



Supplement of

Influence of bank slope on sinuosity-driven hyporheic exchange flow and residence time distribution during a dynamic flood event

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16 This supporting information contains additional information on how the model used in
 17 our simulations was set up in COMSOL. We use the same methods, equations, and
 18 metrics described in Gomez-Velez et al. (2017), however, we implement here a
 19 deformed geometry method to capture the dynamic evolution of the wetting front along
 20 the sloping banks, while Gomez-Velez et al. (2017) assumed a vertical river bank. Due
 21 to their large file size, the COMSOL model files and the raw data to the figures in the
 22 manuscript are available upon request. For this please contact Yiming Li
 23 (liyim@cug.edu.cn) or Zhang Wen (wenz@cug.edu.cn).

24 **S1 Water flow model**

25 The water flow model is based on that of Gomez-Velez et al. (2017), comprising
 26 an alluvial valley with a sinusoidal meandering river that overlies non-permeable river
 27 deposits, as shown in Fig. S1. To simplify the model, aquifer properties are assumed to
 28 be spatially homogeneous and isotropic. This means they can be modeled by the
 29 commonly used vertical-integrated approach which can reduce a 3-D groundwater flow
 30 problem to a two-dimensional (2-D) one, as shown in Fig. S2a. The model is bounded
 31 by hillslopes and a two-period fully penetrating sinusoidal river. By neglecting the
 32 compression of groundwater, unsteady, 2-D transient groundwater flow through the
 33 deformable aquifer is described by the Boussinesq's equation:

$$34 \quad S_y \frac{\partial h}{\partial t} = \nabla [K(h - z_b) \nabla h] \quad (\text{S1a})$$

$$35 \quad h(\mathbf{x}, t = 0) = h_0(\mathbf{x}) \quad (\text{S1b})$$

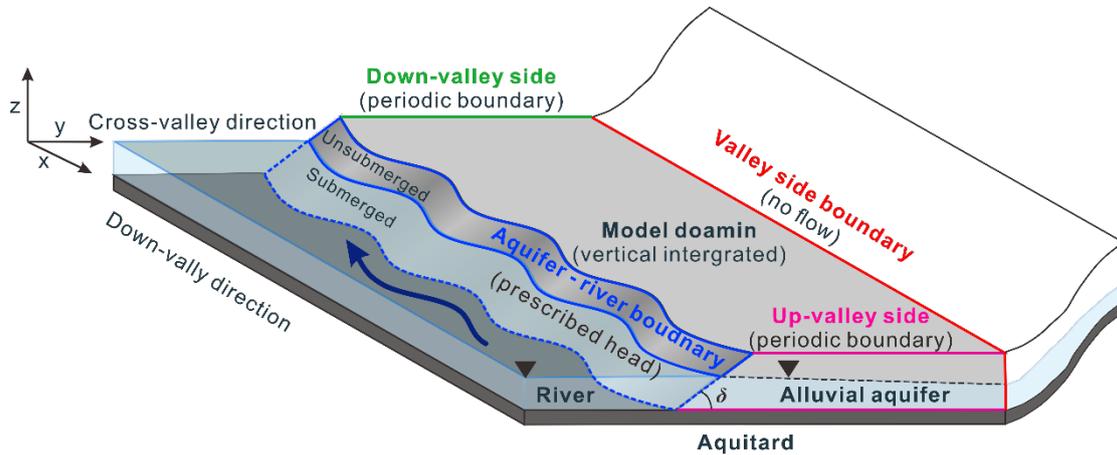
$$36 \quad \mathbf{n} \cdot \nabla [K(h - z_b) \nabla h] = 0 \quad \text{for } \Omega_v \quad (\text{S1c})$$

$$37 \quad h(x_u, y, t) = h(x_d, y, t) + 2 \lambda J_x \quad \text{for } \Omega_u \text{ and } \Omega_d \quad (\text{S1d})$$

