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# A new criterion for determining the representative elementary volume of translucent porous media and inner contaminant

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Abstract. Representative elementary volume (REV) is essential for measuring and quantifying the effective parameters of a complex heterogeneous medium. To overcome the limitations of the existing REV estimation criteria, a new REV estimation criterion  $(\chi^i)$  based on dimensionless range and gradient calculation is proposed in this study to estimate the REV of a translucent material based on light transmission techniques. Three sandbox experiments are performed to estimate REVs of porosity, density, tortuosity, and perchloroethylene (PCE) plume using multiple REV estimation criteria. In comparison with  $\chi^i$ , previous REV estimation criteria based on the coefficient of variation  $(C_{\rm V}^i)$ , the entropy dimension (DI<sup>*i*</sup>) and the relative gradient error ( $\varepsilon_{g}^{l}$ ) are tested in REV quantification of translucent silica and the inner PCE plume to achieve their corresponding effects. Results suggest that the new criterion  $(\chi^i)$  can effectively identify the REV in the materials, whereas the coefficient of variation and entropy dimension  $(F = -2.01 \times 10^{-12} +$  $\frac{1}{\sqrt{2\pi} \times 1.50} e^{-\frac{(\text{REV}-4.35)^2}{2\cdot 1.50^2}}$  are not effective. The relative gradient error can make the REV plateau obvious, while random fluctuations make the REV plateau difficult to identify accurately. Therefore, the new criterion is appropriate for REV estimation of the translucent materials and inner contaminant. Models are built based on a Gaussian equation to simulate the distribution of REVs for media properties, whose frequency of REV is dense in the middle and sparse on both sides. REV estimation of the PCE plume indicates that a high level of porosity leads to a large value of mean and standard deviation for REVs of PCE saturation  $(S_0)$  and PCE–water interfacial area  $(A_{OW})$ . Fitted equations are derived from distribution

of REVs for the PCE plume related to  $d_{\rm m}$  (distances from

mass center to considered point) and  $d_{\rm I}$  (distances from injection position to considered point). Moreover, relationships between REVs of the PCE plume and  $S_{\rm o}$  are fitted using regression analysis. Results suggest a decreasing trend appears for  $S_{\rm o}$ -REV when  $S_{\rm o}$  increases, while  $A_{\rm OW}$ -REV increases with increasing  $S_{\rm o}$ .

# 1 Introduction

Modelling groundwater and contaminant (such as hazardous ions) transport in a subsurface environment is based on the premise that the micro-structure of an aquifer has a representative elementary volume (REV) (Wang et al., 2016; Lei and Shi, 2019). REV acts as a micro-scale characteristic, which is important for improving our understanding of materials, inner fluid flow, and other processes (Brown and Hsieh, 2000; Costanza-Robinson et al., 2011; Wu et al., 2017). Previous studies (Bai et al., 2016; Esfandiar et al., 2017) suggested that even the platinum-nanoparticle-catalyzed hydrogenation reactions and ion transport through Angström-scale slits in cell activity had apparent size effects, implying size effect is a common characteristic for many processes and materials. With the help of REV, a porous medium can be treated as a continuum medium (Brown and Hsieh, 2000; Kang et al., 2003; Müller and Siegesmund, 2010; Teruel and Rizwan-uddin, 2010; Hendrick et al., 2012; Wang et al., 2012; Ukrainczyk and Koenders, 2014; Kim and Mohanty, 2016; Gilevska et al., 2019). A conceptual representation of an "REV curve" (Brown and Hsieh, 2000), characterizing porosity  $(\theta)$  change with measured scale (L) increment, is



Figure 1. (a) Variable changes as a measured scale (L) increment in the conceptual curve (Costanza-Robinson et al., 2011); (b) scale effect and the cuboid image window geometry; (c) system device for acquisition of parameters (porosity and density, etc.) of translucent material; (d) the infinitesimal selected from translucent porous media packed in a 2D sandbox.

presented in Fig. 1a. Based on the characteristic of the REV curve in Fig. 1a, the REV curve can be divided into three regions. When the measured scale (Fig. 1b) is in region I, the porosity fluctuates drastically at small scales. As the measured scale size ranges between  $L_{\min}$  and  $L_{\max}$ , a flat plateau with a constant and steady value is encountered and the property is factored into its average value. Material property in spatial scales less than  $L_{\min}$  is spatially varied in portions with a small scale, which can be easily influenced by individual pores in micro-structure such as region I (Fig. 1a). Likewise, material property is allowed to drift to new values in the spatial scale above  $L_{\text{max}}$  due to additional morphological structures of large field heterogeneity (region III). As a matter of fact, the REV scale of region II can be derived between the small and spatially varied property in region I and large field variability in region III. However, the lower and upper boundaries  $L_{\min}$  and  $L_{\max}$  of the REV plateau are hard to identify in reality (Brown and Hsieh, 2000; Costanza-Robinson et al., 2011).

