to new*
975 *experimental work but that, in our opinion, is the state of the science*" (Hornberger & Boyer 1995).

976 - (ii) However, there remains an important need for research on classical free-surface flow
977 (hydrological or hydraulic) modelling for engineering applications in predicting floods, designing
978 water supply infrastructures and for water resources management, from the headwater catchment to
979 the regional scale. More recently, free-surface flow modelling has become an indispensable tool for
980 many interdisciplinary projects, such as predicting pollution and/or erosion incidents, the impact of
981 anthropogenic and climate change on environmental variables such as water, soil, biology, ecology, or
982 socio-economy and ecosystemic services. The direct consequence is a significant increase of the
983 complexity of the objective function, from simple mono-site (e.g. one-point), mono-variable (e.g. the
984 water depth) and mono-criterion (e.g. the error on peakflow) to complex multi-site (e.g. large number
985 of points within a catchment), multi-variable (e.g. water depth, hydrograph, water table,
986 concentrations, ecological indicators, economic impact) and multi-criteria (e.g. errors on peakflow,
987 volume, RMSE) objective functions.

988 - (iii) There is often a mismatch between model types, site data and objective functions. First,
989 models were developed independently from the specificities of the study site and available data, prior
990 to the definition of any objective function. In using free-surface flow models, the context of their
991 original purpose and development is often lost, so that they may be applied to situations beyond their
992 validity or capabilities. Second, site data are often collected independently of the objectives of the
993 study. Third, the objective function must be specific to the application but also meet standard practices
994 in evaluating model performance, in order to compare modelling results between sites and to
995 communicate the results to other scientists or stakeholders. The known danger is to use flow and
996 erosion equations outside their domains of validity (*i.e.*, breaking the assumptions made during their
997 derivation) then to rely on the calibration of model parameters as for technical compensations of
998 theoretical flaws, at the risk of losing the physical sense of model parameters, creating equifinality and
999 obtaining the “*right results for the wrong reason*” (Klemeš 1986). Choosing the right model for the
1000 right reason is crucial but the identification of the optimal data-model couple to reach a predefined
1001 objective is not straightforward. We need a framework to seek the optimum balance between the
1002 model, data and the objective function as a solution for a hydrological or hydraulic problem, on the

1003 basis of the principle of parsimony. The latter follows a famous quote often attributed to Einstein, that
1004 "*everything should be made as simple as possible, but not simpler*" which somehow originates in the
1005 philosophy of William of Ockham (1317) (*Numquam ponenda est pluralitas sine necessitate*
1006 [*Plurality must never be posited without necessity*]) or may even be traced back to Aristotle's (~350
1007 BCE) *Analytica Posteriora* that already advocated demonstrations relying on the fewest possible
1008 number of conjectures, i.e. the dominant determinisms.

1009 Finally, analytical procedures for free-surface flows and erosion issues necessitates a
1010 comprehensive analysis of the interplay between models (assumptions, accuracy, validity), data
1011 requirements and all contextual information available, encompassed in the "signature" of any given
1012 application: model refinement, spatiotemporal scales, flow typology and scale-independent description
1013 by dimensionless numbers. This review helps the modeller positioning his (or her) case study with
1014 respect to the modelling practices most encountered in the literature, without providing any
1015 recommendation. A complementary step and future research challenge is to decipher relevant
1016 modelling strategies from the available theoretical and practical material, resorting to the same objects,
1017 the previously defined signatures. Its purpose clearly is to address the "*which model, for which scales*
1018 *and objectives?*" question. A complete analytical framework, comprised of both loops, would provide
1019 references and guidelines for modelling strategies. Its normative structure in classifying theoretical
1020 knowledge (the mathematics world, equations and models) and contextual descriptions (real-life
1021 physical processes, scales and typologies) hopefully makes it also relevant for other Earth Sciences.

1022

1023

1024

1025 **Appendix A. References used in the Figures.**

1026

1027 Abrahams & Parsons (1994), Afzalimehr & Anctil (2000), Afzalimehr et al. (2007), Akan & Yen
1028 (1981), Alonso et al. (2002), Audusse et al. (2008), Aziz & Scott (1989), Bajracharya & Barry (1997),
1029 Bates & De Roo (2000), Bathurst (2006), Belaud & Paquier (2001), Beltaos et al. (2012), Berger &
1030 Stockstill (1995), Blandford & Meadows (1990), Booker et al. (2001), Bounvilay (2003), Burguete et
1031 al. (2008), Camacho & Lees (1999), Canovaro & Solari (2007), Cao et al. (2004), Cassan & Belaud
1032 (2012), Cassan et al. (2012), Chahinian et al. (2005), Charlier (2007), Charlier et al. (2009), Charpin &
1033 Myers (2005), Charru et al. (2004), Chartrand & Whiting (2000), Chen & Wu (2000), Chiari (2008),
1034 Chin (1999), Chinnarasri and Wongwise (2006), Choi & Molinas (1993), Church & Zimmermann
1035 (2007), Davies et al. (1997), Devauchelle et al. (2007), Dunkerley (2003), Dunkerley (2004), Einstein
1036 (1950), Elhanafy et al. (2008), Emmett (1970), Engelund & Fredsoe (1976), Fan & Li (2006), Feng &
1037 Michaelides (2002), Fovet et al. (2013), Gao & Abrahams (2004), Garcia & Parker (1993), Gaur &
1038 Mathur (2003), Gejadze & Copeland (2006), Gerbeau & Perthame (2001), Ghavasieh et al. (2001),
1039 Gimenez & Govers (2001), Gimenez et al. (2004), Gomez-Delgado et al. (2011), Govers (1992), Grant
1040 et al. (1990), Guinot & Cappelaere (2009), Hallema & Moussa (2013), Hauke (2002), Hayami (1951),
1041 Henine et al. (2014), Hessel (2006), Hessel et al. (2003), Horritt & Bates (2002), Hromadka &
1042 DeVries (1988), Jain & Singh (2005), Järvelä (2005), Keshavarzy & Ball (1997), Keskin &
1043 Agiraliloglu (1997), Keulegan (1938), Koussis (1978), Lajeunesse et al. (2010), Lamarre & Roy
1044 (2008), Lamb et al. (2008), Lane et al. (1994), Lawless & Robert (2001), Lawrence (2000), Leopold et
1045 al. (1960), Liang et al. (1996), Liu et al. (2003), Liu et al. (2007), Lobkovsky et al. (2008), Lyn
1046 (1992), Malverti et al. (2008), Mangeney et al. (2007), McDonald et al. (1995b), Meile (2007),
1047 Métivier & Meunier (2003), Meyer-Peter & Müller (1948), Mizanur & Chaudhry (1995), Mohapatra
1048 & Ballamudi (1996), Morgan et al. (1998), Morin et al. (2009), Moussa (1996), Moussa & Bocquillon
1049 (1996a), Moussa & Bocquillon (1996b), Moussa & Bocquillon (2009), Moussa & Chahinian (2009),
1050 Moussa et al. (2002), Moussa et al. (2007), Mügler et al. (2010), Munier et al. (2008), Nepf (1999),

1051 Nikora et al. (2001), Nikora et al. (2008), Nino et al. (2003), Nord & Esteves (2010), Paiva et al.
1052 (2013), Parsons et al. (1997), Perkins & Koussis (1996), Perumal & Price (2013), Peyras et al. (1992),
1053 Pokrajac et al. (2007), Polyakov & Nearing (2003), Ponce et al. (1978), Ponce et al. (1996), Prahel et al.
1054 (2007), Prosser et al. (1994), Rathburn & Wohl (2003), Rauws (1980), Reddy et al. (2007a), Reddy et
1055 al. (2007b), Rodellar et al. (1993), Roux & Dartus (2006), Rutschmann & Hager (1996), Saleh et al.
1056 (2013), Sau et al. (2010), Savat (1980), Schindler & Robert (2004), Schmeeckle & Nelson (2003),
1057 Schmeeckle et al. (2007), Sear (1996), Sen & Garg (2012), Shields (1936), Simpson & Castelltort
1058 (2006), Sivakumaran & Yevyevich (1987), Sivapalan et al. (1997), Smart (1984), Smith & McLean
1059 (1977), Stevenson et al. (2002), Tatard et al. (2008), Tiemeyer & al. (2007), Todini & Bossi (1986),
1060 Trigg et al. (2009), van Maren (2007), Vieux et al. (2004), Wang & Chen (2003), Wang et al. (2006),
1061 Wang et al. (2014), Weichert (2006), Williams (1970), Wainwright et al. (2008), Wu & Lee (2001),
1062 Yager et al. (2007), Zhou (1995), Zimmermann & Church (2001), Zoppou & O'Neill (1982).

1063

1064

1065 **Acknowledgements**

1066 The authors would like to thank Gilles Belaud (SupAgro Montpellier), Claude Bocquillon
1067 (University of Montpellier II) and Pierre-Olivier Malaterre (IRSTEA Montpellier) for fruitful
1068 discussions. [We are also grateful to the two anonymous peer reviewers for their recommendations and](#)
1069 [to HESS for financial support.](#)

1070

1071 **References**

- 1072 | Abbott, M. B. (1979). Computational Hydraulics. Pitman, London, 324pp.
- 1073 | Abrahams, A. D. and Parsons, A. J. (1994). Hydraulics of interrill overland flow on stone-covered
1074 desert surfaces, *Catena*, 23, pp.111-140.
- 1075 | Achdou, Y, Pironneau, O. and Valentin, F. (1998). Effective boundary conditions for laminar flows
1076 over periodic rough boundaries, *Journal of Computational Physics*, 147, pp.187-218.
- 1077 | Afzalimehr, H. and Anctil, F. (1999). Velocity distribution and shear velocity behavior of
1078 decelerating flows over a gravel bed, *Canadian Journal of Civil Engineering*, 26 (4), pp.468-475.
- 1079 | Afzalimehr, H., Dey, S. and Rasoulianfar, P. (2007). Influence of decelerating flow on incipient
1080 motion of a gravel-bed stream, *Sadhana-Academy Proceedings in Engineering Sciences*, 32 (5),
1081 pp.545-559.
- 1082 | Akan, A. O. and Yen, B. C. (1981). Diffusion-wave flood routing in channel networks. *Journal of*
1083 *the Hydraulic Division, ASCE*, vol 107, N°HY6, pp.719-732.
- 1084 | Aksoy, H. and Kavvas, M. L. (2005) A review of hillslope and watershed scale erosion and
1085 sediment transport models, *Catena*, 64 (2-3), pp.247-271.
- 1086 | Alavian, V., Jirka, G. H., Denton, R. A., Johnson, M. A., Stefan, H. G. (1992). Density currents
1087 entering lakes and reservoirs, *Journal of Hydraulic Engineering*, 118 (11), pp.1464-1489.
- 1088 | Allen, P. A. (2008). Time scales of tectonic landscapes and their sediment routing systems,
1089 *Geological Society of London, Special Publications*, 296, pp.7-28.
- 1090 | Alonso, C. V., Bennett, S. J. and Stein, O.R. (2002). Predicting head cut erosion and migration in
1091 concentrated flows typical of upland areas, *Water Resources Research*, 38 (12), 1303.
- 1092 | Alonso, R., Santillana, M. and Dawson, C. (2008). On the diffusive wave approximation of the
1093 shallow water equations, *European Journal of Applied Mathematics*, 19 (5), pp.575-606.
- 1094 | Ascough II, J. C., Baffaut, C., Nearing, M. A. and Liu, B.Y. (1997). The WEPP watershed model:
1095 I. Hydrology and erosion, *Transactions of the ASAE*, 40(4), pp.921–933.