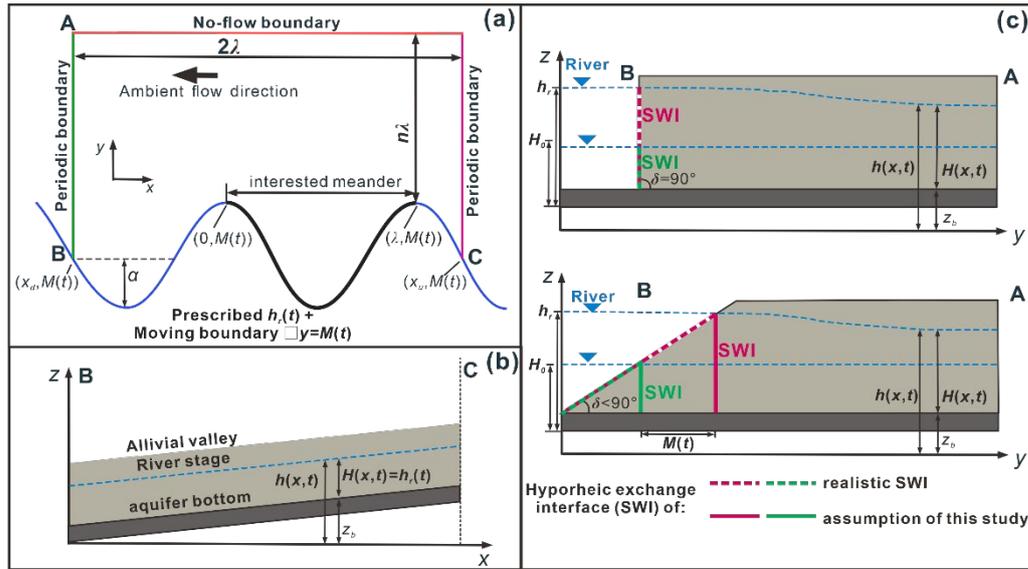
$$38 \quad h(\mathbf{x}, t) = \left(\frac{J_x}{\sigma}\right)s(x) + H_s(t) + 2 \lambda J_x \quad \text{for } \Omega_{in} \cup \Omega_{out} \quad (\text{S1e})$$

39 where $\mathbf{x} = (x, y)$ [L] is the spatial coordinate with x positive in the upstream direction, t
 40 [T] is simulation time, S_y [-] is specific yield, K [LT⁻¹] is the hydraulic conductivity, ∇
 41 is the Laplace operator, $h(\mathbf{x}, t)$ and $h_0(\mathbf{x})$ [L] represent the hydraulic head at t and $t = 0$,

42 while $z_b(\mathbf{x})$ [L] is the elevation of the underlying impermeable layer with respect to the
 43 reference datum $z = 0$ (see Fig. S2b and S2c), respectively. $H(\mathbf{x}, t) = h(\mathbf{x}, t) - z_b(\mathbf{x})$ [L]
 44 is the thickness of the saturated aquifer, \mathbf{n} is the outward normal vector along the model
 45 boundary, Ω_v , Ω_u and Ω_d are the valley, upstream and downstream boundaries,
 46 respectively, while Ω_{in} and Ω_{out} are the inlet and outlet boundaries along the river. The
 47 fluxes are calculated by Darcy's law via $\mathbf{q} = -K\nabla h$ [LT^{-1}]. Here, \mathbf{q} is the specific
 48 discharge or Darcy flux, \mathbf{q}/θ [LT^{-1}] is the pore water velocity with θ [-] as effective
 49 porosity, and $\mathbf{Q} = \mathbf{q}(h - z_b)$ [L^2T^{-1}] is the aquifer-integrated discharge in our 2-D model.
 50 The valley boundary (Ω_v) is assigned as a no-flow boundary and located at $y = n\lambda$, with
 51 the scaling number $n = 4.5$, which has proven to be sufficiently large for this simulation
 52 based on a series of pre-simulation tests while λ [L] is the wavelength of the river
 53 sinusoid. The river has been assigned a known transient hydraulic head, $h_r(x, t) =$
 54 $(J_x/\sigma)s(x) + H_r(t)$ [m], where J_x [-] is the base head gradient of ambient flow along the
 55 valley in positive x direction, $H_r(t)$ [L] is the elevation of river stage above the
 56 impermeable deposit at the downstream end.
 57



58
 59 **Figure S1.** Conceptual model of the study area. Colored lines represent the river, up-
 60 valley, down-valley and valley side boundary conditions set in the model. Modified
 61 from Schmadel et al. (2016)
 62



63

64 **Figure S2.** Modified after Gomez-Velez et al. (2017): (a) Schematic representation of
 65 the boundary conditions for the non-submerged alluvial system. The colors of the
 66 boundaries correspond to those in Fig. S1. (b) Representation of the stream stage
 67 variation along the channel thalweg. (c) Cross-section of unconfined aquifer and
 68 floodplain of vertical ($\delta = 90^\circ$) and sloping riverbank ($\delta < 90^\circ$). Green and red lines
 69 refer to the sediment-water interface (SWI) during base flow condition and flood event,
 70 respectively; the dashed lines on the riverbank surface and the vertical bold lines in Fig.
 71 S2c indicate the realistic SWIs and SWIs of this study, respectively.