As technology advanced and progressed, non-destructive and non-invasive techniques of X-ray and gamma ray microtomography were utilized for micro-structure characteristic measurement in materials (Ghilardi, 1993; Brown and Hsieh, 2000; Niemet and Selker, 2001; Bob et al., 2008; Al-Raoush and Papadopoulos, 2010; Costanza-Robinson et al., 2011; Al-Raoush, 2012; Borges and Pires, 2012; Fernandes et al., 2012; Rozenbaum and Roscoat, 2014; Pereira Nunes et al., 2016; Piccoli et al., 2019). Generally, REV estimation for material properties, inner gas and fluid was also usually implemented by micro-visualization and scanning of Xrays and gamma rays in the laboratory (Brown and Hsieh, 2000; Razavi et al., 2007; Nordahl and Ringrose, 2008; Al-Raoush and Papadopoulos, 2010; Costanza-Robinson et al., 2011; Rozenbaum and Roscoat, 2014; Borges et al., 2018), while different criteria were utilized to quantify REV (Brown and Hsieh, 2000; San José Martínez et al., 2007; Nordahl and Ringrose, 2008; Costanza-Robinson et al., 2011). A lower boundary scale  $L_{\min}$  of REV was identified by means of entropy dimension (DI<sup><math>i</sup>) for eight soil samples (San José Martínez et al., 2007). Further, the REV scale of permeability for ripple-laminated sandstone intercalated with mudstone was estimated using the coefficient of variation  $(C_V^l)$ , where the REV scale is identified by the variability among the 10 samples to achieve an average REV scale (Nordahl and Ringrose, 2008). As a result, only one REV boundary was identified, and not every sample can be estimated effectively (Nordahl and Ringrose, 2008). More interestingly, REVs of material property (porosity), moisture saturation, and air-water interfacial areas in porous media were estimated by a criterion named the relative gradient error  $(\varepsilon_{\sigma}^{i})$  (Costanza-Robinson et al., 2011). REVs of permeability of translucent material, perchloroethylene (PCE) saturation, and PCE-water interfacial area also can be estimated using the relative gradient error (Wu et al., 2017). In summary, the REV estimation was made up of multiple kinds of criteria, while the REV identification effects of these criteria were not clear. What is more, these previous criteria estimate REV scale unsatisfactorily in that the beginning and end of the REV plateau cannot be identified simultaneously for translucent porous media based on the light transmission technique. Therefore, a new criterion which can identify the REV plateau accurately is needed.

In this study, a new criterion ( $\chi^i$ ) for REV estimation is proposed to identify the REV scale of the translucent silica and inner contaminant. Three PCE transport experiments are conducted in two-dimensional (2D) sandboxes to test the effect of different REV estimation criteria. Translucent silica is selected for associated REV analysis due to its extensive utilization in laboratory experiments for exploring groundwater flow and contaminant migration behavior in the micro-structure of a sandy aquifer (Niemet and Selker, 2001; Bob et al., 2008; Costanza-Robinson et al., 2011). Moreover, translucent silica is also an important material applied in numerous industries (Bouvry et al., 2016). In laboratory experiments, translucent silica is packed in 2D sandboxes where porosity, density, tortuosity, and PCE saturation are derived



Figure 2. (a) The system sandbox equipment of Experiment I; (b) the system sandbox equipment of Experiment II; (c) the system sandbox equipment of Experiment III.

by the light transmission technique (Fig. 1c). Porosity and PCE saturation are selected as the representative variables to explore corresponding REV estimation by different criteria, which is very essential and significant for REV identification. Previous criteria such as the coefficient of variation ( $C_V^i$ ), entropy dimension (DI<sup>*i*</sup>), the relative gradient error ( $\varepsilon_g^i$ ), and the new criterion- $\chi^i$  are tested in REV estimation. Associated effects are analyzed to achieve the best criterion of effective and appropriate quantification of REV.

# 2 Experiment procedure and method

# 2.1 Experiment

Three sandboxes (Fig. 2a–c) packed by a translucent silica medium are prepared in the laboratory to test different criteria of REV quantification. PCE is selected as a typical dense nonaqueous-phase liquid (DNAPL) contaminant used in experiments. A 2D sandbox is composed of three aluminum interior frames and two glass walls whose thickness is 1.6 cm. The dimensions of sandboxes used in Experiment I are 20 cm (width)  $\times$  15 cm (height), and the dimensions of Experiments II and III are 60 cm (width)  $\times 45 \text{ cm}$  (height). F40/50 and F20/30 mesh translucent silica sands are used for background material for Experiments I and II, while heterogeneous translucent silica with low porosity and permeability are packed in a sandbox for Experiment III. To make the translucent silica fully saturated by water in a flow field close to a natural groundwater environment (Erning et al., 2012), water flow at a flow velocity of  $0.5 \text{ m d}^{-1}$  is set from left to right in laboratory sandbox experiments (Fig. 2a-c). Water is restricted in a sandbox to the top and bottom layers of the sandbox packed with F70/80 mesh translucent silica as capillary barriers. A light source is placed behind the sandbox to make light penetrate through translucent media and capture emergent light intensity using a thermoelectrically air-cooled

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Experiment	Ι	II	III
Sandbox dimensions (cm)	$20 \times 15$	$60 \times 45$	$60 \times 45$
Background translucent silica sand	F40/50	F20/30	F20/30
Medium condition	Homogeneity	Homogeneity	Heterogeneity
Median grain diameter (mm)	0.36	0.72	0.72
Permeability (m <sup>2</sup> )	$4.25 \times 10^{-11}$	$1.35 \times 10^{-10}$	$1.35 \times 10^{-10}$
$V_{\rm PCE} ({\rm mL})$	9	32	40
Injection rate (mL min <sup><math>-1</math></sup> )	0.5	0.5	0.5

charge-coupled device (CCD) camera (Fig. 1c). Afterward, PCE is injected into sandboxes from the injection needle at a constant rate of  $0.5 \text{ mL min}^{-1}$  for the three experiments. Detailed experimental conditions are listed in Table 1.

## 2.2 Light transmission technique

By means of the light transmission technique (Fig. 1c), DNAPL and water saturation can be obtained rapidly and in real time, which greatly helps to explore the mechanism of groundwater flow and contaminant migration in porous media. When light passes through translucent materials of a given thickness, the emergent light intensity after the absorptive and interfacial losses can be expressed as (Niemet and Selker, 2001; Bob et al., 2008; Wu et al., 2017)

$$I_{\rm T} = C I_0 \left( \prod \tau_b \right) \exp\left(-\sum \alpha_a d_a \right), \tag{1}$$

where *a* represents the phase number; *b* represents the number of the interface between phase *a* and a+1;  $I_0$  is the original light intensity; *C* is a constant of correction for light emission and light observation;  $\tau_b$  is the transmittance when light penetrates from phase *a* to a + 1;  $\alpha_a$  is the absorption coefficient when light penetrates in phase *a*;  $d_a$  is the length of the light penetration path in phase *a*.