1096 Audusse, E., Bristeau, M. O. and Decoene, A. (2008). Numerical simulations of 3D surface flows
1097 by a multilayer Saint-Venant model, *International Journal for Numerical Methods in Fluids*, 56 (3),
1098 pp.331-350.

1099 Aziz, N. M. and Scott, D. E. (1989). Experiments on sediment transport in shallow flows in high
1100 gradient channels, *Hydrological Sciences Journal*, 34, pp.465–478.

1101 Bagnold, R. A. (1954). Experiments on the gravity-free dispersion of large solid spheres in a
1102 Newtonian fluid under shear, *Proceedings of the Royal Society of London-Series A*, 225 (964), pp.49-
1103 63.

1104 Bagnold, R. A. (1956). The flow of cohesionless grains in fluids, *Philosophical Transactions of the*
1105 *Royal Society of London*, 249 (964), pp.235-297.

1106 Bagnold, R. A. (1966). An approach to the sediment transport problem from general physics, US
1107 Geological Survey Professional Paper 442-I, Washington D.C., US Government Printing Office,
1108 42 pp.

1109 Bajracharya, K. and Barry, D. A. (1997). Accuracy criteria for linearised diffusion wave flood
1110 routing, *Journal of Hydrology*, 195 (1-4), pp.200-217.

1111 Barenblatt, G. I. (1987) *Dimensional Analysis*, New York: Gordon and Breach Science Publishers,
1112 135 pp.

1113 Basson, A. and Gerard-Varet, D. (2008). Wall laws for fluid flows at a boundary with random
1114 roughness, *Communications on Pure and Applied Mathematics*, 61 (7), pp.941-987.

1115 Bates, P. D. and De Roo, A. P. J. (2000). A simple raster-based model for flood inundation
1116 simulation, *Journal of Hydrology*, 236, pp.54-77.

1117 Bates, P.D., Horritt, M.S. and Fextrell, T.J. (2010). A simple inertial formulation of the shallow
1118 water equations for efficient two-dimensional flood inundation model. *Journal of Hydrology*, 387, 33-
1119 45.

1120 Bathurst, J. C. (1985). Flow resistance estimation in mountain rivers, *Journal of Hydraulic*
1121 *Engineering – ASCE*, 111 (4), pp.625-643.

1122 Bathurst, J. C. (2006). At-a-site variation and minimum flow resistance for mountain rivers, *Journal*
1123 *of Hydrology*, 269 (1-2), pp.11-26.

- 1124 Bayazit, M. (1976). Free surface flow in a channel of large relative roughness, *Journal of Hydraulic*
1125 *Research*, 14 (2), pp.115-126.
- 1126 Beasley, D.B., Huggins, L.F. and Monke, E.J. (1980). ANSWERS: a model for watershed
1127 planning. *Transactions of the ASAE*, pp.938-944.
- 1128 Belaud, G. and Paquier, A. (2001). Sediment diversion through irrigation outlets, *Journal of*
1129 *Irrigation and Drainage Engineering*, 127 (1), pp.35-38.
- 1130 Beltaos, S., Tang, P. and Rowsell, R. (2012). Ice jam modelling and field data collection for flood
1131 forecasting in the Saint John River, Canada. *Hydrological Processes*, 26, pp.2535- 2545.
- 1132 Bennett, J. P. (1974). Concepts of mathematical modelling of sediment yield, *Water Resources*
1133 *Research*, 10, pp.485-492.
- 1134 Bennett, S. J., Alonso, C. V, Prasad, S. N. and Römkens, M. J. M. (2000). Experiments on headcut
1135 growth and migration in concentrated flows typical of upland areas, *Water Resources Research*, 36
1136 (7), pp.1911-1922.
- 1137 Berger, R. C. and Stockstill, R. L. (1995). Finite-element model for high-velocity channels. *J*
1138 *Hydraulic Engineering*, 121, 10, pp.710-715.
- 1139 Bertrand, J. (1878). Sur l'homogénéité dans les formules de physique, *Comptes Rendus de*
1140 *l'Académie des Sciences*, 86 (15), pp.916-920.
- 1141 Best, J.L. (1992), On the entrainment of sediment and initiation of bed defects: insights from recent
1142 developments within turbulent boundary layer research, *Sedimentology*, 39, pp.797-811.
- 1143 Beven K. J. (2000). *Rainfall-runoff modelling, the primer*, John Wiley & Sons, Chichester, UK
1144 360 pp
- 1145 Blandford, G. E. and Meadows, M. E. (1990). Finite element simulation of nonlinear kinematic
1146 surface runoff. *Journal of Hydrology*, 119, pp.335-356.
- 1147 Blöschl, G. and Sivapalan, M. (1995). Scale issues in hydrological modeling: a review,
1148 *Hydrological Processes*, 9, pp.251-290.
- 1149 Boardman, J. (2006). Soil erosion science: Reflections on the limitation of current approaches,
1150 *Catena*, 68 (2-3), pp.73-86.

1151 Booker, D. J., Sear, D. A. and Payne, A. J. (2001). Modelling three-dimensional flow structures
1152 and patterns of boundary shear stress in a natural pool-riffle sequence, *Earth Surface Processes and*
1153 *Landforms*, 26, pp.553-576.

1154 Bouchut, F., Mangeney-Castelnau, A., Perthame, B., Vilotte, J.-P. (2003). A new model of Saint-
1155 Venant and Savage-Hutter type for gravity driven shallow water flows, *Comptes-Rendus de*
1156 *l'Académie des Sciences de Paris*, 336, pp.531-536.

1157
1158 Bouchut, F. and Westdickenberg, M. (2004). Gravity-driven shallow water models for arbitrary
1159 topography, *Communications in Mathematical Sciences*, 2, pp.359-389.

1160 Bounvilay, B. (2003). Transport velocities of bedload particles in rough open channels flows, PhD
1161 Thesis, Department of Civil Engineering, Colorado State University, USA, 2003.

1162 Boussinesq, J. (1877). Essai sur la théorie des eaux courantes. Mémoires à l'Académie des
1163 Sciences. T23-24, pp.1-680.

1164 Bracken, L. J., Wainwright, J., Ali, G. A., Tetzlaff, D., Smith, M. W., Reaney, S. M. and Roy, A.
1165 G. (2013). Concepts of hydrological connectivity: research approaches, pathways and future agendas,
1166 *Earth Science Reviews*, 119, pp.17-34.

1167 Bridgman, P.W. (1922). *Dimensional Analysis*, Yale University Press, New Haven.

1168 Bridgman, P.W. (1963). *Dimensional Analysis*, Yale Paperbound, New Haven, 113 pp.

1169 Buckingham, E. (1914). On physically similar systems; illustrations of the use of dimensional
1170 equations, *Physical Review*, 4, pp.345-376.

1171 Bucur, D., Feireisl, E. and Necasova, S. (2010). Boundary behavior of viscous fluids: influence of
1172 wall roughness and friction-driven boundary conditions, *Archive for Rational Mechanics Analysis*,
1173 197 (1), pp.117-138.

1174 Bucur, D., Feireisl, E., Necasova, S. and Wolf, J. (2008). On the asymptotic limit of the Navier-
1175 Stokes system on domains with rough boundaries, *Journal of Differential Equations*, 244 (11),
1176 pp.2890-2908.

1177 Buffington, J. M. and Montgomery, D. R. (1997). A systematic analysis of eight decades of
1178 incipient motion studies, with special reference to gravel-bedded rivers. *Water Resources Research*,
1179 33, pp.1993-2029.

- 1180 Burguete, J., Garcia-Navarro, P. and Murillo, J. (2008). Friction-term Discretization and limitation
1181 to preserve stability and conservation in the 1D shallow-water model: application to unsteady
1182 irrigation and river flow, *International Journal for Numerical Methods in Fluids*, 58, pp.403-425.
- 1183 Camacho, L. A. and Lees, M. J. (1999). Multilinear discrete lag-cascade model for channel routing,
1184 *Journal of Hydrology*, 226, pp.30-47.
- 1185 Campisano, A., Creaco, E. and Modica, C. (2004). Experimental and numerical analysis of the
1186 scouring effects of flushing waves on sediment deposits, *Journal of Hydrology*, 299 (3-4), pp.324-334.
- 1187 Canovaro, F. and Solari, L. (2007). Dissipative analogies between a schematic macro-roughness
1188 arrangement and step-pool morphology, *Earth Surface Processes and Landforms*, 32 (11), pp.1628-
1189 1640.
- 1190 Cao, Z., Pender, G., Wallis, S. and Carling, P. (2004). Computational dam-break hydraulics over
1191 erodible sediment bed, *Journal of Hydraulic Engineering-ASCE*, 130 (7), pp.689-703.
- 1192 Carlier M. (1980). *Hydraulique générale et appliquée*. Eyrolles, Paris, 565 pp.
- 1193 Carollo, F.G., V. Ferro, and D. Termini (2005), Analyzing turbulence intensity in gravel bed
1194 channels, *Journal of Hydraulic Engineering*, 131 (12), pp.1050-1061.
- 1195 Casado-Diaz, J., Fernandez-Cara, E. and Simon, J. (2003). Why viscous fluids adhere to rugose
1196 walls: a mathematical explanation, *Journal of Differential Equations*, 189 (2), pp.526-537.
- 1197 Cassan, L. and Belaud, G. (2012). Experimental and numerical investigation of flow under sluice
1198 gates, *Journal of Hydraulic Engineering*, 138, pp.367-373.
- 1199 Cassan, L., Belaud, G., Baume, J. and Dejean, C. (2012) Seasonal Variation of Velocity Fields in
1200 Lined Channels: Impact on Flow Measurement. World Environmental and Water Resources Congress
1201 [20-22 May 2012](#); pp.-2188-2197.
- 1202 Chahinian, N., Moussa, R., Andrieux, P. and Voltz, M. (2005). Comparison of infiltration models
1203 to simulate flood events at the field scale, *Journal of Hydrology*, 306, pp.191-214.
- 1204 Charlier, J.-B. (2007). *Fonctionnement et modélisation hydrologique d'un petit bassin versant*
1205 *cultivé en milieu volcanique tropical*. Thèse de Doctorat, Université de Montpellier 2, [Montpellier,](#)
1206 [France](#), 246 pp.