72

73 The river ($\Omega_{in} \cup \Omega_{out}$) is implemented as a sinusoid, following the
 74 conceptualizations of Boano et al. (2006), Cardenas (2009a, 2009b), and Gomez-Velez
 75 et al. (2017). The initial condition is represented as: $y_0(x) = \alpha \cos(2\pi x/\lambda) - \alpha$, where α [L]
 76 is the amplitude of the river boundary. Left- and right-bottom vertices in initial
 77 condition are located at $x_d = -3\lambda/4$ and $x_u = 5\lambda/4$, respectively.

78

79 The impermeable bottom deposit $z_b = J_x(x - x_d)$ [L] is assumed to be parallel to
 80 the alluvial valley. Ω_u and Ω_d are periodic with a known variable hydraulic head drop
 81 $h(x = x_u, y, t) = h(x = x_d, y, t) + 2\lambda J_x$ [L] to eliminate any boundary effects. Thus, the
 82 model domain can represent two periodic parts in horizontal direction of the infinite
 aquifer. The river stage fluctuates during the dynamic flood event following (Cooper

83 and Rorabaugh, 1963):

$$84 \quad H_s = \begin{cases} H_0 + H_p \exp \left\{ -\eta(t-t_p) \frac{[1 - \cos(\omega t)]}{[1 - \cos(\omega t_p)]} \right\} & \text{if } 0 < t < t_p \\ H_0 & \text{if } t_p < t \end{cases} \quad (\text{S2})$$

85 where $H_0(\mathbf{x})$ [L] is the initial river stage, H_p [L] is the maximum (peak) river stage
 86 during the flood event, while t_d and t_p [T] are the duration of flood event and the time-
 87 to-peak river stage, respectively. $\omega = 2\pi/t_d$ [T⁻¹] is the flood event frequency, $\eta =$
 88 $\omega \cot(\omega t_p/2)$ [T⁻¹] represents the degree of flood event asymmetry. The peak river stage
 89 and time-to-peak are assumed to be linearly correlated with the base flow stage ($H_p =$
 90 $n_0 H_0$) and the duration of the event ($t_p = n_d t_d$), respectively. Constants n_0 [-] and n_d [-]
 91 represent river stage hydrograph intensity and skewness.

92

93 **S2 Conservative solute transport model and calculation of HZ area (extent)**

94 In this work, we adopt the mathematical model used by Gomez-Velez et al. (2017),
 95 where the transport of a conservative solute within the vertically integrated system is
 96 given by:

$$97 \quad \frac{\partial(H\theta C)}{\partial t} = \nabla \cdot (\mathbf{D}\nabla C - \mathbf{Q}C) \quad (\text{S3a})$$

$$98 \quad C(\mathbf{x}, t=0) = C_0(\mathbf{x}) \quad (\text{S3b})$$

$$99 \quad \mathbf{n} \cdot (\mathbf{Q}C - \mathbf{D}\nabla C) = 0 \text{ for } \Omega_v \quad (\text{S3c})$$

$$100 \quad C(x_u, y, t) = C(x_d, y, t) \text{ for } \Omega_u \text{ and } \Omega_d \quad (\text{S3d})$$

$$101 \quad C(\mathbf{x}, t) = C_s(\mathbf{x}, t) \text{ for } \Omega_{in} \quad (\text{S3e})$$

$$102 \quad \mathbf{n} \cdot (\mathbf{Q}C - \mathbf{D}\nabla C) = 0 \text{ for } \Omega_{out} \quad (\text{S3f})$$

103 where $C(\mathbf{x}, t)$, $C_0(\mathbf{x})$, and $C_s(\mathbf{x}, t)$ are the solute concentrations [ML⁻³] in the aquifer,
 104 initial concentration and concentration in the river, respectively. The dispersion-
 105 diffusion tensor $\mathbf{D} = \{D_{ij}\}$ [L²T⁻¹] is defined according to Bear and Cheng (2010) as:

106
$$D_{ij} = \alpha_T |\mathbf{Q}| \delta_{ij} + (\alpha_L - \alpha_T) \frac{Q_i Q_j}{|\mathbf{Q}|} + H\theta \epsilon D_L \quad (\text{S4})$$

107 where α_T and α_L [L] are the transverse and longitudinal dispersivity, respectively, D_L
 108 [L^2T^{-1}] is the water diffusivity, $\epsilon = \theta^{1/3}$ [-] represents tortuosity (Millington and Quirk,
 109 1961), and δ_{ij} [-] is the Kronecker delta function.