To derive the porosity, the 2D translucent porous medium should be only saturated by water. Consequently, the emergent light intensity can be expressed as (Niemet and Selker, 2001; Bob et al., 2008; Wu et al., 2017)

$$I_{\rm s} = C I_0 \tau_{\rm s,w}^{2k_0} \exp\left(-\alpha_{\rm s} k_{\rm s} d_{\rm s}\right),\tag{2}$$

where  $\tau_{s,w}^{2k_0}$  is the transmittance of the solid–water interface;  $\alpha_s$  is the solid particle absorption coefficient;  $d_s$  is the median diameter of the solid particles;  $k_0$  is the number of pores across the light penetration path; and  $k_s$  is the number of solid particles across the light penetration path.

If we arbitrarily select an infinitesimal element, its area  $A_o$  approaches zero ( $A_o \rightarrow 0$ ) from the 2D translucent porous media (Fig. 1d), and suppose the infinitesimal element with thickness  $L_T$  containing solid particles and pores can be regarded as being of lamellar structure (Fig. 1d), we can obtain the following relationships (Wu et al., 2017):

$$\theta A_{\rm o} L_{\rm T} = A_{\rm o} k_{\rm o} d_{\rm o},\tag{3}$$

$$k_{\rm s}d_{\rm s} + k_{\rm o}d_{\rm o} - L_{\rm T},\tag{4}$$

where  $d_0$  is the median diameter of pores;  $\theta$  is porosity.

Substituting Eqs. (3) and (4) into Eq. (2), the relationship between emergent light intensity and porosity can be achieved (Wu et al., 2017):

$$\ln I_{\rm s} = \beta + \theta \gamma, \tag{5}$$

where  $\beta = \ln\left(\frac{CI_s}{e^{\alpha_s d_s L_T}}\right)$  and  $\gamma = \ln(\tau_{s,w}^{\frac{2L_T}{d_0}} e^{\alpha_s L_T})$ .  $\beta$  and  $\gamma$  can be determined from experimental data, and then porosity can be obtained.

The density and tortuosity are derived as (Wu et al., 2018)

$$\rho = \theta \rho_{\rm w} + (1.0 - \theta) \rho_{\rm s},\tag{6}$$

$$\tau = 1 + \frac{\pi - 2}{\sqrt{\frac{\pi}{1 - \theta}}},\tag{7}$$

where  $\rho$  is the density of translucent porous media;  $\rho_w$  is the density of water;  $\rho_s$  is the density of solid particles; and  $\tau$  is tortuosity.

The saturation of DNAPL was quantified by the light transmission technique based on light passing through translucent materials (Niemet and Selker, 2001; Bob et al., 2008):

$$S_{\rm o} = \frac{\ln I_{\rm s} - \ln I_{\rm T}}{\ln I_{\rm s} - \ln I_{\rm oil}},\tag{8}$$

where  $S_0$  is the saturation of DNAPL;  $I_s$  is the light intensity after light penetration through translucent pores when all pores are fully saturated by water;  $I_{oil}$  is the light intensity when all pores are fully saturated by DNAPL; and  $I_T$  is the light intensity after penetration through translucent materials. After quantification of PCE saturation, the PCE–water interfacial area ( $A_{OW}$ ) can be obtained based on the method proposed by Wu et al. (2017), where the unit of  $A_{OW}$  is cm<sup>-1</sup>.

Emergent light intensity for the three experiments is captured by a thermoelectrically air-cooled CCD camera (Niemet and Selker, 2001; Bob et al., 2008). Every pixel with a small scale could be approximated as an infinitesimal element in a high-resolution image to apply light transmission techniques. As a consequence, porosity of translucent silica was derived using light transmission techniques through Eq. (5). The whole 2D translucent silica area was numerically discretized so that every cell had the uniform dimensions of  $0.015 \text{ m} \times 0.015 \text{ m}$ . The cuboid window (Fig. 1b) was utilized to quantify the variables (porosity, density, tortuosity, PCE saturation, PCE-water interfacial area) of every cell as the measured scale was increased. In detail, the measured cuboid window scale was increased from the center of each cell, and the associated value of the variable was calculated from the high-resolution porosity of 2D translucent silica derived by the light transmission technique. Observation cells were selected from the discretized cells (Fig. 3b), of which cells I-1-2, II-1-2, and III-1-2 belong to Experiments I-III, respectively.

To analyze the regularity of the REV distribution for the PCE plume, the mass center coordinate and the ganglia-topool ratio (GTP) of the PCE plume are quantified for Experiments I–III. The mass center coordinate and GTP are calculated as

$$X_{\rm m} = \frac{M_{10}}{M_{00}},\tag{9}$$

$$Z_{\rm m} = \frac{M_{01}}{M_{00}},\tag{10}$$

$$\text{GTP} = \frac{V_{\text{ganglia}}}{V_{\text{pool}}},\tag{11}$$

where  $X_{\rm m}$  is the *x* coordinate of the mass center for the PCE plume;  $Z_{\rm m}$  is the *z* coordinate of the mass center for the PCE plume; GTP is the ganglia-to-pool ratio, which equals the ratio of  $V_{\rm ganglia}$  to  $V_{\rm pool}$ ;  $V_{\rm ganglia}$  is the PCE volume under ganglia state;  $V_{\rm pool}$  is the PCE volume under pool state;  $M_{00}$ ,  $M_{10}$ , and  $M_{01}$  are computed using the definition of the spatial moment  $(M_{ij})$ ,  $M_{ij} = \int_{x_0}^{x_1} \int_{z_0}^{z_0} \theta(x, z) S_0(x, z, t) x^i z^j d_x d_z$ ;  $x_0$  and  $z_0$  are minimum limits of the *x* axis and *z* axis;  $x_1$  and  $z_1$  are maximum limits of the *x* axis and *z* axis;  $\theta(x, z)$  is the porosity at point (x, z);  $S_0(x, z, t)$  is PCE saturation of point (x, z)

#### 2.3 Criteria of REV quantification

at time t.