- 1207 Charlier, J.-B., Cattan, P., Voltz, M. and Moussa, R. (2009). Transport of a nematicide in surface
1208 and ground waters in a tropical volcanic catchment, *Journal of Environmental Quality*, 38, pp.1031-
1209 1041.
- 1210 Charpin, J. P. F. and Myers, T. G. (2005). Modelling thin film flow with erosion and deposition,
1211 *Advances in Water Resources*, 28, pp.761-722.
- 1212 Charru, F. (2006). Selection of the ripple length on a granular bed sheared by a liquid flow,
1213 *Physics of Fluids*, 18, 121508.
- 1214 Charru, F., Mouilleron-Arnould, H. and Eiff, O. (2004). Erosion and deposition of particles on a
1215 bed sheared by a viscous flow, *Journal of Fluid Mechanics*, 519, pp.59-80.
- 1216 Chartrand, S. M. and Whiting, P. J. (2000). Alluvial architecture in headwater streams with special
1217 emphasis on step-pool topography, *Earth Surface Processes and Landforms*, 25, pp.583-600.
- 1218 Chen, C. L. (1992). Momentum and energy coefficients based on power-law velocity profile.
1219 *Journal of Hydraulic Engineering*, ASCE, 118 (11), pp.1571-1584.
- 1220 Chen, R. C. and Wu, J. L. (2000). The flow characteristics between two interactive spheres,
1221 *Chemical Engineering Science*, 55, pp.1143-1158.
- 1222 Chézy, A. de (1775). Mémoire sur la vitesse de l'eau conduite dans une rigole donnée, *Fonds*
1223 *Ancien de l'Ecole Nationale des Ponts et Chaussées – No.847*. Reprinted in: *Annales des Ponts et*
1224 *Chaussées*, 60, 1921.
- 1225 Chiari, M. (2008). Numerical modelling of bedload transport in torrents and mountain streams, Phd
1226 Thesis, University of natural resources and applied life sciences, Vienna, [Austria](#), 212 pp.
- 1227 Chin, A. (1999). The morphologic structure of step-pools in mountain streams, *Geomorphology*,
1228 27, pp.191-204.
- 1229 Chin, A. and Wohl, E. (2005). Toward a theory for step pools in stream channels, *Progress in*
1230 *Physical Geography*, 29 (3), pp.275-296.
- 1231 Chinnarasri, C. and Wongwise, S. (2006). Flow patterns and energy dissipation over various
1232 stepped chutes, *Journal of Irrigation and Drainage Engineering – ASCE*, 132, pp.70-76.

- 1233 Choi, G. W. and Molinas, A. (1993). Simultaneous solution algorithm for channel network
1234 modelling. *Water Resources Research*, 29 (2), pp.321-328.
- 1235 Chow, V. T. (1959). *Open-Channel Hydraulics*. McGraw Hill, New York, 680 pp.
- 1236 Church, M. (2002). Geomorphic thresholds in riverine landscapes, *Freshwater Biology*, 47,
1237 pp.541-557.
- 1238 Church, M. and Zimmermann, A. (2007). Form and stability of step-pool channels: research
1239 progress, *Water Resources Research*, 43 (3), W03415, doi:10.1029/2006WR005037.
- 1240 Colebrook, C. F. and White, C. M. (1937). Experiments with fluid friction in roughened pipes,
1241 *Proceedings of the Royal Society of London – Series A: Mathematical and Physical Sciences*, 161
1242 (906), pp.367-381.
- 1243 Coleman, N. L. (1967). A theoretical and experimental study of drag and lift forces acting on a
1244 sphere resting on a hypothetical streambed, in *Proceedings of 12th IAHR Congress*, vol. 3, pp. 185–
1245 192, Int. Assoc. of Hydraul. Eng. and Res., Madrid.
- 1246 Cooper, J. R., Aberle, J., Koll, K. and Tait, S. J. (2013). Influence of relative submergence on
1247 spatial variance and form-induced stress of gravel-bed flows, *Water Resources Research*, 49, pp.5765-
1248 5777.
- 1249 Croke, J. and Mockler, S. (2001). Gully initiation and road-to-stream linkage in a forested
1250 catchment, south-eastern Australia, *Earth-Surface Processes and Landforms*, 26, pp.205-217.
- 1251 Cunge, J. A. (1969). On the subject of a flood propagation computation method (Muskingum
1252 method), *Journal of Hydraulic Research*, 7 (2), pp.205-230.
- 1253 Cunge, J., Holly, F. M. and Verwey, A. (1980). *Practical aspects of computational river hydraulics*.
1254 Pitman Advanced Publishing Program, London, 420 pp.
- 1255 Daluz Vieira, J.H. (1983). Conditions governing the use of approximations for the Saint-Venant
1256 equations for shallow water flow, *Journal of Hydrology*, 60, pp.43-58.
- 1257 Darcy, H. (1857). *Recherches expérimentales relatives au mouvement de l'eau dans les tuyaux*,
1258 Mallet-Bachelier, Paris. 268 pages and atlas (in French).

- 1259 Davies, A. G., Ribberink, J. S., Temperville, A. and Zyserman, J. A. (1997). Comparisons between
1260 sediment transport models and observations made in wave and current flows above plane beds,
1261 Coastal Engineering, 31, pp. 163-198.
- 1262 de Marsily, G. (1986). Quantitative hydrogeology, Academic Press, Inc., Orlando, FL, USA,
1263 464 pp.
- 1264 de Marsily, G. (1994). On the use of models in hydrology (free opinion), Revue des Sciences de
1265 l'Eau, 7, pp.219-234.
- 1266 de Roo, A. P. J., Wesseling, C. G. and Ritsema, C. J. (1996). LISEM: a single event physically-
1267 based hydrologic and soil erosion model for drainage basins. I: Theory, input and output, Hydrological
1268 Processes, 10, pp.1107-1117.
- 1269 Devauchelle, O., Josserand, C., Lagrée, P.-Y. and Zaleski, S. (2007). Morphodynamic modelling of
1270 erodible laminar channels, Physical Review E, 76, 056318, doi: [10.1103/PhysRevE.76.056318](https://doi.org/10.1103/PhysRevE.76.056318).
- 1271 Devauchelle, O., Malverti, L., Lajeunesse, E., Josserand, C., Lagrée, P.-Y. and Métivier, F. (2010).
1272 Rhomboid beach pattern : a laboratory investigation, Journal of Geophysical Research, 115, F02017,
1273 doi: [10.1029/2009JF001471](https://doi.org/10.1029/2009JF001471).
- 1274 Dey, S. and Papanicolaou, A. (2008). Sediment threshold under stream flow: a state-of-the-art
1275 review, KSCE Journal of Civil Engineering, 12 (1), pp.45-60.
- 1276 Dittrich, A., and K. Koll (1997). Velocity field and resistance of flow over rough surfaces with
1277 large and small relative roughness, International Journal of Sediment Research, 12 (3), pp.21-33.
- 1278 Drake, T.G., R.L. Shreve, W.E. Dietrich, P.J. Whiting, and L.B. Leopold (1988). Bedload transport
1279 of fine gravel observed by motion picture photography, Journal of Fluid Mechanics, 192, pp.193-217.
- 1280 Du Boys, D. (1879). Le Rhône et les rivières à lit affouillables. Annales des Ponts et Chaussées,
1281 Série 5, 18, pp.141-195.
- 1282 Du Buat, P. (1779). Principes d'hydraulique et de pyrodynamique vérifiés par un grand nombre
1283 d'expériences faites par ordre du gouvernement, Firmin Didot Ed., Paris
- 1284 Dunkerley, D. (2003). Determining friction coefficients for interrill flows: the significance of flow
1285 filaments and backwater effects, Earth Surface Processes and Landforms, 28 (5), pp.475-491.

Mis en forme : Français (France)

- 1286 Dunkerley, D. (2004). Flow threads in surface run-off: implications for the assessment of flow
1287 properties and friction coefficients in soil erosion and hydraulics investigations, *Earth Surface*
1288 *Processes and Landforms*, 29 (8), pp.1011-1026.
- 1289 Einstein, H. A. (1950). The bed-load function for sediment transportation in open channel flows,
1290 US Department of Agriculture, Soil Conservation Service, Technical Bulletin No. 1026, 74 pp.
- 1291 Einstein, H. A. and Banks, R. B. (1950). Fluid resistance of composite roughness, *Transactions of*
1292 *the American Geophysical Union*, 31 (4), pp.603-610.
- 1293 Elhanafy, H., Copeland, G. J. M. and Gejadze, I. Y. (2008). Estimation of predictive uncertainties
1294 in flood wave propagation in a river channel using adjoint sensitivity analysis, *International Journal*
1295 *for Numerical Methods in Fluids*, 56 (8), pp.1201-1207.
- 1296 Emmett, W. W. (1970). The hydraulics of overland flow on hillslopes. United States Geological
1297 Survey, Professional Paper 662-A, US Government Printing Office: Washington, DC.
- 1298 Engelund, F. and Fredsoe, J. (1976). Sediment transport model for straight alluvial channels,
1299 *Nordic Hydrology*, 7 (5), pp.293-306.
- 1300 Exner, F. M. (1925). Über die Wechselwirkung zwischen Wasser und Geschiebe in Flüssen,
1301 *Sitzungsberichte der kaiserlichen Akademie der Wissenschaften Wien, Abteilung IIa*, 134, pp.165-205.
- 1302 Fan, P. and Li, J. C. (2006). Diffusive wave solutions for open channel flows with uniform and
1303 concentrated lateral inflow, *Advances in Water Resources*, 29, pp.1000-1019.
- 1304 Feng, Z. G. and Michaelides, E. E. (2002). Interparticle forces and lift on a particle attached to a
1305 solid boundary in suspension flow, *Physics of Fluids*, 14 (1), pp.49-60.
- 1306 Ferguson, R.I. (2007). Flow resistance equations for gravel- and boulder-bed streams. *Water*
1307 *Resources Research*, 43, W05427, doi: 10.1029/2006WR005422.
- 1308 Fernandez-Luque, R., and van Beek, R.(1976). Erosion and Transport of Bed-Load Sediment,
1309 *Journal of Hydraulic Research*, 14(2), pp.127-144.
- 1310 Ferrick, M. G. (1985). Analysis of wave types, *Water Resources Research*, 21, pp.209-212.
- 1311 Ferro, V. (2003). Flow resistance in gravel-bed channels with large-scale roughness, *Earth Surface*
1312 *Processes and Landforms*, 28 (12), pp.1325-1339.