110 In order to mimic a periodical repetition of the meanders in x direction and
 111 eliminate potential boundary effects, a periodic boundary condition (Eq. (S3d)) is used
 112 at Ω_u and Ω_d . This type of boundary condition can produce the periodic nature of the
 113 model domain, flow field as well as the HZ that repeats for each meander bend (Gomez-
 114 Velez et al., 2017). However, in order to explore the local HZ that is caused by the HEFs
 115 at the studied meander: $0 < x < \lambda$ (i.e., bold black line along the meander in Fig. S2a),
 116 the conservative in-stream concentration is given by

117
$$C_s(t) = \begin{cases} 1 & \text{if } x \in [0, \lambda] \cap \Omega_m \\ 0 & \text{else} \end{cases} \quad (\text{S5})$$

118 According to Eq. (S5), the river concentrations are assigned as an open boundary
 119 condition along the studied meander with the external concentration mimicking the
 120 concentration of the tracer (100% of stream water). Then the concentration of the pore
 121 water within the aquifer represents the fraction of water inflow from the river at any
 122 given location and time.

123 **S3 Model of residence time distribution**

124 The residence time (also terms as travel time or age) (RT) in the HZ describes the
 125 characteristic time scale over which water or solute molecule are exposed to the
 126 biogeochemical conditions within the hyporheic sediment. For HEF process, RT is
 127 controlled by the advective and dispersive characteristics of the system, thereby it is
 128 hard to calculate the RT of each molecule due extremely large computational demand.
 129 Thus, the residence time distribution equation was proposed (Ginn, 2000; Gomez-Velez
 130 et al., 2012.), and has been widely applied to calculate the mean residence time

131 distribution (RTD) in HEF models (Gomez-Velez et al., 2017; Singh et al., 2019).
 132 Similar to Gomez-Velez et al. (2017), here we focus on the orders of moment of RTD
 133 that represent the mean residence time distribution:

$$134 \quad \mu_n(\mathbf{x}, t) = \int_0^\infty \tau^n \rho(\mathbf{x}, t, \tau) d\tau \quad n = 0, 1, \dots \quad (S6)$$

135 here, $\mu_n(\mathbf{x}, t)$ [T^n] is the n -th moment, $\rho(\mathbf{x}, t, \tau)$ [T^{-1}] is the residence time distribution, τ
 136 [T] is residence time, and $\mu_0 = 1$. The first moment of RTD ($\mu_1(\mathbf{x}, t)$) is the mean
 137 residence time distribution at a given location and time, which can be used to evaluate
 138 the transient variation of RTD. Here, we used the approach provided by Gomez-Velez
 139 et al. (2017) where the moments of RTD are calculated by a form of the advection-
 140 dispersion equation following

$$141 \quad \frac{\partial(H\theta\mu_n)}{\partial t} = \nabla \cdot (\mathbf{D}\nabla\mu_n - \mathbf{Q}\mu_n) + nH\theta\mu_{n-1} \quad (S7a)$$

$$142 \quad \mu_n(\mathbf{x}, t=0) = \mu_{n,0}(\mathbf{x}) \quad (S7b)$$

$$143 \quad \mathbf{n} \cdot (\mathbf{Q}\mu_n - \mathbf{D}\nabla\mu_n) = 0 \text{ for } \Omega_v \quad (S7c)$$

$$144 \quad \mu_n(x_u, y, t) = \mu_n(x_d, y, t) \text{ for } \Omega_u \text{ and } \Omega_d \quad (S7d)$$