The REV is defined as the volume range in which all material characteristics are factored into the average and associated values approach single and constant (Brown and Hsieh, 2000). In the range of REV, the value of one associated property will meet the condition

$$\frac{\partial Y(L_i)}{\partial L}|_{L_i=L_0} = 0, \tag{12}$$

where  $Y(L_i)$  is the value of an associated property when the system scale is  $L_i$ ;  $L_i$  is the value of the system scale;  $L_0$  is the scale range of REV,  $L_{\min} < L_0 < L_{\max}$ ;  $L_{\max}$  is the upper boundary of REV;  $L_{\min}$  is the lower boundary of the REV scale. According to Eq. (12), when the measured scale ( $L_i$ ) reaches the REV range,  $\frac{\partial Y(L_i)}{\partial L} \rightarrow 0$ . As a matter of fact, most previously used criteria were applied to estimate REV based on this requirement (Brown and Hsieh, 2000; San José Martínez et al., 2007; Nordahl and Ringrose, 2008; Costanza-Robinson et al., 2011).

To evaluate the REV of porosity, the coefficient of variation  $(C_V^i)$  is utilized to estimate the variability (Nordahl and Ringrose, 2008):

$$C_{\rm V}^i = \frac{\hat{s}_i}{\overline{\varphi_i}},\tag{13}$$

where *i* is the cuboid window (Fig. 1b) increment number;  $\varphi$  is the measured variable, such as porosity;  $\hat{s}_i$  is the standard deviation of sub-grid variables in different measured volumes or scales;  $\overline{\varphi_i}$  is the arithmetic average of the variable values in the sub-grids. When the number of sub-grids (*N*) is less than 10, a correction is utilized to replace Eq. (13). According to Nordahl and Ringrose (2008),  $0 < C_V^i < 0.5$  is defined as homogeneous and  $C_V^i = 0.5$  can be used as a criterion to identify the REV scale.

Similarly, for porosity of translucent silica, entropy dimension ( $DI^i$ ) is utilized for REV analysis and estimation (San José Martínez et al., 2007), which is defined as

$$\mathrm{DI}^{i} \approx \frac{\sum_{j=1}^{m(i)} \mu_{j}\left(L_{\varepsilon}\right) \log \mu_{j}\left(L_{\varepsilon}\right)}{\log L_{\varepsilon}},$$
(14)

where  $L_{\varepsilon}$  is the scale of the sub-grid; " $\approx$ " indicates the asymptotic equivalence as  $L_{\varepsilon} \rightarrow 0$  (San José Martínez et al., 2007); *j* is the ordinal number of sub-grids in measured cuboid windows (Fig. 1b) of increment number i; m(i) is the total number of sub-grids in measured cuboid windows (Fig. 1b) of increment number i;  $\mu_i(\varepsilon)$  is the proportion of the variable of sub-grid j in the whole measured cuboid window i. The right-hand side of Eq. (14) is the simplification of Shannon entropy of all sub-grids, in which  $DI^{i}$  can be considered the average of logarithmic values of the variable distribution weighted by  $\mu_i(L_{\varepsilon})$  to quantify the degree of medium heterogeneity. Using Eq. (14), a series of values of DI<sup>*i*</sup> (i = 1, 2, 3...) are obtained for each measured cuboid window (Fig. 1b) of increment number *i*. For estimation of the REV in a porous medium, the relative increment of entropy dimension and associated criterion of REV identification are, respectively, expressed as

$$\mathbf{RI}^{i} = |\frac{\mathbf{DI}^{j} - \mathbf{DI}^{j-1}}{\mathbf{DI}^{j-1}}| \times 100,$$
(15)



Figure 3. (a) The emergent light intensity, porosity, permeability, and tortuosity of 2D translucent silica sand for Experiments I–III; (b) the PCE saturation of Experiments I–III during 0-1523 min and observation cells.

 $\mathbf{RI}^i \le 0.2 \mathbf{CV}_{\mathrm{DI}},\tag{16}$ 

where  $CV_{DI}$  is the coefficient of variation of  $DI^i$  series (i = 1, 2, 3...), which is calculated through  $CV_{DI} = (\hat{s}_{DI}/\overline{DI}) \times 100$ ;  $\overline{DI}$  is the mean value of the  $DI^i$  series;  $\hat{s}_{DI}$  is the standard deviation of the  $DI^i$  series.

To achieve the REV for multiple system variables, such as porosity, moisture saturation and air-water interfacial areas in an unsaturated porous medium, a criterion named the relative gradient error was applied (Costanza-Robinson et al., 2011):

$$\varepsilon_{g}^{i} = |\frac{\varphi^{i+1} - \varphi^{i-1}}{\varphi^{i+1} + \varphi^{i-1}}| \frac{1}{\Delta L},$$
(17)

where  $\varepsilon_g^i$  is relative gradient error;  $\Delta L$  is the measured cuboid window size increment length for REV estimation. Usually,  $\varepsilon_g^i$  less than 0.2 (Costanza-Robinson et al., 2011) is utilized to identify REV sizes.

According to the requirement in Eq. (12), a new criterion based on the required condition of REV is proposed to estimate the REV range of the translucent silica in this study:

$$\chi^{i} = \frac{|\delta^{i+1} - \delta^{i-1}|}{\delta^{i} \Delta L},\tag{18}$$

where  $\delta^i$  is the dimensionless range,  $\delta^i = \frac{\varphi(L_i)_{\max} - \varphi(L_i)_{\min}}{\varphi(L_i)}$ ;  $\varphi(L_i)_{\max}$  is the maximum value of the variable on the volume scale  $L_i$ ;  $\varphi(L_i)_{\min}$  is the minimum value of the variable on the volume scale  $L_i$ ; and  $\overline{\varphi(L_i)}$  is the mean value of the variable on the volume scale  $L_i$ . Brown and Hsieh (2000) suggested  $\delta^i = \frac{\varphi(L_i)_{\max} - \varphi(L_i)_{\min}}{\varphi(L_i)} \ll 1$  can be used for REV estimation. In fact, the calculated value of  $\delta^i$  mostly is less than 1, while  $\delta^i \ll 1$  is hard to use to identify the REV scale for realistic materials, such as the translucent silica used in this study. The value limit of  $\chi^i$  used for REV estimation also is explored in this study.

In this study, criteria for the coefficient of variation  $(C_V^i)$ , entropy dimension  $(DI^i)$ , the relative gradient error  $(\varepsilon_g^i)$  and the new criterion  $(\chi^i)$  are all applied to REV estimation for porosity and PCE saturation. Corresponding REV plateau identification effects are compared to select the best criterion for REV quantification.