- 1313 Fourier, J. B. (1822) *Théorie analytique de la chaleur*, Paris.
- 1314 Fovet, O., Litrico, X., Belaud, G. and Genthon, O. (2013). Adaptive control of algae detachment in
1315 regulated canal networks, *Journal of Hydroinformatics*, 15 (2), pp.321-334.
- 1316 French, R. H. (1985). *Open-channel hydraulics*, New York, McGraw-Hill, 705 pp.
- 1317 Froude, W. (1868). Observations and suggestions on the subject of determining by experiment the
1318 resistance of ships, *Correspondence with the Admiralty*, Chelston Cross, December 1868. Reprinted in
1319 “The papers of William Froude”, The Institution of Naval Architects, London, 1955, pp.120-128.
- 1320 Gao, P, Abrahams, A. D. (2004). Bedload transport resistance in rough open-channel flows. *Earth*
1321 *Surface Processes and Landforms*, 29, pp.423–435.
- 1322 Garcia, M. and Parker, G. (1993). Experiments on the entrainment of sediment into suspension by a
1323 dense bottom current, *Journal of Geophysical Research-Oceans*, 98 (C3), pp.4793-4807.
- 1324 Gaur, M. L. and Mathur, B. S. (2003). Modeling event-based temporal variability of flow
1325 resistance coefficient, *Journal of Hydrologic Engineering*, 8 (5), pp.266-277.
- 1326 Gejadze, I. Y. and Copeland, G. J. M. (2006). Open boundary control problem for Navier-Stokes
1327 equations including a free surface: adjoint sensitivity analysis, *Computer and Mathematics with*
1328 *Applications*, 52 (8-9), pp.1243-1268.
- 1329 Gerard-Varet, D. and Masmoudi, N. (2010). Relevance of the slip condition for fluid flows near an
1330 irregular boundary, *Communications in Mathematical Physics*, 295 (1), pp.99-137.
- 1331 Gerbeau, J.-F. and Perthame, B. (2001). Derivation of a viscous Saint-Venant system for laminar
1332 shallow water; numerical validation, *Discrete and Continuous Dynamical Systems-Series B*, 1 (1),
1333 pp.89-102.
- 1334 Ghavasieh, A.-R., Poulard, C. and Paquier, A. (2001). Effect of roughened strips on flood
1335 propagation: assessment on representative virtual cases and validation, *Journal of Hydrology*, 318 (1-
1336 4), pp.121-137.
- 1337 Gimenez, R. and Govers, G. (2001). Interaction between bed roughness and flow hydraulics in
1338 eroding rills, *Water Resources Research*, 37 (3), pp.791-799.

- 1339 Gimenez, R., Planchon, O., Silvera, N. and Govers, G. (2004). Longitudinal velocity patterns and
1340 bed morphology interaction in a rill, *Earth Surface Processes and Landforms*, 29 (1), pp.105-114.
- 1341 Gioia, G. and Bombardelli, F. A. (2001). Scaling and similarity in rough channel flows, *Physical*
1342 *Review Letters*, 88, 014501, doi: [10.1103/PhysRevLett.88.014501](https://doi.org/10.1103/PhysRevLett.88.014501).
- 1343 Govers, G. (1992). Evaluation of transporting capacity formulae for overland flow. In: *Overland*
1344 *Flow Hydraulics and Erosion Mechanics*, Parsons A. J. and Abrahams, A. D. (eds), UCL Press,
1345 London, pp.243-273.
- 1346 Grant, G. E. (1997). Critical flow constrains flow hydraulics in mobile-bed streams: a new
1347 hypothesis, *Water Resources Research*, 33 (2), pp.349-358.
- 1348 Grant, G. E., Swanson, F. J. and Wolman, M. G. (1990). Pattern and origin of stepped-bed
1349 morphology in high-gradient streams, Wetern Cascades, Oregon, *Geological Society of America*
1350 *Bulletin*, 102 (3), pp.340-352.
- 1351 Gresho, P. M. and Sani, R. L. (1998). *Incompressible Flow and the Finite Element Method*.
1352 [Wiley](#) John Wiley & Sons, Inc., New York, NY, USA.
- 1353 Guinot, V. and Cappelaere, B. (2009). Sensitivity analysis of 2D steady-state shallow water flow.
1354 Application to the surface flow model calibration, *Advances in Water Resources*, 32 (4), pp.540-560.
- 1355 Gumiere, S. J., Raclot, D., Cheviron, B., Davy, G., Louchart, X., Fabre, J. C., Moussa, R., Le
1356 Bissonnais, Y. (2011). MHYDAS-Erosion : a distributed single-storm water erosion model for
1357 agricultural catchment, *Hydrological Processes*, 25 (11), pp.1717-1728.
- 1358 Hairsine, P. B. and Rose, C. W. (1992a). Modeling water erosion due to overland flow using
1359 physical principles. 1. Sheet flow, *Water Resources Research*, 28 (1), pp.237-243.
- 1360 Hairsine, P. B. and Rose, C. W. (1992b). Modeling water erosion due to overland flow using
1361 physical principles. 2. Rill flow, *Water Resources Research*, 28 (1), pp.245-250.
- 1362 Hallema, D. and Moussa, R. (2013). A model for distributed GIUH-based flow routing on natural
1363 and anthropogenic hillslopes, *Hydrological Processes*, 28, pp.4877-4895, doiDOI: [10.1002/hyp.9984](https://doi.org/10.1002/hyp.9984).
- 1364 Härtel, C. (1996). *Turbulent flows: direct numerical simulation and large-eddy simulation*, In:
1365 *Handbook of Computational Fluid Mechanics*, R. Peyret Ed., Elsevier, New York, pp.283-338.

Mis en forme : Espace Avant : 0 pt,
Après : 0 pt

- 1366 Hauke, G. (2002). A stabilized finite element method for the Saint-Venant equations with
1367 application to irrigation, *International Journal for Numerical Methods in Fluids*, 38 (10), pp.963-984.
- 1368 Hayami, S. (1951). On the propagation of flood waves, *Disaster Prevention Research Institute*
1369 *Bulletin*, 1, pp.1-16.
- 1370 Henderson, F.M. (1966). *Open Channel Hydraulics*. MacMillan and Co., New York.
- 1371 Hénine, H., Nédélec, Y. and Ribstein, P. (2014). Coupled modelling of the effect of overpressure
1372 on water discharge in a tile drainage system. *Journal of Hydrology*, 511, pp.39–48.
- 1373 Hessel, R. (2006). Consequences of hyperconcentrated flow for process-based soil erosion
1374 modelling on the Chinese Loess Plateau, *Earth Surface Processes and Landforms*, 31, pp.1100-1114.
- 1375 Hessel, R., Jetten, V. and Ganghui, Z. (2003). Estimating Manning's n for steep slopes, *Catena*, 54,
1376 pp.77-91.
- 1377 Hino, M. (1963) Turbulent flows with suspended particles, *Journal of the Hydraulic Division*
1378 *ASCE*, 89(2a), pp.161-185.
- 1379 Hjulström, F. (1935), Studies of the morphological activity of rivers as illustrated by the river
1380 Fyris, *Bulletin of the Geology Institute of Uppsala*, 25, pp.221-527.
- 1381 Hornberger, G.M. and Boyer, E. W. (1995). Recent advances in watershed modelling. US National
1382 report to the IUGG 1991–1994. *Reviews of Geophysics*, 33, Supplement S2, pp.949–957.
- 1383 Horritt, M. S. and Bates, P. D. (2002). Evaluation of 1D and 2D numerical models for predicting
1384 river flood inundation, *Journal of Hydrology*, 268 (1-4), pp.87-99.
- 1385 Horton, R. E. (1945). Erosional development of streams and their drainage basins: hydrological
1386 approach to quantitative morphology, *Bulletin of the Geological Society of America*, 56 (3), pp.275-
1387 370.
- 1388 Howard, A. D. (1994). A detachment-limited model of drainage basin evolution, *Water Resources*
1389 *Research*, 30, pp.2261-2285.
- 1390 Ikeda, S. (1982). Incipient motion of sand particles on side slopes, *Journal of Hydraulic*
1391 *Engineering*, 108, pp.95–114.

- 1392 Iwagaki Y. (1955). Fundamental studies on the runoff analysis by characteristics, Disaster
1393 Prevention Research Institute Bulletin, 10, Kyoto, 25 pp.
- 1394 Izumi, N. and Parker, G. (1995). Linear stability of channel inception: downstream-driven theory,
1395 Journal of Fluid Mechanics, 283, pp.341-363.
- 1396 Jäger, W. and Mikelic, A. (2001). On the roughness-induced effective boundary conditions for an
1397 incompressible viscous flow, Journal of Differential Equations, 170 (1), pp.96-122.
- 1398 Jäger, W. and Mikelic, A. (2003). Couette flow over a rough boundary and drag reduction,
1399 Communications in Mathematical Physics, 232, pp.429-455.
- 1400 Jain, M. K. and Singh, V. P. (2005). DEM-based modelling of surface runoff using diffusive wave
1401 equation, Journal of Hydrology, 302, pp.107-126.
- 1402 Jansons, K. M. (1988). Determination of the macroscopic (partial) slip boundary condition for a
1403 viscous flow over a randomly rough surface with a perfect slip microscopic boundary condition,
1404 Physics of Fluids, 31 (1), pp.15-17.
- 1405 Järvelä, J. (2005). Effect of flexible vegetation on flow structure and resistance, Journal of
1406 Hydrology, 307, pp.233-241.
- 1407 Jetten, V., de Roo, A. and Favis-Mortlock, D. (1999). Evaluation of field-scale and catchment-scale
1408 soil erosion models, Catena, 37 (3-4), pp.521-541.
- 1409 Jetten, V., Govers, G. and Hessel, R. (2003). Erosion models: quality of spatial predictions,
1410 Hydrological Processes, 17, pp.887-900.
- 1411 Johnson, R.W. (1998). Handbook of fluid ~~dynamics~~~~mechanics~~, CRC Press, Boca Raton, FL, USA,
1412 1952 pp.
- 1413 Julien, P.Y. (2010). Erosion and sedimentation. Cambridge, ~~UK-University Press~~.
- 1414 Julien, P. Y. and Bounvilay, B. (2013). Velocity of rolling bed load particles, ASCE - Journal of
1415 Hydraulic Engineering, 139 (2), pp.177-186.
- 1416 Julien, P. Y. and Simons, D. B. (1985). Sediment transport capacity of overland flow, Transactions
1417 of the American Society of Agricultural Engineers, 28 (3), pp.755-762.