$$145 \quad \mu_n(\mathbf{x}, t) = 0 \text{ for } \Omega_{in} \quad (S7e)$$

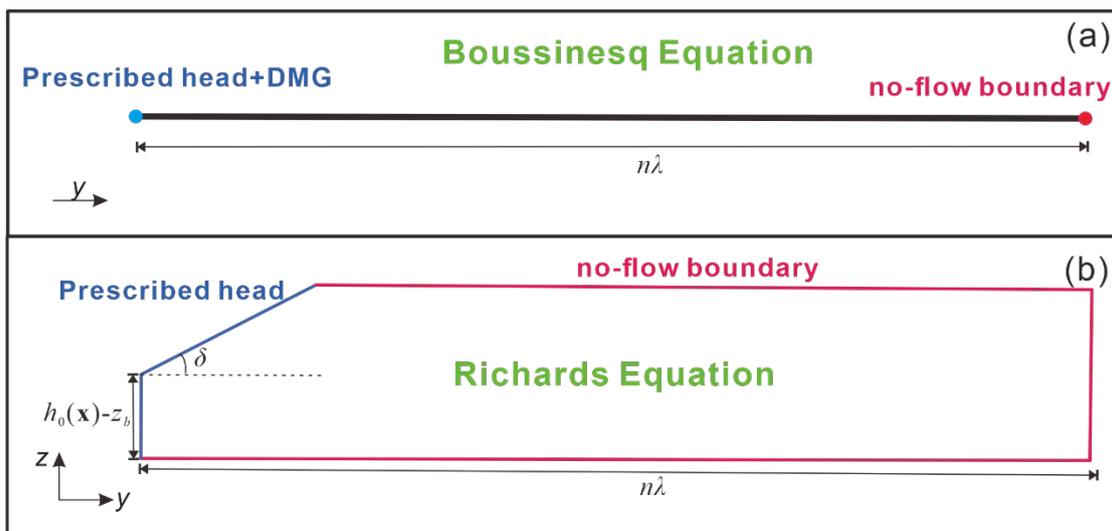
$$146 \quad \mathbf{n} \cdot (\mathbf{Q}\mu_n - \mathbf{D}\nabla\mu_n) = 0 \text{ for } \Omega_{out} \quad (S7f)$$

147 where $\mu_{n,0}(\mathbf{x}, t)$ is the initial condition of the n -th RTD that is calculated by the base
 148 flow condition (steady forcing before the arrival of flood event), while the upstream
 149 and downstream boundaries are assigned periodic boundary conditions (Eq. (S7d). As
 150 we ignore the vadose zone, the RT is defined as the time since the water entered the
 151 model domain from the river (i.e., the travel time of the water). Thus, n -th RTD at the
 152 inflow river boundary is zero (Eq. (S7e). A flow boundary is used for the region where
 153 the water exits the model domain (Eq. (S7f)).

154 **S4 Validation of assumptions and using of DGM**

155 The two main assumptions of this study are: (1) The SWI is always vertical and (2)

156 the DGM can be used to represent the displacement of the SWI. In order to test the
 157 appropriateness of these assumptions, we built a 1-D horizontal model (from river
 158 channel to valley). Its flow field was calculated by the Boussinesq equation coupled
 159 with the DGM. We then built a 2-D vertical model (river channel to valley) the flow
 160 field of which was simulated by the Richard's equation. The diagrams of the 1-D and
 161 2-D vertical models for validation are shown in Fig. S3a and Fig. S3b, respectively.
 162 Thus, the horizontal 1-D model (Fig. S3a) represents a reduced 2-D horizontal model
 163 (Fig. S2a, the model we use in our main manuscript) while the vertical 2-D model (Fig.
 164 3b) represents a reduced 3-D model (described as Fig. S1), respectively. Our
 165 comparison neglects river sinuosity and the ambient groundwater flow gradient, but can
 166 efficiently prove the reliability of the vertical SWI assumption and usability of the
 167 DGM.



168
 169 **Figure S3.** Diagrams of the 1-D horizontal and 2-D vertical model that we used to
 170 evaluate the appropriateness of the assumptions of this study. Blue lines and dot indicate
 171 the river boundary, red lines and dot indicate the no-flow boundary

172
 173 The upper boundary of the 2-D model was higher than the peak river stage to avoid
 174 its submergence. The model length ($n\lambda$) has been shown to have no impact on the flow
 175 field (see S1), so the valley boundaries were set to be no-flow boundaries. The river
 176 boundary (SWI) of the 1-D model was assumed to be always vertical and the

177 displacement of the SWI was calculated by Eq. (1) coupled with DGM. The model
 178 length and parameters were the same as shown in Table 1 in the main manuscript. The
 179 model mesh was refined near the river boundary with the total number of elements
 180 being 4168. The time step was 0.1 d with a total run time of $2*t_d$.