# 3 Results and discussion

# 3.1 REV identification effect of different criteria

#### 3.1.1 The coefficient of variation

Emergent light intensity distributions of translucent silica for the three experiments, which had been fully saturated by water, were obtained by a thermoelectrically air-cooled CCD camera (Niemet and Selker, 2001; Bob et al., 2008). The porosity, density, tortuosity and PCE saturation for the three experiments are derived by the light transmission technique as shown in Fig. 3a and b. The PCE spreads from the injecting point shaped like a drop of water at  $t = 5 \min$  (Fig. 3b). In 2D sandboxes for the three experiments, the PCE plume infiltrates translucent silica sands and reaches the bottom after  $t = 80 \min$ .

Porosity and PCE saturation variation curves of all observation cells with increasing measured cuboid window scale are shown in Fig. 4a and b. However, for all observation cells from translucent silica, the REV plateaus were not certain to be objectively judged visually, which made REV plateaus hard to identify effectively by original variation curves for porosity and PCE saturation (Fig. 4a and b).

To make the REV plateau more explicit, different criteria of REV quantification are utilized. The coefficient of variation  $(C_{\rm V}^{l})$  of porosity and PCE saturation fluctuating with increasing measured cuboid window size is shown in Fig. 4. The measured cuboid window scale is limited to the dimensions of cells in discretization of 2D translucent silica. The observation cells show various characteristics of variation tendency for the coefficient of variation  $(C_V^i)$ . The  $\theta$  and  $S_0$  variation curves of the coefficient of variation  $(C_V^i)$  for all observation cells do not reach stable values like those shown in Fig. 4a and b; the beginning of the REV flat plateau is not easy to identify, and the coefficient of variation  $(C_{y}^{i})$  is not suitable for REV estimation. According to the heterogeneity definition by Corbett and Jensen (1992), the heterogeneity of materials is defined by  $C_V^i$  magnitude such that  $0 < C_V^i < 0.5$ is classed as a homogeneous medium,  $0.5 < C_V^i < 1.0$  is classed as a heterogeneous medium and  $1.0 < C_V^i$  is classed as a strong heterogeneous medium. For the coefficient of variation  $(C_V^l)$  magnitude in Fig. 4a and b, the  $C_V^l$  values are all below 0.5, so that the criterion  $C_{\rm V}^i = 0.5$  is unable to identify the REV scale for translucent silica.

#### 3.1.2 Entropy dimension

Entropy dimension  $(DI^{i})$  is utilized by San José Martínez et al. (2007) for multifractal analysis of a porous medium porosity and REV estimation. In this study, entropy dimension (DI<sup>i</sup>) is tested to avoid an unclear REV plateau in porosity curves. The entropy dimension  $(DI^{i})$  of porosity is calculated by Eq. (14). Variation curves of entropy dimension (DI<sup>i</sup>) for all observation cells (Fig. 2a) are presented in Fig. 4. The curves of entropy dimension  $(DI^{i})$  of porosity and PCE saturation generally result in the increasing trend curves, which makes REV estimates become very difficult and invalid. Entropy dimension (DI<sup>*l*</sup>) was quickly increased with increasing measured cuboid window size. Compared to the coefficient of variation  $(C_V^i)$  of porosity and PCE saturation, entropy dimension (DI<sup>i</sup>) increased step by step without opposite fluctuation tendency in the variation curves as the length scale of the measured cuboid window increased simultaneously. In general, the REV plateau in region II (Fig. 1a) of porosity is not obvious for the entropy dimension (DI<sup>*i*</sup>) curves of all observation cells from the three experiments, which suggests REV scales are difficult to identify for translucent silica using entropy dimension  $(DI^{i})$  by the light transmission technique.

#### 3.1.3 The relative gradient error

The relative gradient error  $(\varepsilon_g^i)$  of porosity and PCE saturation is calculated by Eq. (17). The variation of  $\varepsilon_g^i$  at different measured cuboid window scales is shown in Fig. 4 for all observation cells in the 2D translucent silica. For all  $\varepsilon_g^i$  curves at observation cells from experiments, the REV plateaus in region II (Fig. 1a) are clearer than the variation curves based



**Figure 4. (a)** The change in porosity ( $\theta$ ), associated coefficient of variation ( $C_V^i$ ), entropy dimension ( $DI^i$ ), relative gradient error ( $\varepsilon_g^i$ ), and new criterion- $\chi^i$  for observation cells as the cuboid window scale (*L*) increases; (**b**) the change in PCE saturation ( $S_0$ ), associated  $C_V^i$ ,  $DI^i$ ,  $\varepsilon_g^i$ , and  $\chi^i$  for observation cells as the cuboid window scale (*L*) increases.

on the criteria of  $C_{\rm V}^i$  and  ${\rm DI}^i$ . Apparently, erratic variations of the relative gradient error  $(\varepsilon_g^i)$  at small measured cuboid window scales are observed for all  $\varepsilon_g^i$  curves as the characteristic of REV region I in Fig. 1a. When the measured cuboid window scale further increases for all observation cells, the variability and magnitude of the relative gradient error  $(\varepsilon_{\alpha}^{i})$ decrease well and factor into the average, which can be identified as the REV plateau in region II (Fig. 1a). The relative gradient error  $(\varepsilon_g^i)$  makes the REV plateau quantification convenient for all observation cells. At the measured cuboid window size above the REV plateau,  $\varepsilon_{g}^{i}$  curves result in large variability for observation cells I-1-2. These findings suggest that the relative gradient error  $(\varepsilon_g^i)$  can make the REV plateau more obvious, which greatly contributes to convenient and applicable REV quantification for translucent silica by the light transmission technique. However, random fluctuations exist in  $\varepsilon_{g}^{i}$  curves visually, which makes the REV plateau difficult to identify accurately.