1418 Katopodes, N.D. (1982). On zero-inertia and kinematic waves, Journal of Hydraulic Engineering,
1419 American Society of Civil Engineers, 108(HY11), pp.1380-1385.

1420 Katopodes, N. D. and Bradford, S. F. (1999). Mechanics of overland flow, Proceedings of the
1421 International Workshop on Numerical Modelling of Hydrodynamic Systems, Zaragoza, Spain, 21-24
1422 June 1999.

1423 Keshavarzy, A. and Ball, J. E. (1997). Analysis of the characteristics of rough bed turbulent shear
1424 stresses in an open channel, Stochastic Hydrology and Hydraulics, 11 (3), pp.193-210.

1425 Keskin, M. E. and Agiralioglu, N. (1997). A simplified dynamic model for flood routing in
1426 rectangular channels, Journal of Hydrology, 202, pp.302-314.

1427 Keulegan, G. H. (1938). Laws of turbulent flow in open channels, paper RP1151, Journal of
1428 Research, USA National Bureau of Standards, 21, pp.707-741.

1429 Kim, J., Moin, P. and Moser, R. (1987). Turbulence statistics in fully-developed channel flow at
1430 low Reynolds number, Journal of Fluid Mechanics, 177, pp.133-166.

1431 King, H. W. and Brater, E. F. (1963). Handbook of Hydraulics, 5th Ed., McGraw-Hill Book
1432 Company, New York.

1433 Kirkby M.J. (1978). Hillslope Hydrology, John Wiley & Sons, Chichester, UK, 389 pp.

1434 Kirkby, M. J. (1980). The stream head as a significant geomorphic threshold, in: D.R. Coates and
1435 J.D. Vitek (eds) Threshold in Geomorphology, London: George Allen and Unwin, pp.53-73.

1436 Klemes, V. (1986). Dilletantism in hydrology: Transition or destiny?, Water Resources Research,
1437 22, pp.177-188.

1438 Kline, S. J., Reynolds, W. C., Schraub, F. A. and Runstadler, P. W. (1967). The structure of
1439 turbulent boundary layers, Journal of Fluid Mechanics, 30 (4), pp.741-773.

1440 Kneller, B. and Buckee, C. (2001). The structure and fluid mechanics of turbidity currents: a
1441 review of some recent studies and their geological implications, Sedimentology, 47 (Suppl.1), pp.62-
1442 94.

1443 Knisel, W.G. (1980). Creams, a field scale model for chemicals, runoff and erosion from
1444 agricultural management systems. U. C. R. Report USDA no26.

Mis en forme : Espace Avant : 0 pt,
Après : 0 pt

- 1445 | Koomey, J., Berard, S., Sanchez, M. and Wong, H. (2010). Implications of historical trends in the
1446 | electrical efficiency of computing, *IEEE Annals of the history of computing*, 33 (3), pp.46-54.
- 1447 | Koussis, A. D. (1978). Theoretical estimation of flood routing parameters, *Journal of the Hydraulic*
1448 | *Division - ASCE*, 104, HY1, pp.109-115.
- 1449 | Kramer, C. and Papanicolaou, A. (2005) The Effects of Relative Submergence on Cluster
1450 | Formation in Gravel Bed Streams. *Impacts of Global Climate Change*: pp. 1-12.
- 1451 | Kuchment, L.S. (1972). *Matematicheskoye modelirovaniye rechnogo stoka (Mathematical Models*
1452 | *of River Flow)*. Gidrometeoizdat, Leningrad, 190 pp. (in Russian).
- 1453 | Laflen, J. M., Lane, L. J. and Foster, G. R. (1991). A new generation in erosion-prediction
1454 | technology, *Journal of Soil and Water Conservation*, 46, pp.34-38.
- 1455 | Lagrée, P.-Y. (2003). A triple-deck model of ripple formation and evolution, *Physics of Fluids*, 15
1456 | (8), pp.2355-2368.
- 1457 | Lajeunesse, E., Malverti, L., Lancien, P., Armstrong, L., Métivier, F., Coleman, S., Smith, C. E.,
1458 | Davies, T., Cantelli, A. and Parker, G. (2010). Fluvial and submarine morphodynamics of laminar and
1459 | near-laminar flows, *Sedimentology*, 57, pp.1-26.
- 1460 | Lamarre, H. and Roy, A. (2008). The role of morphology on the displacement of particles in a step-
1461 | pool river system, *Geomorphology*, 99, pp.270-279.
- 1462 | Lamb, M. P., Dietrich, W. E. and Venditti, J. G. (2008a). Is the critical Shields stress for incipient
1463 | sediment dependent on channel-bed slope?, *Journal of Geophysical Research-Earth Surface*, 113,
1464 | F02008.
- 1465 | Lamb, M. P., Dietrich, W. E. and Sklar, L. S. (2008b). A model for fluvial bedrock incision by
1466 | impacting suspended and bed load sediment, *Journal of Geophysical Research*, 113, F03025.
- 1467 | Lane, S. N., Richards, K. S. and Chandler, J. H. (1994). Application of distributed sensitivity
1468 | analysis to a model of turbulent open-channel flow in a natural river channel, *Proceedings of the Royal*
1469 | *Society of London Series A- Mathematical Physical and Engineering Sciences*, 447 (1929), pp.49-63.
- 1470 | Langhaar, H. L. (1951). *Dimensional Analysis and the Theory of Models*, New York: Wiley,
1471 | 166 pp.

1472 Lawless, M. and Robert, A. (2001). Scales of boundary resistance in coarse-grained channels,
1473 turbulent velocity profiles and implications, *Geomorphology*, 39 (3-4), pp.221-238.

1474 Lawrence, D. S. L. (1997). Macroscale surface roughness and frictional resistance in overland
1475 flow, *Earth Surface Processes and Landforms*, 22 (4), pp.365-382.

1476 Lawrence, D. S. L. (2000). Hydraulic resistance in overland flow during partial and marginal
1477 surface inundation: experimental observations and modeling, *Water Resources Research*, 36 (8),
1478 pp.2381-2393.

1479 Leonard, A. (1974). Energy cascade in large-eddy simulation of turbulent channel flow, *Advances*
1480 *in Geophysics*, 18, pp.237-248.

Mis en forme : Espace Avant : 0 pt,
Après : 0 pt

Mis en forme : Anglais (États-Unis)

1481 Leopold, L. B., Bagnold, R. A., Wolman, M. G. and Brush, L. M. Jr, (1960). Flow resistance in
1482 sinuous or irregular channels, U.S. Geological Survey Professional Paper 282-C, 134 pp.

1483 Liggett, J. A. and Woolhiser, D. A. (1967). Difference solutions of the shallow-water equation,
1484 *Journal of the Engineering Mechanics Division, ASCE-American Society of Civil Engineers*, 93,
1485 pp.39-71.

1486 Lighthill, M. J. and Whitham, G. B. (1955). On kinematic waves, 1. Flood movement in long
1487 rivers, *Proceedings of the Royal Society – Series A*, 229, pp.281-316.

1488 Lilburne, L. (2002). The scale matcher: A framework for assessing scale compatibility of
1489 environmental data and models, PhD Thesis, University of Otago, Dunedin, New Zealand, 392 pp.

1490 Lobkovsky, A. E., Orpe, A. V., Molloy, R., Kudrolli, A. and Rothman, D. H. (2008). Erosion of a
1491 granular bed driven by laminar fluid flow, *Journal of Fluid Mechanics*, 605, pp.47-58.

1492 Loucks, D. P. and van Beek, E. (2005). *Water Resources Systems Planning and Management - An*
1493 *Introduction to Methods, Models and Applications*, Studies and Reports in Hydrology series,
1494 UNESCO Publishing / WL - Delft Hydraulics, 680 pp.

1495 Lyn, D. A. (1992). Turbulence characteristics of sediment-laden flows in open channels, *Journal of*
1496 *Hydraulic Engineering*, 118 (7), pp.971-988.

1497 Manes, C., Pokrajac, D., Coceal, O. and McEwan, I. (2008). On the significance of form-induced
1498 stress in rough wall turbulent boundary layers, *Acta Geophysica*, 56 (3), pp.845-861.

1499 Mangeney, A., Bouchut, F., Thomas, N., Vilotte, J. P. and Bristeau, M. O. (2007). Numerical
1500 modelling of self-channeling granular flows and of their levee-channel deposits, Journal of
1501 Geophysical Research, 112, F02017.

1502 Manning, R. (1871). On the flow of water in open channels and pipes, Transactions of the
1503 Institution of Civil Engineers of Ireland, 20, pp.161-207.

1504 Marche, F. (2007). Derivation of a new two-dimensional viscous shallow water model with varying
1505 topography, bottom friction and capillary effects, European Journal of Mechanics B/Fluids, 26 (1),
1506 pp.49-63.

1507 [Mavriplis, D. \(1998\). On Convergence Acceleration Techniques for Unstructured Meshes.](#)
1508 [Technical Report ICASE No. 98-44 and NASA/CR-1998-208732, Institute for Computer Applications](#)
1509 [in Science and Engineering, NASA Langley Research Center.](#)

Mis en forme : Anglais
(Royaume-Uni)

Mis en forme : Anglais (États-Unis)

1510 Meile, T. (2007). Influence of macro-roughness of walls on steady and unsteady flow in a channel,
1511 PhD Thesis, Ecole Polytechnique Fédérale de Lausanne, 414 pp.