181 For the vertical 2-D model corresponding to the 3-D model the flow field was
 182 simulated by the Richard's equation for unsaturated conditions which utilizes the van
 183 Genuchten equation for hydraulic properties (Van Genuchten, 1980):

$$184 \quad k = K S_e^{0.5} \left[1 - (1 - S_e^{1/m})^m \right]^2 \quad (S7a)$$

$$185 \quad S_e = \left[1 + \frac{1}{(\alpha \phi)^n} \right]^m \quad (S7b)$$

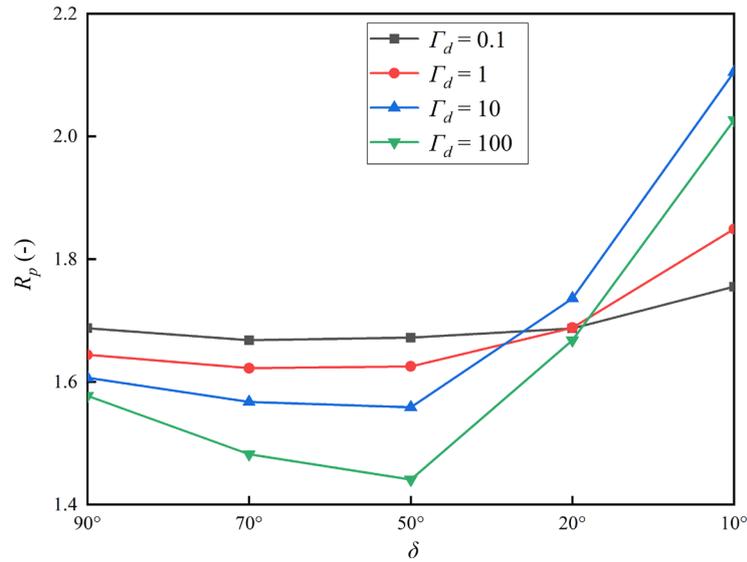
$$186 \quad S_e = \frac{S - S_r}{S_s - S_r} \quad (S7c)$$

$$187 \quad m = 1 - 1/n \quad (S7d)$$

188 Here, k is hydraulic conductivity with specific saturation [LT^{-1}]; S is soil moisture
 189 content [-]; S_r is residual water content [-], which is assumed to be 0.01; S_s is saturated
 190 water content [-], which is equal to the porosity (0.3); S_e is the effective saturation [-];
 191 α [L^{-1}] and n [-] are empirical coefficients, and assumed to be constant at 7.5 and 1.89,
 192 respectively while ϕ is pressure head [L]. The mesh was refined near the river boundary
 193 with a mesh size of 0.001 m, and a total number of mesh elements of 0.48 million. The
 194 total simulation time was $2*t_d$ with a time step of 0.1 d.

195 To evaluate the accuracy of our model approach for the scenarios tested in this
 196 study, the saturated hydraulic conductivity (parameterized by Γ_d) and bank slope angle
 197 were assigned to be the same as shown in Table 1 in the main manuscript. Fig. S4 shows
 198 the ratio of peak net flux of the 2-D to 1-D models ($R_p = Q_{max,2-D}^*/Q_{max,1-D}^*$). We could
 199 observe differences in net flux estimates between the 2-D vertical model (reduced 3-D)
 200 and the 1-D horizontal model even for vertical riverbank condition ($\delta = 90^\circ$), which
 201 resulted from the effect of the vadose zone (Liang et al., 2020). For sloping riverbank
 202 conditions these differences in net flux estimates remain albeit a change in magnitude.

203 Overall, the differences in peak net flux between horizontal 1-D and vertical 2-D
 204 models are mostly caused by the effect of the vadose zone. This implies that using the
 205 DGM has a minor influence on the prediction of the net flux in most of scenarios tested
 206 in this study.
 207



208
 209 **Figure S4.** Ratio of maximum net flux of 2-D vertical model to 1-D horizontal model
 210 $R_p = Q_{max,2-D}^*/Q_{max,1-D}^*$ and aquifer transmissivities.
 211

212 Figure S5 shows the relative difference in the bank storage between the 1-D
 213 horizontal and 2-D vertical model (R_s) for the scenarios tested in this study, which is
 214 calculated by:

$$215 \quad R_s = (S_{2-D} - S_{1-D}) / S_{2-D} \quad (S8a)$$

$$216 \quad S_{1-D} = \int h dx \quad (S8b)$$

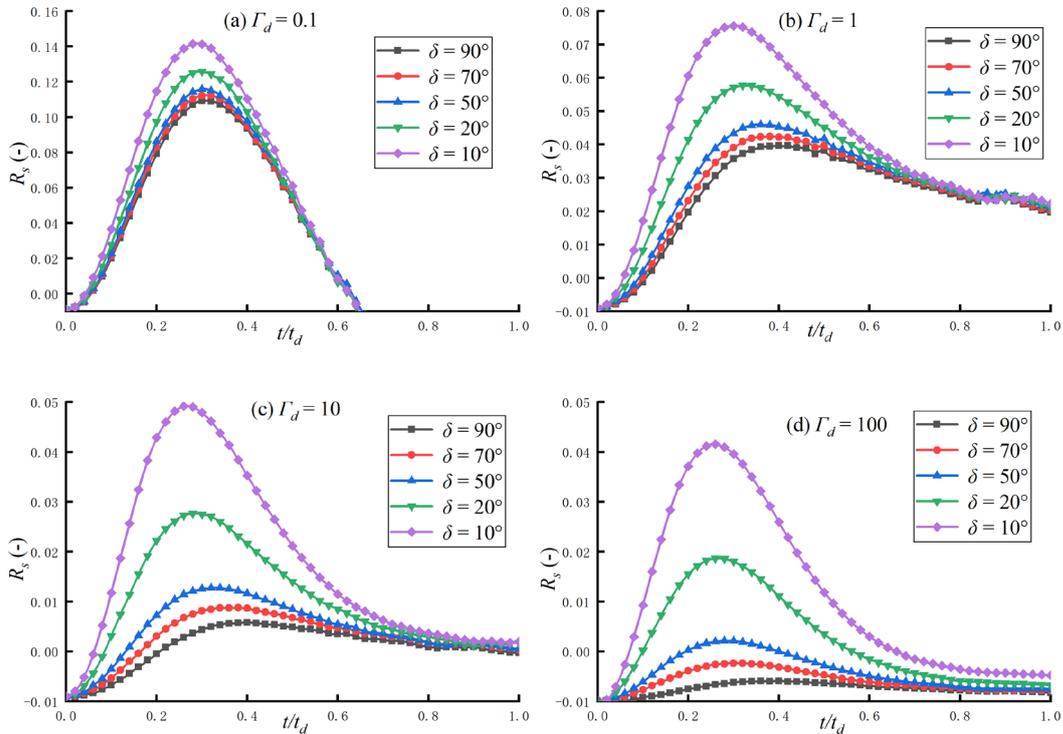
$$217 \quad S_{2-D} = \iint dx dz \quad \text{if } S_e = 1 \quad (S8c)$$

218 where S_{1-D} and S_{2-D} are bank storage of 1-D and 2-D models [L^2], respectively. Fig. S5
 219 indicates that the maximum difference in prediction of bank slope between 1-D and 2-
 220 D model were less than 15% under the condition with highest aquifer transmissivity
 221 ($\Gamma_d = 0.1$), however, the maximum R_p for $\Gamma_d = 0.1 \delta = 90^\circ$ was 11%, indicating the

222 model approach of this study has minor influence on the model performance.

223

224



225

226 **Figure S5.** Relative difference in bank storage between vertical 2-D and horizontal 1-
227 D model for various bank slope angle and aquifer transmissivity conditions.

228

229 While the using the Boussinesq equation neglects the influence of the vadose zone,
230 this equation as well as the assumption of vertical integrated distribution of hydraulic
231 head and solute have been widely used in the literature and proven adequate when
232 simulating sinuosity-driven HEF patterns (Boano et al., 2006; 2010., Cardenas. 2008;
233 2009a, b; Gomez-Velez et al., 2012; 2017, Kruegler et al., 2020). Despite that, Fig. S4
234 and Fig. S5 show that differences in the prediction of HEF patterns exist between the
235 Boussinesq model and Richard's model for all types of slope angle including a vertical
236 riverbank indicating a discrepancy between both mathematical approaches. This
237 discrepancy needs to be studied further to better understand the advantages and
238 limitations of either approach, e.g., in terms of computability or efficiency in predicting
239 HEF under various conditions.

240

241

242 S5 Metrics

243 We used the following dimensionless metrics to quantify the effects of bank
244 slope angle on the response of the dynamic hyporheic zone: (i) hyporheic exchange flux
245 along the river, (ii) in-valley penetration distance (i.e., the distance the river water
246 penetrates into the aquifer), (iii) the area of the HZ (i.e., the area of the aquifer exposed
247 to river water), and (iv) RTD and flux-weighted relative RT of HZ water discharging
248 into the river. In this section, we briefly define and describe each of these terms.