# **3.1.4** The new criterion $(\chi^i)$

 $\chi^i$  of porosity and PCE saturation changing with measured cuboid window size is shown in Fig. 4. Like region I

(Fig. 1a), erratic and random fluctuations appear at small measured cuboid window sizes, and  $\chi^i$  increases with the increase in the measured cuboid window size. When the measured scale increases, the values of  $\chi^i$  for all observation cells appear to decrease quickly and tend to steady rapidly, which exhibits the characteristic of the REV plateau as the measured scale reaches region II. The  $\chi^i$  for observation cells restores the erratic variation state of the increasing trend after the measured cuboid window size exceeds the REV plateau. As shown in the variation curves of  $\chi^i$  for all observation cells, the beginning of the REV flat plateaus can be identified easily, suggesting  $\chi^i$  is more convenient and reliable than other methods for REV estimation. All observation cells show similar variation curves of  $\chi^i$  that low value intervals are quite apparent, indicating that  $\chi^i$  is very effective at making the REV plateau obvious for the translucent silica used in this study. Using the criterion of  $\chi^i$ , the REV plateau of region II is flat, which is easily identified compared with other criteria for observation cells (Fig. 4a and b).

#### **3.2 REVs of material properties**

Based on the REV plateau identifications using the coefficient of variation ( $C_V^i$ ), entropy dimension ( $DI^i$ ), the relative gradient error ( $\varepsilon_g^i$ ), and the proposed new criterion  $\chi^i$  in Fig. 4a and b, the new criterion  $\chi^i$  appears to be the most appropriate criterion for REV plateau identification. Even though the relative gradient error ( $\varepsilon_g^i$ ) can also make the REV plateau obvious, various random fluctuations weaken the method and reduce the associated accuracy. Therefore, REVs of porosity, density, tortuosity and the PCE plume are estimated using the new criterion  $\chi^i$  in the following study.

In fact, a large number of discretized cells in the 2D translucent silica for the three experiments are quantified using the new criterion  $\chi^i$ , which is convenient for examining the regularities for REV sizes and related factors. Using the new criterion  $\chi^i$ , the REV estimation is conducted based on Eq. (18). Figure 5a shows minimum REV sizes of porosity, density and tortuosity quantified by  $\chi^i$  for all cells of the three experiments. Associated statistical analysis for REVs is illustrated in Fig. 5b, where circular points represent frequency and triangular points represent cumulative frequency. Frequency of REVs is dense in the middle and sparse on both sides, so the distribution of REVs can be fitted by a Gaussian equation:

$$F = \omega + \frac{1}{\sqrt{2\pi\epsilon}} e^{-\frac{(\text{REV}-\nu)^2}{2\epsilon^2}},$$
(19)

where F is the frequency of REV;  $\omega$ ,  $\epsilon$  and  $\nu$  are fitted parameters of the model.

After regression analysis, the derived models for REV frequency are listed in Table 2. The coefficients of determination ( $R^2$ ) of models for REVs of porosity and density for the three experiments all exceed 0.85.  $R^2$  for the REV of tortuosity for the three experiments exceeds 0.76. Moreover, the computed cumulative frequency based on models fits cumulative frequency from experimental results well in Fig. 5b.

The minimum REV size frequency of porosity appears as a peak between 4.0 and 5.0 mm for Experiment I. As the minimum REV size of porosity increases, the corresponding frequency continuously decreases. Further, a smooth convex shape of cumulative frequency is observed for the minimum REV size of porosity (Fig. 5b). Most minimum REV sizes of translucent silica are distributed within 0.0-7.0 mm. For density of translucent silica sand, the associated REV frequency has high values between 2.0 and 3.0 mm. For the REV sizes of tortuosity, minimum REV sizes are distributed within 0.0-6.0 mm. Compared with Experiment I (F40/50 mesh translucent silica sand), the frequency of REVs for Experiment II (F20/30 mesh translucent silica sand with larger porosity) shows a flat shape and has a larger value of standard deviation, especially for REVs of porosity. Figure 5b shows that translucent silica with larger porosity will achieve a border distribution of minimum REV size distribution compared to translucent silica with a relatively lower porosity. Moreover, **Table 2.** The fitted parameters of the models of frequency for REVs of porosity, density, and tortuosity.

Experin	nent	Ι	II	III
	ω	$-2.11 \times 10^{-4}$	$-1.45  imes 10^{-3}$	$7.63  imes 10^{-4}$
0 DEV	$\epsilon$	1.73	3.45	3.18
0-KEV	ν	4.35	7.90	6.50
	$R^2$	0.938	0.924	0.907
-	ω	$-6.51 \times 10^{-4}$	$-2.51 \times 10^{-4}$	$1.51 \times 10^{-3}$
o DEV	$\epsilon$	1.08	1.66	2.40
p-REV	ν	2.89	2.97	3.70
	$R^2$	0.967	0.990	0.859
	ω	$-3.36 \times 10^{-4}$	$-2.04 \times 10^{-4}$	$1.29 \times 10^{-3}$
7 DEV	$\epsilon$	1.39	2.15	1.20
ι-INE V	ν	3.65	4.20	1.05
	$R^2$	0.769	0.875	0.919

 $\theta$  represents porosity,  $\rho$  represents density, and  $\tau$  represents tortuosity;  $\omega$ ,  $\epsilon$ , and  $\nu$  are fitted parameters of the model.

the frequency of REVs of porosity and permeability for Experiment III (background material is F20/30 mesh translucent silica sand with larger porosity; five lenses with lower porosity are packed in a sandbox to create heterogeneity) is similar to the frequency of REVs for Experiment II. However, the frequency of  $\tau$ -REV for Experiment III is different from the frequency of  $\tau$ -REV for Experiment II under a homogeneous condition. The mean REV sizes of porosity, density and tortuosity for Experiment I are 4.35, 2.89 and 3.65 mm, respectively. All mean REV sizes of these variables for Experiment II are larger than REVs of Experiment I, whose corresponding mean REV sizes are 8.05, 2.97 and 4.30 mm. These results suggest translucent porous media with higher porosity lead to larger values of mean and standard deviation for REV sizes.