1512 Mendoza, C. and Zhou, D. (1997). Energetics of sediment-laden streamflows, Water Resources
1513 Research, 33(1), pp.227-234.

1514 Merritt, W.S., Letche, R.A. and Jakeman, A.J. (2003). A review of erosion and sediment transport
1515 models, Environmental Modelling and Software, 18, 761-799.

1516 Métivier, F. and Meunier, P. (2003). Input and output flux correlations in an experimental braided
1517 stream: implications on the dynamics of the bed load transport, Journal of Hydrology, 271, pp.22-38.

1518 Meyer-Peter, E. and Müller, R. (1948). Formulas for bed-load transport, Proceedings of the Second
1519 Meeting of IAHR, Stockholm, pp.39-64.

1520 Milliman, J.D. and Syvitski, J.P.M. (1992). Geomorphic/tectonic control of sediment discharge to
1521 the ocean: the importance of small mountainous rivers, Journal of Geology, pp.525-544.

1522 [Moore, G. E. \(1965\). Cramming More Components onto Integrated Circuits, Electronics, pp. 114–](#)
1523 [117.](#)

Mis en forme : Anglais (États-Unis)

Mis en forme : Anglais (États-Unis)

Mis en forme : Anglais (États-Unis)

1524 Montgomery, D. R. and Buffington, J. M. (1997). Channel-reach morphology in mountain drainage
1525 basins, Geological Society of America Bulletin, 109 (5), pp.596-611.

- 1526 Morgan, R. P. C., Quinton, J. N., Smith, R. E., Govers, G., Poesen, J., Auerwald, K., Chisci, G.,
1527 Torri, D. and Styczen, M. E. (1998). The European Soil Erosion Model (EUROSEM): a dynamic
1528 approach for predicting sediment transport from fields and small catchments, *Earth Surface Processes*
1529 *and Landforms*, 23, pp.527-544.
- 1530 Moussa, R. (1996). Analytical Hayami solution for the diffusive wave flood routing problem with
1531 lateral inflow, *Hydrological Processes*, 10 (9), pp.1209-1227.
- 1532 Moussa, R. and Bocquillon, C. (1996a). Criteria for the choice of flood-routing methods in natural
1533 channels, *Journal of Hydrology*, 186 (1-4), pp.1-30.
- 1534 Moussa, R. and Bocquillon, C. (1996b). Algorithms for solving the diffusive wave flood routing
1535 equation, *Hydrological Processes*, 10 (1), pp.105-124.
- 1536 Moussa, R. and Bocquillon, C. (2000). Approximation zones of the Saint-Venant equations for
1537 flood routing with overbank flow, *Hydrology and Earth System Sciences*, 4(2), pp.251-261.
- 1538 Moussa, R. and Bocquillon, C. (2009). On the use of the diffusive wave for modelling extreme
1539 flood events with overbank flow in the floodplain, *Journal of Hydrology*, 374, pp.116-135.
- 1540 Moussa, R., Voltz, M. and Andrieux, P. (2002). Effects of the spatial organization of agricultural
1541 management on the hydrological behaviour of a farmed catchment during flood events, *Hydrological*
1542 *Processes*, 16, pp.393-412.
- 1543 Moussa, R., Chahinian, N. and Bocquillon, C. (2007). Distributed hydrological modelling of a
1544 Mediterranean mountainous catchment - model construction and multi-site validation, *Journal of*
1545 *Hydrology*, 337, pp.35-51.
- 1546 Mügler, C., Planchon, O., Patin, J., Weill, S., Silvera, N., Richard, P. and Mouche, E. (2010).
1547 Comparison of roughness models to simulate overland flow and tracer transport experiments under
1548 simulated rainfall at plot scale, *Journal of Hydrology*, 402, pp.25-40.
- 1549 Mulder, T. and Alexander, J. (2001). The physical character of subaqueous sedimentary density
1550 flows and their deposits, *Sedimentology*, 48, pp.269-299.
- 1551 Munier, S., Litrico, X., Belaud, G. and Malaterre, P.-O. (2008). Distributed approximation of open-
1552 channel flow routing accounting for backwater effects, *Advances in Water Resources*, 31, pp.1590-
1553 1602.

- 1554 Muñoz-Carpena, R., Parsons, J. E. and Gillian, J. W. (1999). Modelling hydrology and sediment
1555 transport in vegetative filter strips, *Journal of Hydrology*, 214, pp.111-129.
- 1556 Myers, T. G. (2003). Unsteady laminar flow over a rough surface, *Journal of Engineering*
1557 *Mathematics*, 46 (2), pp.111-126.
- 1558 Nakagawa, H. and Nezu, I. (1977). Prediction of the contributions to the Reynolds stress from
1559 bursting events in open channel flows, *Journal of Fluid Mechanics*, 80, pp.99–128.
- 1560 Nash, J. E. and Sutcliffe, J. V. (1970). River flow forecasting through conceptual models, part I - A
1561 discussion of principles, *Journal of Hydrology*, 10 (3), pp.282-290.
- 1562 Navier, C. L. M. H. (1822). Mémoire sur les lois du mouvement des fluides, *Mémoires de*
1563 *l'Académie Royale des Sciences de l'Institut de France*, 6, pp.389-440.
- 1564 Navier, C. L. M. H. (1827). Sur les lois de l'équilibre et du mouvement des corps élastiques,
1565 *Mémoires de l'Académie Royale des Sciences de l'Institut de France*, 7, pp.375-393.
- 1566 Neill, C. R (1968). A re-examination of the beginning of movement for coarse granular bed
1567 materials. Report No. 68, Hydraulic Research Station, Wallingford, England.
- 1568 Nelson, J.M., R.L. Shreve, D.C. McLean, and T.G. Drake (1995). Role of near-bed turbulence
1569 structure in bed load transport and bed form mechanics, *Water Resources Research*, 31(8), pp.2071-
1570 2086.
- 1571 Nepf, H. (1999). Drag, turbulence, and diffusion in flow through emergent vegetation, *Water*
1572 *Resources Research*, 35 (2), pp.479-489.
- 1573 Newton, I. (1687). *Philosophiæ Naturalis – Principia Mathematica*, 1st ed., London, UK, 512 pp.
- 1574 Nezu, I. and Nekagawa, H. (1993). *Turbulence in open-channel flows*, Balkema, Rotterdam, The
1575 Netherlands, 286 pp.
- 1576 Nikora, V., and D. Goring (2000), Flow turbulence over fixed and weakly mobile gravel beds,
1577 *Journal of Hydraulic Engineering*, 126 (9), pp.679-690.
- 1578 Nikora, V., Goring, D., McEwan, I. and Griffiths, G. (2001). Spatially-averaged open-channel flow
1579 over a rough bed, *Journal of Hydraulic Engineering-ASCE*, 127 (2), pp.123-133.

1580 Nikora, V., Larned, S., Nikora, N., Debnath, K., Cooper, G. and Reid, M. (2008). Hydraulic
1581 resistance due to aquatic vegetation in small streams: field study, *Journal of Hydraulic Engineering*,
1582 134 (9), pp.1326-1332.

1583 Nino, Y., Lopez, F. and Garcia, M. (2003). Threshold for particle entrainment into suspension,
1584 *Sedimentology*, 50 (2), pp.247-263.

1585 Nord, G. and Esteves, M. (2010). The effect of soil type, meteorological forcing and slope gradient
1586 on the simulation of internal erosion processes at the local scale, *Hydrological Processes*, 24 (13),
1587 pp.1766-1780.

1588 Ouriemi, M., Aussillous, P., Medale, M., Peysson, Y. and Guazzelli, E. (2007). Determination of
1589 the critical Shields number for particle erosion in laminar flow, *Physics of Fluids*, 19, 061706.

1590 Paiva, R. C. D., Collischonn, W. and Buarque, D. C. (2013). Validation of a full hydrodynamic
1591 model for large-scale hydrologic modelling in the Amazon, *Hydrological Processes*, 27, pp.333-346.

1592 Panton, R. L. (1984). *Incompressible Flow*, John Wiley and Sons, New York, 780 pp.

1593 Paola, C., Heller, P. L. and Angevine, C. L. (1992). The large-scale dynamics of grain-size
1594 variation in alluvial basins. I: Theory, *Basin Research*, 4, pp.73-90.

1595 Paola, C., Straub, K., Mohrig, D. and Reinhardt, L. (2009). The “unreasonable effectiveness” of
1596 stratigraphic and geomorphic experiments, *Earth-Science Reviews*, 97, pp.1-43.

1597 Papanicolaou, A. N., Diplas, P., Dancey, C. L. and Balakrishnan, M. (2001). Surface roughness
1598 effects in near-bed turbulence: implications to sediment entrainment, *Journal of Engineering*
1599 *Mechanics*, 127, pp.211-218.

1600 Parker, G. (1976). On the cause and characteristic scales of meandering and braiding in rivers,
1601 *Journal of Fluid Mechanics*, 76, pp.457-480.

1602 Parker, G. (1978a). Self-formed straight rivers with equilibrium banks and mobile bed. Part 1: the
1603 sand-silt river, *Journal of Fluid Mechanics*, 89 (1), pp.109-125.

1604 Parker, G. (1978b). Self-formed straight rivers with equilibrium banks and mobile bed. Part 2: the
1605 gravel river, *Journal of Fluid Mechanics*, 89 (1), pp.127-146.