249 S5.1 Hyporheic exchange flux

250 Exchange flux from the river to the HZ ($Q_{in, HZ}$) and from the aquifer to the river
251 ($Q_{out, HZ}$) was defined as:

$$252 \quad Q_{in, HZ}(t) = - \int_{\partial\Omega_{in, HZ}(t)} \mathbf{Q}(\mathbf{x}, t) \cdot \mathbf{n} ds \quad (S9a)$$

$$253 \quad Q_{out, HZ}(t) = - \int_{\partial\Omega_{out, HZ}(t)} \mathbf{Q}(\mathbf{x}, t) \cdot \mathbf{n} ds \quad (S9b)$$

254 where $\Omega_{in, HZ}(t)$ and $\Omega_{out, HZ}(t)$ correspond to the inflow and outflow boundaries along
255 the meander of interest (black line along the river boundary in Fig. S2a). The net flux
256 from the aquifer into the river ($Q_{net, HZ} = Q_{out, HZ} - Q_{in, HZ}$) can be expressed in
257 dimensionless terms following Gomez-Velez et al. (2017) using $Q_{in, HZ}^*(t) = Q_{in, HZ}$
258 $(t)/(K\bar{H}_s^2)$, $Q_{out, HZ}^*(t) = Q_{out, HZ}(t)/(K\bar{H}_s^2)$, and $Q_{net, HZ}^*(t) = Q_{net, HZ}(t)/(K\bar{H}_s^2)$. Note that
259 these dimensionless fluxes are proportional to the integrated head gradient between the
260 river stage and the adjacent aquifer along the river boundary.

261 S5.2 Hyporheic zone area

262 Dynamic changes of the river-aquifer interface and pressure distribution along
263 the SWI induce variations of the flow field and changes to the HZ as represented by
264 area (i.e., the aquifer area exposed to river water) and penetration distance (i.e., how far
265 river water travels into the aquifer) during the flood event. These are useful metrics for

266 assessing the opportunity for biogeochemical and geochemical reactions induced by
 267 hyporheic exchange. Here we use a geochemical definition of HZ proposed by Triska
 268 et al. (1989), that defines the HZ as the area within the alluvial valley that contains more
 269 than 50% stream water ($C(\mathbf{x}, t) > 0.5$). It can be calculated using

$$270 \quad A(t) = \iint a(\mathbf{x}, t) dx dy \quad (\text{S10a})$$

$$271 \quad a(\mathbf{x}, t) = \begin{cases} 1 & \text{if } C(\mathbf{x}, t) \geq 0.5 \\ 0 & \text{if } C(\mathbf{x}, t) < 0.5 \end{cases} \quad (\text{S10b})$$

272 where $A(t)$ [L^2] is the area of the HZ. The dimensionless area is then defined similar to
 273 Gomez-Velez et al. (2017) as $A^*(t) = A(t)/\lambda^2$ and the dimensionless variation of the HZ
 274 area relative to base flow conditions can be calculated by $A^{**}(t) = A^*(t) - A^*(0)$, where
 275 $A^*(0)$ is the initial area of HZ under baseflow condition.

276 S5.3 Penetration distance of the hyporheic zone

277 The maximum penetration distance $d(t)$ of river water into the HZ in the direction
 278 perpendicular to the axis of the river can be calculated by the maximum y coordinate of
 279 the HZ. Similar to Gomez-Velez et al. (2017), we focus on the evolution of the
 280 dimensionless term of $d^{**}(t) = d^*(t) - d^*(0)$, where $d^*(t) = d(t)/\lambda$.

281 S5.4 Residence time

282 The difference in mean residence time distribution between a sloping and a vertical
 283 riverbank model was calculated by $\mu_r^*(\mathbf{x}, t) = \log_{10}(\mu_{\tau-s}(\mathbf{x}, t)/\mu_{\tau-v}(\mathbf{x}, 0))$. $\mu_r^* < 0$ indicating
 284 that RT was overestimated in these areas when ignoring the bank slope while $\mu_r^* > 0$
 285 indicating the opposite. Furthermore, a representative value of the flux-weighted ratio
 286 of mean RT to mean RT under baseflow conditions along the river boundary is given
 287 by: $\mu_{out}^*(x, t) = \mathbf{n} \cdot \mathbf{Q}_{out}(x, t) \log_{10}(\mu_{\tau}(x, t)/\mu_{\tau}(x, 0))$, which indicates aquifer discharge of
 288 younger water with relatively short travel times (values smaller than zero) or older
 289 water with longer travel times within the alluvial aquifer as compared to the baseflow
 290 conditions.

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