# **3.3** REVs of S<sub>0</sub> and A<sub>OW</sub> for the PCE plume

Based on the new criterion  $\chi^i$  and the light transmission technique, the real-time distributions of So-REV and AOW-REV for the PCE plume can be obtained over the entire experimental period. The minimum REV sizes of  $S_0$  and  $A_{OW}$  obtained using the new criterion  $\chi^i$  are shown in Fig. 6a and b. When PCE migrates in the sandbox, the REV of the PCE plume is changed over time (Fig. 6). The REVs of the PCE plume for Experiment I mostly are lower than the REVs of the PCE plume for Experiments II and III. Moreover, when heterogeneous porous media are packed in the sandbox, the REV distribution of Experiment III becomes more heterogeneous compared with the REV distribution of Experiment II under a homogeneous condition. Based on REV distributions of the PCE plume for the three experiments, statistical analysis is conducted to explore the regularity of the REV distribution for the PCE plume.



Figure 5. (a) The distributions of minimum REV sizes of porosity, sand density, and tortuosity for Experiments I–III; (b) the frequency of minimum REV sizes of experiments and fitted models.

The mass center coordinates of the PCE plume, GTP and plume area are shown in Fig. 7a. The values of  $X_m$ ,  $Z_m$  and GTP for Experiments II and III are higher than the  $X_m$ ,  $Z_m$  and GTP of Experiment I (lower porosity). Moreover, the plume area of Experiment II is larger than the plume of Experiment I. When packed material is heterogeneous, the plume area of PCE is increased further for Experiment III. Besides, the mean and standard deviation of REVs of the PCE plume during 0–1523 min are derived by statistical analysis (Fig. 7a). Compared with REVs of the PCE plume for Experiment I, Experiment II (F20/30 mesh translucent silica sand with higher porosity) has larger values of the mean and standard deviation of REVs. The mean value of  $A_{OW}$ -REV for Experiment III is much higher than  $A_{OW}$ -REV for Experiments I and II.



Figure 6. (a) The distributions of  $S_0$ -REV sizes during 0–1523 min for Experiments I–III; (b) the distributions of  $A_{OW}$ -REV sizes during 0–1523 min for Experiments I–III.

The average value of REVs ( $\overline{\text{REV}}$ ) and the associated distance ( $d_{\rm m}$ ) from the mass center to the corresponding cells contained in the PCE plume at t = 1523 min are presented in Fig. 7b. Regression analysis is performed for average REVs of the PCE plume and  $d_{\rm m}$ , where fitted models and the associated  $R^2$  for Experiments I–III are listed in Table 3. Simultaneously, the fitted equations between  $\overline{\text{REV}}$  and  $d_{\text{I}}$  (the distance from injection point to cell contained in the PCE plume) are also derived by regression analysis. From the results in Fig. 7a,  $\overline{\text{REV}}$  of  $S_{\text{o}}$  and  $A_{\text{OW}}$  appear as a peak and then decrease with increasing  $d_{\text{m}}$  and  $d_{\text{I}}$  for Experiment I.  $\overline{\text{REV}}$  of  $S_{\text{o}}$  and  $A_{\text{OW}}$  for Experiment I all firstly increase and



Figure 7. (a) The mass center coordinate of the PCE plume, GTP, plume area and the mean and standard deviation of  $S_0$ -REV and  $A_{OW}$ -REV during 0–1523 min; (b) the change in average REV size as the distance  $d_{I}$ ,  $d_{m}$  increases and the fitted relationship between REV sizes and  $S_0$  for Experiments I and II.

then decrease with the increasing  $d_m$  and  $d_I$ . However,  $\overline{\text{REV}}$  of  $S_o$  presents an apparent decreasing tendency as  $d_m$  and  $d_I$  increase for Experiment II, and  $\overline{\text{REV}}$  of  $A_{OW}$  just slightly increases first and then decreases for Experiment II. In addition, the value of  $A_{OW}$ -REV mostly is higher than the value of  $S_o$ -REV for the three experiments. Compared with the  $R^2$  of the fitted relationship between average REVs of the PCE plume and  $d_m$ ,  $d_I$  for Experiments I and II, the values of  $R^2$  achieved by Experiment III are much lower (Table 3).

Besides, the relationship between REVs and PCE saturation is fitted by regression analysis, where the fitted equation and  $R^2$  for the three experiments are listed in Table 4 and Fig. 7b. With increasing PCE saturation, REVs of  $S_0$  show a declining trend for the three experiments. However, the REV of  $A_{OW}$  increases when  $S_0$  increases for all three experiments (Fig. 7b). On the other hand, the REV of  $S_0$  for Experiment II is higher than the corresponding REV for Experiment I, while Experiments I and II have similar values of  $A_{OW}$ -REV (Fig. 7b). Moreover, the REVs of  $S_0$  and  $A_{OW}$  for Experiment III are higher than the REVs of  $S_0$  and  $A_{OW}$  for Experiments I and II. These results suggest higher porosity will lead to a high value of  $S_0$ -REV and the relationship between