- 1606 Parker, G. and Coleman, N. L. (1986). Simple model of sediment-laden flows, *Journal of Hydraulic*
1607 *Engineering*, 112(2b), pp.356-375.
- 1608 Parker, G., Fukushima, Y. and Pantin, H. M. (1986). Self-accelerating turbidity currents, *Journal of*
1609 *Fluid Mechanics*, 171, pp.145-181.
- 1610 Parsons, A. J. and Abrahams, A. D. (1992). *Overland flow: hydraulics and erosion mechanics*,
1611 Chapman & Hall, New-York, 438 pp.
- 1612 Parsons, A. J., Wainwright, J., Abraham, A. D. and Simanton, J. R. (1997). Distributed dynamic
1613 modelling of interrill overland flow. *Hydrological Processes*, 11, pp.1833-1859.
- 1614 Parsons, A. J., Brazier, R. E., Wainwright, J. and Powell, M. E. (2003). Scale relationships in
1615 hillslope runoff and erosion, *Earth Surface Processes and Landforms*, 31 (11), pp.1384-1393.
- 1616 Perkins, S. P. and Koussis, A. D. (1996). Stream-aquifer interaction model with diffusive wave
1617 routing, *Journal of Hydraulic Engineering*, 122 (4), pp.210-218.
- 1618 Perumal, M. and Price, R. K. (2013). A fully mass conservative variable parameter McCarthy–
1619 Muskingum method: Theory and verification, *Journal of Hydrology*, 502, pp.89–102.
- 1620 Peyras, L., Royet, P. and Degoutte, G. (1992). Flow and energy dissipation over stepped gabion
1621 weirs, *Journal of Hydraulic Engineering – ASCE*, 118 (5), pp.707-717.
- 1622 Pickup, G. and Marks, A. (2001). Regional scale sedimentation process models from airborne
1623 gamma ray remote sensing and digital elevation data, *Earth Surface Processes and Landforms*, 26,
1624 pp.273-293.
- 1625 Pokrajac, D., Campbell, L. J., Nikora, V., Manes, C. and McEwan, I. (2007). Quadrant analysis of
1626 persistent spatial velocity perturbations over square-bar roughness, *Experiments in Fluids*, 42 (3),
1627 pp.413-423.
- 1628 Polyakov, V. O. and Nearing, M. A. (2003). Sediment transport in rill flow under deposition and
1629 detachment conditions, *Catena*, 51 (1), pp.33-43.
- 1630 Ponce, V. M. (1990). Generalized diffusive wave equation with inertial effects, *Water Resources*
1631 *Research*, 26 (5), pp.1099-1101.

- 1632 Ponce, V. M. and Simons, D. B. (1977). Shallow wave propagation in open channel flow, Journal
1633 of the Hydraulic Division, American Society of Civil Engineers, 103(HY12), pp.1461-1476.
- 1634 Ponce, V. M., Li, R. M. and Simons, D. B. (1978). Applicability of kinematic and diffusion
1635 models, Journal of the Hydraulic Division, American Society of Civil Engineers, 104 (HY3), pp.353-
1636 360.
- 1637 Ponce, V. M., Lohani, A. K. and Scheyhing, C. (1996). Analytical verification of Muskingum-
1638 Cunge routing, Journal of Hydrology, 174, pp.235-241.
- 1639 Powell, D. M. (2014). Flow resistance in gravel-bed rivers: progress in research, Earth Science
1640 Reviews, 136, pp.301-338.
- 1641 Prah, L., Holzer, A., Arlov, D., Revstedt, J., Sommerfeld, M. and Fuchs, L. (2007), On the
1642 interaction between two fixed spherical particles, International Journal of Multiphase Flow, 33 (7),
1643 pp.707-725.
- 1644 Priezjev, N. and Troian, S. (2006). Influence of periodic wall roughness on the slip behaviour at
1645 liquid/solid interfaces: molecular-scale simulations versus continuum predictions, Journal of Fluid
1646 Mechanics, 554, pp.25-46.
- 1647 Prosser, I. P. and Rustomji, P. (2000). Sediment transport capacity for overland flow, Progress in
1648 Physical Geography, 24 (2), pp.179-193.
- 1649 Prosser, I. P., Dietrich, W. E. and Stevenson, J. (1995). Flow resistance and sediment transport by
1650 concentrated overland flow in a grassland valley, Geomorphology, 13, pp.71-86.
- 1651 Rathburn, S. and Wohl, E. (2003). Predicting sediment dynamics along a pool-riffle mountain
1652 channel, Geomorphology, 55, pp.111-124.
- 1653 Raupach, M. R. (1981). Conditional statistics of Reynolds stress in rough-wall and smooth-wall
1654 turbulent boundary layers, Journal of Fluid Mechanics, 108, pp.363-382.
- 1655 Rauws, G. (1980). Laboratory experiments on resistance to overland flow due to composite
1656 roughness, Journal of Hydrology, 103, 37-52.
- 1657 | Rayleigh (1877). Theory of sound, Macmillan, London, [UK](#).

1658 Reddy, K. V., Eldho, T. I., Rao, E. P. and Hengade, N. (2007). A kinematic-wave-based distributed
1659 watershed model using FEM, GIS and remotely sensed data, *Hydrological Processes*, 21 (20),
1660 pp.2765–2777.

1661 Reynolds, O. (1883). An experimental investigation of the circumstances which determine whether
1662 the motion of water shall be direct or sinuous, and of the law of resistance in parallel channels.
1663 *Philosophical Transactions of the Royal Society of London (A)*, 174, pp.935–982.

1664 Reynolds, O. (1895). On the dynamical theory of incompressible viscous fluids and the
1665 determination of the criterion, *Philosophical Transactions of the Royal Society of London (A)*, 186,
1666 pp.123-164.

1667 Riabouchinsky, D. (1911). Méthode des variables de dimension zero et son application en
1668 aérodynamique, *L'aérophile*, 19, pp.407-408.

1669 Richardson, S. (1973). On the no-slip boundary condition, *Journal of Fluid Mechanics*, 59 (4),
1670 pp.707-719.

1671 Ritchie, J.C. and McHenry, J.R (1990). Application of Radioactive Fallout Cesium-137 for
1672 Measuring soil erosion and sediment accumulation rates and patterns: A review. *Journal of*
1673 *Environmental Quality*, 19, pp.215-233.

1674 Roche, N. (2006). Modélisation du ruissellement sur surfaces rugueuses, PhD Thesis, Université
1675 Joseph Fourier, Grenoble, France, 213 pp.

1676 Rodellar, J., Gomez, M. and Bonet, L. (1993). Control method for on-demand operation of open-
1677 channel flow, *Journal of Irrigation and Drainage Engineering*, 119, pp.225-241.

1678 Rödi, W. (1988). [Turbulence models and their application in hydraulics - a state of the art review,](#)
1679 [International Association for Hydraulic Research, Delft, The Netherlands, 47p.](#)

1680 Romanowicz, R.J., Dooge, J.C.I., and Kundzewicz, Z.W. (1988). Moments and cumulants of
1681 linearized St. Venant equation, *Advances in Water Resources*, 11, 92-100.

1682 Rosgen, D. L. (1994). A classification of natural rivers, *Catena*, 22 (3), pp.169-199.

1683 Roux, H. and Dartus, D. (2006). Use of parameter optimization to estimate a flood wave: Potential
1684 applications to remote sensing of rivers, *Journal of Hydrology*, 328, pp.258– 266.

Mis en forme : Espace Avant : 0 pt,
Après : 0 pt

Mis en forme : Anglais

1685 Rutschmann, P. and Hager, W. H. (1996). Diffusion of floodwaves, *Journal of Hydrology*, 178,
1686 pp.19-32.

1687 Saint-Venant, A. J.-C. B. de (1871). Théorie du mouvement non permanent des eaux, avec
1688 application aux crues des rivières et à l'introduction des marées dans leur lit, *Comptes-Rendus de*
1689 *l'Académie des Sciences*, 73, pp.147-154 and 237-240.

1690 Saleh, F, Ducharme, A., Flipo, N., Oudin, L. and Ledoux, E. (2013). Impact of river bed
1691 morphology on discharge and water levels simulated by a 1D Saint-Venant hydraulic model at
1692 regional scale, *Journal of Hydrology*, 476, pp.169–177.

1693 Sander, G. C., Parlange, J. Y., Barry, D. A., Parlange, M.B. and Hogarth, W. L. (2007). Limitation
1694 of the transport capacity approach in sediment transport modeling, *Water Resources Research*, 43 (2),
1695 W02403, doi: [10.1029/2006WR005177](https://doi.org/10.1029/2006WR005177).

1696 Sau, J., Malaterre, P.-O. and Baume, J.-P. (2010). Sequential Monte-Carlo state estimation of an
1697 irrigation canal, *Comptes-Rendus Mécanique*, 338, pp.212-219.

1698 Savage, S. B. and Hutter, K. (1989). The motion of a finite mass of granular material down a rough
1699 incline, *Journal of Fluid Mechanics*, 199, pp.177-215.

1700

1701 Savage, S. B. and Hutter, K. (1991). The dynamics of avalanches of granular materials from
1702 initiation to runout. Part I: Analysis, *Acta Mechanica*, 86, pp.201-223.

1703 Schlichting H. (1936). Experimentelle Untersuchungen zum Rauigkeitsproblem *Ing.Arch.* 7:1–34.
1704 (Engl. transl. 1937. Experimental investigation of the problem of surface roughness. NACA TM 823)

1705 Schmeeckle, M. W. and Nelson, J. M. (2003). Direct numerical simulation of bedload transport
1706 using a local, dynamic boundary condition, *Sedimentology*, 50 (2), pp.279-301.

1707 Schmeeckle, M. W., Nelson, J. M. and Shreve, R. L. (2007). Forces on stationary particles in near-
1708 bed turbulent flows, *Journal of Geophysical Research-Earth Surface*, 112 (F2), F02003, doi:
1709 [10.1029/2006JF00536](https://doi.org/10.1029/2006JF00536).

1710 Sear, D. A. (1996). Sediment transport processes in pool-riffle sequences, *Earth Surface Processes*
1711 *and Landforms*, 21, pp.241-262.

1712 Sen, D. J. and Garg, N. K. (2002). Efficient algorithm for generally varied flows in channel
1713 networks, *Journal of Irrigation and Drainage Engineering*, 128 (6), pp.351-357.