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Experiment		d <sub>m</sub>	dI		
Ι	So-REV	$\overline{\text{REV}} = -1.67 \times 10^{-3} d_{\text{m}}^2 + 0.193 d_{\text{m}} + 2.72$ $(R^2 = 0.807)$	$\overline{\text{REV}} = -1.97 \times 10^{-3} d_{\text{I}}^2 + 0.245 d_{\text{I}} + 1.12$ $(R^2 = 0.832)$		
	A <sub>OW</sub> -REV	$\overline{\text{REV}} = -6.10 \times 10^{-4} d_{\text{m}}^2 + 5.82 \times 10^{-2} d_{\text{m}} + 7.20$ $(R^2 = 0.401)$	$\overline{\text{REV}} = -1.47 \times 10^{-3} d_{\text{I}}^2 + 0.205 d_{\text{I}} + 1.84$ $(R^2 = 0.733)$		
П	<i>S</i> <sub>o</sub> -REV	$\overline{\text{REV}} = -4.08 \times 10^{-5} d_{\text{m}}^2 + 1.50 \times 10^{-2} d_{\text{m}} + 7.54$ $(R^2 = 0.655)$	$\overline{\text{REV}} = -3.94 \times 10^{-5} d_{\text{I}}^2 + 7.80 \times 10^{-3} d_{\text{I}} + 8.50$ (R <sup>2</sup> = 0.327)		
	A <sub>OW</sub> -REV	$\overline{\text{REV}} = -1.92 \times 10^{-5} d_{\text{m}}^2 + 4.47 \times 10^{-3} d_{\text{m}} + 9.46$ (R <sup>2</sup> = 0.616)	$\overline{\text{REV}} = -1.92 \times 10^{-5} d_{\text{I}}^2 + 4.47 \times 10^{-3} d_{\text{I}} + 9.46$ (R <sup>2</sup> = 0.616)		
III	S <sub>o</sub> -REV	$\overline{\text{REV}} = -6.06 \times 10^{-6} d_{\text{m}}^2 + 2.27 \times 10^{-3} d_{\text{m}} + 7.76$ $(R^2 = 0.153)$	$\overline{\text{REV}} = 1.69 \times 10^{-5} d_{\text{I}}^2 - 1.21 \times 10^{-2} d_{\text{I}} + 9.62$ $(R^2 = 0.236)$		
	A <sub>OW</sub> -REV	$\overline{\text{REV}} = -8.71 \times 10^{-6} d_{\text{m}}^2 + 5.66 \times 10^{-3} d_{\text{m}} + 11.5$ (R <sup>2</sup> = 0.115)	$\overline{\text{REV}} = -1.50 \times 10^{-5} d_{\text{I}}^2 + 7.88 \times 10^{-3} d_{\text{I}} + 11.4$ $(R^2 = 0.150)$		

**Table 3.** The fitted equations between average values of REV and  $d_{I}$ ,  $d_{m}$ .

 $\overline{\text{REV}}$  is the average value of REV size,  $d_m$  is the distance from the mass center of the PCE plume to the cell contained in the PCE plume, and  $d_I$  is the distance from the injection point to the cell contained in the PCE plume.

Table 4. The fitted relationship between REV and So.

Experiment	So-REV	$A_{\rm OW}$ -REV
I	$REV = -2.13 \times \ln S_0 + 0.876$ $(R^2 = 0.466)$	$REV = 2.27e^{2.70 \times S_0}$ $(R^2 = 0.366)$
II	$REV = -0.961 \times \ln S_0 + 1.09$ $(R^2 = 0.415)$	$REV = 1.70e^{3.30 \times S_0}$ $(R^2 = 0.500)$
III	$REV = -1.40 \times \ln S_0 + 3.96$ $(R^2 = 0.538)$	$REV = 2.05e^{3.22 \times S_0}$ $(R^2 = 0.573)$

REVs of the PCE plume and  $d_{\rm m}$ ,  $d_{\rm I}$ .  $S_{\rm o}$ -REV and  $A_{\rm OW}$ -REV are increased under heterogeneous conditions.

#### 4 Conclusions

In this study, a new criterion  $\chi^i$  is proposed to identify the REVs of translucent porous media and inner contaminant transformation based on previous criteria. The REV plateaus of observation cells selected from the three experiments of PCE transport are hard to judge visually from the porosity and PCE saturation curves. From the REV identification effects of different criteria, the REV flat plateau is difficult to identify by the coefficient of variation  $(C_{\rm V}^i)$  and entropy dimension (DI<sup>*i*</sup>), indicating the coefficient of variation ( $C_V^i$ ) and entropy dimension (DI<sup>i</sup>) are not suitable for REV estimation of translucent porous media. The relative gradient error  $(\varepsilon_{g}^{i})$  can make REV plateaus of all kinds of translucent silica explicit in variation curves, but random fluctuations weaken REV plateau identification. In comparison with these previous criteria, the beginning and ending of the REV flat plateaus could be easily and directly identified in the

curves based on the new criterion  $\chi^i$ , suggesting the new criterion  $\chi^i$  is more convenient and effective for REV estimation. In this study, REVs of porosity, density, tortuosity, and the PCE plume are estimated using the new criterion  $\chi^i$ .

Statistical results of minimum REV scales quantified by the new criterion  $\chi^i$  reveal that cumulative frequencies of porosity, density and tortuosity all have smooth convex shapes. Models based on the Gaussian equation are built for the distribution of REVs of porosity, density and tortuosity, whose porous media with larger porosity lead to larger values of mean and standard deviation for REV sizes of media properties. For REVs of the PCE plume, the result suggested larger porosity leads to a larger value of mean and standard deviation. Regression analysis is performed to study the regularity for distribution of REVs, where the fitted relationship between REVs and  $d_{\rm m}$ ,  $d_{\rm I}$  is derived for the PCE plume. REV of  $S_0$  and  $A_{OW}$  firstly increases and then decreases with the increasing  $d_{\rm m}$  and  $d_{\rm I}$  for Experiment I, whose sandbox is packed with translucent porous media with relatively lower porosity. However,  $\overline{\text{REV}}$  of  $S_0$  and  $A_{OW}$  directly decreases with the increment of  $d_{\rm m}$  and  $d_{\rm I}$  when porosity becomes larger for Experiment II. The values of  $R^2$  of the fitted relationship between average REVs of the PCE plume and  $d_{\rm m}$ ,  $d_{\rm I}$  for Experiment III are much lower under heterogeneous conditions. Significantly, the REV size of  $S_0$  presented a decreasing trend as  $S_0$  increases, while an increasing tendency appeared for an REV size of  $A_{OW}$ . Through regression analysis, the fitted equations between REVs of the PCE plume and PCE saturation are derived for the three experiments. Implications of these findings are essential for quantitative investigation of the scale effect of porous media and contaminant transformation. The fluid migration and transform in porous media can be simulated accurately according to the REV estimation results using the light transmission technique and the appropriate criterion  $\chi^i$ .

*Code and data availability.* The codes and data for this paper are available by contacting the corresponding author at jfwu@nju.edu.cn.

*Author contributions.* MW and JiaW were responsible for the conceptualization, methodology, and writing of the paper. JicW was responsible for its conceptualization. BXH was responsible for the conceptualization and writing of the paper.

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