- 1714 Sharifi, S., Sterling, M. and Knight, D. W. (2009). A novel application of a multi-objective
1715 evolutionary algorithm in open channel flow modelling, *Journal of Hydroinformatics*, 11(1), pp.31-50.
- 1716 Sheets, B., Hickson, T.A., and Paola, C. (2002). Assembling the stratigraphic record: Depositional
1717 patterns and time-scales in an experimental alluvial basin, *Basin Research*, 14, pp.287–301.
- 1718 Shields, A. (1936). Anwendung der Ähnlichkeitsmechanik und der Turbulenzforschung auf die
1719 Geschiebebewegung, Mitteilungen der Preußischen Versuchsanstalt für Wasserbau und Schiffbau
1720 No.6, Berlin, Germany. (English translation by W. P. Ott and J. C. van Uchelen, Hydrodynamics
1721 Laboratory Publication No.167, Hydrodynamics Laboratory of the California Institute of Technology,
1722 Pasadena, USA).
- 1723 Simpson, G. and Castellort, S. (2006). Coupled model of surface water flow, sediment transport
1724 and morphological evolution, *Computers & Geosciences*, 32 (10), pp.1600-1614.
- 1725 Sivakumaran, N. S. and Yevjevich, V. (1987). Experimental verification of the Dressler curved-
1726 flow equations, *Journal of Hydraulic Research*, 25 (3), pp.373-391.
- 1727 Sivapalan, M., Bates, B. C. and Larsen, J. E. (1997). A generalized, non-linear, diffusion wave
1728 equation: theoretical development and application, *Journal of Hydrology*, 192, pp.1-16.
- 1729 Sklar, L. S. and Dietrich, W. E. (2004). A mechanistic model for river incision into bedrock by
1730 saltating bed load, *Water Resources Research*, 40, W06301, doi: [10.1029/2003WR002496](https://doi.org/10.1029/2003WR002496).
- 1731 Slaymaker, O. (2006). Towards the identification of scaling relations in drainage basin sediment
1732 budgets, *Geomorphology*, 80, pp.8–19.
- 1733 Smagorinsky, J. (1963). General circulation experiments with the primitive equations, *Monthly*
1734 *Weather Review*, 91 (3), pp.99-164.
- 1735 Smart, G. M. (1984). Sediment transport formula for steep channels, *Journal of Hydraulic*
1736 *Engineering*, 110 (3), pp.267-276.
- 1737 Smith, J. D. and McLean S. R. (1977). Spatially averaged flow over a wavy surface, *Journal of*
1738 *Geophysical Research*, 82, pp.1735-1746.
- 1739 Smith, M. W., Cox, N. J. and Bracken, L. J. (2007). Applying flow resistance equations to overland
1740 flows, *Progress in Physical Geography*, 31, pp.363-387.

Mis en forme : Espace Avant : 0 pt,
Après : 0 pt

Mis en forme : Anglais
(Royaume-Uni)

1741 Smith, R. E., Goodrich, D. C., Woolhiser, D. A. and Unkrich, C. L. (1995). KINEROS – a
1742 kinematic runoff and erosion model. In: Computer Models of Watershed Hydrology, Singh, V. P.
1743 (ed.), Water Resources: Littleton, CO, pp.697-732.

1744 Stein, O.R., Alonso, C.V. and Julien, P.Y. (1993). Mechanics of jet scour downstream of a headcut,
1745 Journal of Hydraulic Research, 31 (6), pp.723-738.

1746 Stevenson, P., Thorpe, R. B. and Davidson, J. F. (2002). Incipient motion of a small particle in the
1747 viscous boundary layer at a pipe wall, Chemical Engineering Science, 57 (21), pp.4505-4520.

1748 Stoker, J. J. (1957~~8~~). Water waves, the mathematical theory with application, Wiley, [Interscience](#)
1749 [Publishers, New York, USA, 357~~600~~ pp.](#)

1750 Stokes, G. G. (1845). On the theories of internal friction of fluids in motion, Transactions of the
1751 Cambridge Philosophical Society, 8, pp.287-319.

1752 Strahler, A.N. (1956). The nature of induced erosion and aggradation. In: Thomas, W.L. (Ed.),
1753 Man's Role in Changing the Face of the Earth. University of Chicago Press, Chicago, pp.621–638.

1754 Sundaresan, S., Eaton, J., Koch, D. L. and Ottino, J. M. (2003). Appendix 2: Report of study group
1755 on disperse flow, International Journal of Multiphase Flow, 29, pp.1069-1087.

1756 Sutherland, A.J. (1967), Proposed mechanism for sediment entrainment by turbulent flows, *Journal*
1757 *of Geophysical Research*, 72, pp.6183-6194.

1758 Syvitski, J.P.M. and Milliman, J.D. (2007). Geology, geography, and humans battle for dominance
1759 over the delivery of fluvial sediment to the coastal ocean, *Journal of Geology*, 115 (1), pp.1-19.

1760 Tatard, L., Planchon, O., Wainwright, J., Nord, G., Favis-Mortlock, D., Silvera, N., Ribolzi, O.,
1761 Esteves, M., Huang, C. H. (2008). Measurement and modeling of high-resolution flow-velocity data
1762 under simulated rainfall on a low-slope sandy soil, *Journal of Hydrology*, 348 (1-2), pp1-12.

1763 Tiemeyer, B., Moussa, R., Lennartz, B. and Voltz, M. (2007). MHYDAS-DRAIN : a spatially
1764 distributed model for small, artificially drained lowland catchments, *Ecological Modelling*, 209 (1),
1765 pp.2-20.

1766 Todini, E. and Bossi, A. (1986). PAB (Parabolic and Backwater) : an unconditionnally stable flood
1767 routing scheme particularly suited for real time forecasting and control, *Journal of Hydraulic*
1768 *Research*, 24 (5), pp.405-424.

- 1769 Trigg, M. A., Wilson, M. D., Bates, P. D., Horritt, M. S., Alsdorf, D. E., Forsberg, B. R. and Vega,
1770 M. C. (2009). Amazon flood wave hydraulics, *Journal of Hydrology*, 374 (1-2), pp.92-105.
- 1771 Turek, S. (1999). *Efficient Solvers for Incompressible Flow Problems*. Springer, Berlin, Germany,
1772 352 pp.
- 1773 Van Heijst, M. W. I. M., Postma, G., Meijer, X. D., Snow, J. N. and Anderson, J. B. (2001).
1774 Quantitative analogue flume-model study of River-shelf systems: principles and verification
1775 exemplified by the late quaternary Colorado River-delta evolution, *Basin Research*, 13, pp.243-268.
- 1776 Van Maren, D. S. (2007). Grain size and sediment concentration effects on channel patterns of silt-
1777 laden rivers, *Sedimentary Geology*, 202, pp.297-316.
- 1778 Van Rijn, L. C. (1984a). Sediment transport, part I: bed load transport, *Journal of Hydraulic*
1779 *Engineering*, 110, pp.1431-1456.
- 1780 Van Rijn, L. C. (1984b). Sediment transport, part II: suspended load transport, *Journal of Hydraulic*
1781 *Engineering*, 110, pp.1613-1641.
- 1782 Vanoni, V.A. (1946). Transportation of suspended sediment by water, *Transactions of the A.S.C.E.*,
1783 111, pp.67-133.
- 1784 Vaschy, A. (1892). Sur les lois de similitude en physique, *Annales Télégraphiques*, 19, pp.25-28.
- 1785 Vieux, B. E., Cui, Z. and Gaur, A. (2004). Evaluation of a physics-based distributed hydrologic
1786 model for flood forecasting, *Journal of Hydrology*, 298, pp.155-177.
- 1787 Wainwright, J., Parsons, A. J., Müller, E. N., Brazier, R. E., Powell, D. M. and Fenti, B. (2008). A
1788 transport-distance approach to scaling erosion rates: I. Background and model development, *Earth*
1789 *Surface Processes and Landforms*, 33 (5), pp.813-826.
- 1790 Walling, D.E. (1983). The sediment delivery problem, *Journal of Hydrology*, 65, pp.209-237.
- 1791 Wang, G. T and Chen, S. (2003). A semi-analytical solution of the Saint-Venant equations for
1792 channel flood routing, *Water Resources Research*, 39 (4), 1076, doi:10.1029/2002WR001690.
- 1793 Wang, G. T., Yao, C., Okoren, C. and Chen, S. (2006). 4-Point FDF of Muskingum method based
1794 on the complete St Venant equations, *Journal of Hydrology*, 324, pp.339-349.

- 1795 Wang, L, Wu, J. Q., Elliot, W. J., Fiedler, F. R. and Lapin, S. (2014). Linear diffusion wave
1796 channel routing using a discrete Hayami convolution method, *Journal of Hydrology*, 509, pp.282-294.
- 1797 Wang, Y, Straub, K. M. and Hajek, E. A. (2011). Scale-dependent compensational stacking: an
1798 estimate of time scales in channelized sedimentary deposits, *Geology*, 39 (9), pp.811-814.
- 1799 Weichert, R. (2006). Bed morphology and stability of steep open channels, PhD Thesis, Technische
1800 Hochschule, Zürich, 265 pp.
- 1801 Weisbach, J. (1845). *Lehrbuch der Ingenieur- und Maschinen-Mechanik*, Vieweg und Sohn eds.,
1802 Braunschweig.
- 1803 Whitham, G. B. (1999). *Linear and nonlinear waves*. John Wiley & Sons Inc., New York.
- 1804 Wiberg, P. L. and Smith, J. D. (1987). Calculations of the critical shear stress for motion of
1805 uniform and heterogeneous sediments. *Water Resources Research*, 23, pp.1471–1480.
- 1806 Williams G. P. (1970). Flume width and water depth effects in sediment-transport experiments. US
1807 Geological Survey Professional Paper 562-H.
- 1808 Wright, S. and Parker, G. (2004). [Flow Resistance and Suspended Load in Sand-Bed Rivers:
1809 Simplified Stratification Model](#), *Journal of Hydraulic Engineering*, 130 (8), pp.796-805.
- 1810 Wu, R. M. and Lee, D. J. (2001). Hydrodynamic drag on non-spherical floc and free-settling test,
1811 *Water Research*, 35 (13), pp.3226-3234.
- 1812 Yager, E. M., Kirchner, J. W. and Dietrich, W. E. (2007). Calculating bed load transport in steep
1813 boulder bed channels, *Water Resources Research*, 43, W07418, doi: [10.1029/2006WR005432](#).
- 1814 Yalin, M. S. (1977). *Mechanics of sediment transport*. Pergamon Press, Oxford, UK, 2nd ed,
1815 298 pp.
- 1816 Yang, C. T. (1974). Unit stream power and sediment transport, *Journal of the Hydraulics Division-
1817 ASCE*, 100 (NHY9), pp.1269-1272.
- 1818 Zanke, U. C. E. (2003). On the influence of turbulence on the initiation of sediment motion,
1819 *International Journal of Sediment Research*, 18 (1), pp.17-31.

1820 Zhou, J. G. (1995). Velocity-depth coupling in shallow-water flows, *Journal of Hydraulic*
1821 *Engineering*, 121 (10), pp.717-724.

1822 Zimmermann, A. and Church, M. (2001). Channel morphology, gradient profiles and bed stresses
1823 during flood in a step-pool channel, *Geomorphology*, 40 (3-4), pp.311-327.

1824 Zoppou, C. and O'Neill, I. C. (1982). Criteria for the choice of flood routing methods in natural
1825 channels, *Hydrology and Water Resources, Symposium*, 11-13 May 1982, Melbourne, pp.75-81.

1826